

**KARLSRUHER INSTITUT FÜR TECHNOLOGIE (KIT)**

**Extensive Air Shower Simulation  
with CORSIKA:  
A User's Guide  
(Version 7.8010 from April 3, 2025)**

**D. Heck, T. Hüge and T. Pierog**

Institut für Astroteilchenphysik

KIT - Universität des Landes Baden-Württemberg und  
nationales Forschungszentrum in der Helmholtz-Gemeinschaft

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## Abstract

### **Extensive Air Shower Simulation with CORSIKA: A User's Guide**

CORSIKA is a detailed simulation program for extensive air showers initiated by high energy cosmic particles. The user's guide explains the installation of the code, all the necessary input data sets, the selection of simulation parameters, and the structure of the program outputs.

## Zusammenfassung

### **Simulation ausgedehnter Luftschauer mit CORSIKA: Eine Benutzeranleitung**

CORSIKA ist ein Programm zur detaillierten Simulation von ausgedehnten Luftschauern, die durch hochenergetische kosmische Strahlung ausgelöst werden. Die vorliegende Anleitung erläutert die Installation des Programms, alle nötigen Eingabedateien, die Wahl der Simulationsparameter und die Struktur der Ausgaben des Programms.



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# 1 Introduction

CORSIKA (**CO**smic **R**ay **SI**mulations for **KA**scade) is a detailed Monte Carlo program to study the evolution and properties of extensive air showers in the atmosphere. It was developed to perform simulations for the KASCADE experiment [1] at Karlsruhe in Germany. This experiment measured the elemental composition of the primary cosmic radiation in the energy range  $3 \times 10^{14}$  to  $1 \times 10^{17}$  eV and after its upgrade to KASCADE-Grande [2] it reached  $10^{18}$  eV. The first version 1.0 of CORSIKA is dated from October 26, 1989. Since that version it has been extended and improved continuously.

The CORSIKA program [3] allows to simulate interactions and decays of nuclei, hadrons, muons, electrons, and photons in the atmosphere up to energies of some  $10^{20}$  eV. It gives type, energy, location, direction and arrival times of all secondary particles that are created in an air shower and pass a selected observation level.

CORSIKA is a complete set of standard FORTRAN routines. It uses no additional program libraries for the simulation of air showers. Therefore, it runs on (almost) every computer where FORTRAN is available.

The CORSIKA program consists basically of 4 parts. The first part is a general program frame handling the in- and output, performing decay of unstable particles, and tracking of the particles taking into account ionization energy loss and deflection by multiple scattering and the Earth's magnetic field. The second part treats the hadronic interactions of nuclei and hadrons with the air nuclei at higher energies. The third part simulates the hadronic interactions at lower energies and the fourth part describes transport and interaction of electrons, positrons, and photons. CORSIKA contains several models for the latter three program parts that may be activated optionally with varying precision of the simulation and consumption of CPU time.

High-energy hadronic interactions may be treated by one of the models: The Dual Parton Model DPMJET [4], the simple Monte Carlo generator HDPM [5] which is inspired by the Dual Parton Model and tries to reproduce relevant kinematical distributions being measured, the outdated quark-gluon-string model QGSJET01 [6], the mini-jet model SIBYLL [7, 8, 9], or the outdated VENUS [11]. As a sixth alternative there is added a link to the outdated NEXUS model [12], which combines algorithms of VENUS and QGSJET with new ideas, based on H1 and Zeus data. The last models included are EPOS LHC-R [13, 14] (based on the NEXUS framework but with important improvement concerning hard interactions and nuclear and high density effect, and taking into account LHC data) and QGSJET-III-01 [16] which has treatment of Pomeron-Pomeron interaction at any order including loop diagrams (and tuned to LHC). EPOS is the only model available here reproducing heavy ion data from SPS, RHIC, and LHC.

The low-energy hadronic interactions are simulated alternatively with one of the codes: FLUKA [17] which is a very refined model with many details of nuclear effects, GHEISHA [18] that is a outdated detector Monte Carlo program in the energy region up to some hundred GeV, or UrQMD [19] which describes microscopically the low energetic hadron-nucleus collisions.

For all models the hadronic interaction cross-sections at higher and lower energies are

adopted according to the used model.

The interactions of electrons and photons can be treated either with the adapted EGS4 code [20] following each particle and its reactions explicitly, or using the analytic NKG formulae [21] to obtain electron densities at selected locations and the total number of electrons at up to 10 observation levels.

Further on it is optionally possible to explicitly generate Cherenkov light in the atmosphere, to handle electronic and muonic neutrinos and anti-neutrinos, and to simulate showers with flat incidence. Moreover the HERWIG [22] interaction routines have been linked [23] with CORSIKA to handle primary neutrinos.

To shorten the computing times for ultra-high energy showers above  $10^{16}$  eV the thin sampling option exists, by which only a fraction of the secondary particles is followed in the shower development. A different approach to reduce the simulation time solves numerically the cascade equations for the different particle species. This approach is best suited to determine the longitudinal shower profile as measured by the fluorescence telescopes of the Pierre Auger observatory [24] and is realized in the program CONEX [25, 26, ?] now linked with CORSIKA. To enable a full Monte Carlo simulation without thinning the bulk of subshowers may be run in parallel on a cluster of many CPU-slaves as recently realized in the PARALLEL option [27].

There exists as well a program version that is not suited for air shower simulation but for testing the hadronic interaction models.

A detailed description of the CORSIKA program frame, the used cross-sections, the hadronic interaction model HDPM, the electromagnetic interaction models, and the particle decays has been published in Ref. [3]. For details of the DPMJET, EPOS, NEXUS, QGSJET, SIBYLL, VENUS, FLUKA, GHEISHA, UrQMD, and EGS4 programs see Refs. [4, 13, 14, 12, 6, 15, 16, 8, 11, 17, 18, 19, 20]. However, minor modifications were made to these codes to adapt them for simulation of extensive air showers. A comparison of the various hadronic interaction models is given in Refs. [29, 30, 31].

Besides the explanation [3] of the physics implemented in CORSIKA, this CORSIKA GUIDE is a supplementary description of the technical handling and running of CORSIKA7.8010. It contains information about the installation of the program, the required input data, file formats, parameter settings, outputs, and other technical details. This CORSIKA GUIDE is an updated version of Ref. [32].

For **citation of CORSIKA in your publications** you might use Ref. [3] which is available from the CORSIKA www-page <https://www.iap.kit.edu/corsika/70.php> .

Please do not forget also to make a **reference to the hadronic interaction models** which you used in your simulations. The correct references you find in the bibliography (page 187) of this CORSIKA GUIDE and (in L<sup>A</sup>T<sub>E</sub>X format) in the file '*references.tex*' (see directory *doc/*).

If you have problems in installing or running the program, suggestions to improve the code concerning physics, computing, or handling, please contact:

Dr. T. Pierog  
Karlsruhe Institute of Technology  
Institut für Astroteilchenphysik  
Hermann-von-Helmholtz-Platz 1  
D-76344 Eggenstein-Leopoldshafen  
e-mail: <tanguy.pierog@kit.edu>  
Fax: +49-(0)721-608-24075  
Tel.: +49-(0)721-608-28134

Dr. T. Hüge  
Karlsruhe Institute of Technology  
Institut für Astroteilchenphysik  
Hermann-von-Helmholtz-Platz 1  
D-76344 Eggenstein-Leopoldshafen  
e-mail: <dtim.huege@kit.edu>  
Fax: +49-(0)721-608-24075

All users of CORSIKA are kindly asked **not to hand over the program to interested new users**, but rather to send their name and address to the above addresses. By these means new users can be provided with news about the latest program version, error corrections, and updates. Thus problems with outdated versions should be avoided.

## 2 Installation

### 2.1 Download

The CORSIKA code and the files belonging to it can be obtained from the web page [web.iap.kit.edu/corsika/download/](http://web.iap.kit.edu/corsika/download/) of the IAP. The access to this server is protected by password. To fetch the program package you use an internet browser going to the URL: <https://web.iap.kit.edu/corsika/download/> giving username: `corsika` and password: (*which you get by e-mail*) and change to the appropriate subdirectory by clicking on:

`corsika-v780`

Finally you click on:

`corsika-78010.tar.gz`

to get the most recent version of `corsika-78010`. Additional information you find on the CORSIKA Downloads web page <https://www.iap.kit.edu/corsika/79.php> .

For faster transmission we have compressed the `corsika-78010.tar` file. Before using this file first you have to decompress it by applying the 'gunzip' procedure.

For FLUKA two versions exist:

- FLUKACERN maintained by CERN people. The files belonging to this FLUKACERN version are collected in an object file library. This library and the necessary data files are distributed by the FLUKA-CERN organization for compiler versions `gfortran7` and `gfortran9`. They may be downloaded from the web page: <https://www.fluka.cern/> .
- FLUKA maintained by the Fluka collaboration. The files belonging to this FLUKA version are collected in an object file library. This library and the necessary data files are distributed by the FLUKA organization for different computers and compiler versions. You find further details on the web page: <https://www.fluka.org/> .

The files belonging to the HERWIG code<sup>2</sup> may be downloaded from the web page if you wanted a more recent version:

<http://www.hep.phy.cam.ac.uk/theory/webber/Herwig/> .

### 2.2 CORSIKA Files

The CORSIKA 7.8010 set is distributed as a gzipped `.tar` file consisting of several files and directories:

- a file *README* giving these short instructions how to proceed,
- an executable shell script file `./coconut` to be used to install CORSIKA,
- a file *AUTHORS* giving the names of authors,
- a file *ChangeLog* giving the CORSIKA history,

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<sup>2</sup>Version *herwig6521* is currently used.



- a file *COPYING* giving copyright instructions,
- a file *INSTALL* giving detailed instructions how to use *./coconut* to install CORSIKA,
- a file *NEWS* giving latest news about CORSIKA,
- a file *cDO\_NOT\_RUN\_CONFIGURE* to remind users not to use *./configure* but *./coconut* instead,
- files *Makefile.am*, *Makefile.in*, *acinclude.m4*, *aclocal.m4*, *configure.ac*, *configure*. These files and the subdirectory **config/** are needed for the installation and should not be changed unless you know exactly what you are doing. All subdirectories have a *Makefile.am* and a *Makefile.in* which are needed by *configure* to create the proper *Makefiles*.
- a subdirectory **baack/** containing different modules developed by D. Baack<sup>3</sup> to customize the simulation for specific experiments.
- a subdirectory **bernlöhr/** containing 'bernlöhr' package<sup>4</sup>. It is a set of **C**-routines for Imaging Atmospheric Cherenkov Telescopes (IACT) and for use of external atmospheric profiles in the ATMEXT option; several atmospheric profiles *atmprofi.dat* are included together with various auxiliary files. This package is provided by K. Bernlöhr<sup>5</sup>.
- a subdirectory **coast/** in which the COAST interface package is included. COAST is the basis for the COASTUSERLIB, COREAS, INCLINED, and ROOTOUT options and other useful tools available in **coast/CorsikaOptions**.
- a subdirectory **conex/** containing the routines and data sets for the CONEX program (version 7.8010) to treat the essential part of the shower development by solving numerically the cascade equations to describe the development of the various particle species.
- a subdirectory **doc/** containing:
  - a file *CORSIKA\_PHYSICS.pdf* containing the physics description of CORSIKA (Report FZKA 6019),
  - a file *CORSIKA\_GUIDE7.8010.pdf* giving this User's Guide,
  - a file *references.tex* giving the actual references to CORSIKA and the hadronic interaction models in L<sup>A</sup>T<sub>E</sub>X format (to be used in your publications),
  - a file *MPI-Runner\_GUIDE.pdf* giving a description how to run CORSIKA on multi-core computers in parallel using Message Passing Interface (MPI),

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<sup>3</sup><https://tu-dortmund.de/>

<sup>4</sup>*bernlöhr-1.69.tar.gz* is automatically extracted if needed.

<sup>5</sup><https://www.mpi-hd.mpg.de/hfm/~bernlöhr/iact-atmo/>

- a file *MPI-Runner\_PHYSICS.pdf* explains how algorithmically and scientifically the parallel CORSIKA runs using Message Pasing Interface (MPI) are organized,
- a file *coreas\_manual.pdf* describing the functionality of the COREAS option,
- a subdirectory *dpmjet/* containing the source files of the DPMJET III model (*build.f*, *phojet.f*, *dpmjet.f* and *pythia-6427\_cors.f*),
- a subdirectory *epos/* containing the source package for the EPOS model<sup>6</sup>,
- a subdirectory *herwig/* for installation of HERWIG<sup>7</sup> and containing a *README* file inside,
- a subdirectory *include/* containing the *config.h.in* file which is needed by *configure* to do *config.h*. The file *config.h* contains all the preprocessor commands needed by CORSIKA.
- a subdirectory *lib/* together with a script file *libtool* is created when compiling CORSIKA. The subdirectory contains all object files and libraries.
- a subdirectory *nexus/* containing the source package of the NEXUS model<sup>8</sup>,
- a subdirectory *pythia/* containing the modified and extended PYTHIA 6.411 package<sup>9</sup> and a *README.corsika* file inside,
- a subdirectory *run/* to be used to run CORSIKA containing:
  - 19 data sets containing the energy dependent cross-sections for nucleon-nucleus processes (*NUCNUCCS*), for electromagnetic interactions (6 sets *EGSDAT6-x.x*), for QGSJET01 (*QGSDAT01* and *SECTNU*), for QGSJET-II (*sectnu-II-04* and *qgsdat-II-04*), for QGSJET-III (*sectnu-III* and *qgsdat-III.bz2*), numerical data to be used by the VENUS routines (*VENUSDAT*), cross-section table for UrQMD (*UrQMD-⟨VER⟩-xs.dat*), and for the DPMJET routines the Glauber tables (*glaubint.glb*, *glaubtar.glb*) and the data tables *dpmCT14LL.pds* and *dpmjpar.dat*
  - 11 input examples *all-inputs\**, *conex-\*-inputs*, and *parallel-inputs* to steer the simulation with any model (with/without THIN, with STACKIN, with EPOS, with NEXUS, with CONEX, with INCLINED, with CERENKOV or with PARALLEL),
  - 3 data sets *atmabs.dat*, *mirreff.dat*, and *quanteff.dat* to take into account the atmospheric absorption, mirror reflectivity, and quantum efficiency of Cherenkov radiation,

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<sup>6</sup>*epos-lhc-r-cors.tar.gz* is automatically extracted if needed.

<sup>7</sup>Version *herwig6521* is already included here.

<sup>8</sup>*nexus3.97\_cors.tar.gz* is automatically extracted if needed.

<sup>9</sup>*pythia.tar.gz* is automatically extracted if needed.

- the *gr3.txt* data set to give the angular dependent parameters of the GRAPE-3 experiment for the RIGIDITY option,
  - the executable binary files of CORSIKA and of utility programs after compilation.
- a subdirectory `src/` containing:
    - the major part of the source code (*corsika.F* and *corsika.h*) which is a FORTRAN code file with some C-preprocessor commands, that contains the code of CORSIKA including the EGS4 and the HDPM model routines. It contains also the interfaces to FLUKA CERN, GHEISHA, UrQMD, DPMJET, EPOS, NEXUS, QGSJET, SIBYLL, and VENUS, but without those hadronic interaction codes. Also the interface for linking with the CONEX routines is contained within (*corsika.F* and *corsika.h*),
    - 7 separate files with pure FORTRAN code of the GHEISHA, QGSJET01d, QGSJET-II-04, QGSJET-III-01, SIBYLL2.3e, SIBYLL\*, and VENUS routines (*gheisha\_2002d.f*, *qgsjet01d.f*, *qgsjet-II-04.f*, *qgsjet-III.f*, *sibyll2.3e.f*, *sibyll2.3e-star.f*, and *venus.f*),
    - 1 separate files with pure FORTRAN code of the FLUKA interface (*fluka\_interface.f*),
    - two separate files with pure C-code resp. FORTRAN-code of the EFIELD option (*elfield.c* resp. *elfield.f*),
    - 3 special C-routines needed for compilation (*timerc.c*, *tobuf.c*, and *trapfpe.c*),
    - a subsubdirectory `parallel/` containing the shell scripts and C-programs to run CORSIKA distributed on many CPUs in parallel. This directory also houses the `userguide-p11-*.txt` files which describe the parallel handling of CORSIKA on various CPU-clusters using these shell scripts and some FORTRAN programs to handle the results of a parallel simulation.
    - a subsubdirectory `preshower2.0-rev/` containing the C-programs for production of  $\gamma$ -induced preshowers before reaching the top of atmosphere (*preshw.c*, *utils.c*, *utils.h*, *veto.c*, and *veto.h*),
    - a subsubdirectory `utils/` containing utility programs:
      - \* a `coast/` subdirectory containing COAST based utilities BEST CHOICE FOR BEGINNERS (ROOT or ASCII interface, file merger)
      - \* a FORTRAN routine *bcreinpcont.f* or C<sup>++</sup>-routine *bcreinpcont.cpp* for automatic creation of successive steering files to run corsika simulations and all corresponding shell script files,
      - \* a FORTRAN routine *cors2input.f* to reestablish the used input from the binary particle output file,

- \* 3 FORTRAN routines to read the binary particle output files (*corsikahisto.f*, *corsikahisto\_mthin.f*, and *corsikahisto\_thin.f*) and produce some plots in ASCII format (human readable).
- \* a C<sup>++</sup>-routine to read the binary particle output file (*corsikaread.cpp*),
- \* 3 FORTRAN routines to read the binary particle output files (*corsikaread.f*, *corsikaread\_history.f*, and *corsikaread\_thin.f*),
- \* a FORTRAN routine *corsplitevts.f* to split a large particle output file containing several CORSIKA events into files containing single events each,
- \* a FORTRAN routine *corsplitevts.f* to split a large particle output file into several files containing single event each,
- \* a script shell *cskreadme.sh* to check files (to be edited),
- \* a C-routine *map2png.c* to visualize the shower plots established with the option PLOTSH2,
- \* a FORTRAN routine *modelprint.f* which displays the used high-energy and low-energy models from the binary particle output file,
- \* a FORTRAN routine *plottracks3c.f* and its include file *work.inc* to visualize shower plots established with the option PLOTSH,
- \* a C<sup>++</sup>-routine *rcorsik2beok.cpp* to read binary CORSIKA file with any option to test for completeness,
- \* a C<sup>++</sup>-routine *readcorsika.cpp* to read binary CORSIKA file with any option,
- \* a FORTRAN routine *readcsk2ascii.f* converting a binary particle output file to a (readable) ASCII file,
- \* a FORTRAN routine *readtimes.f* to sum up the CPU times of several .lst files (useful at parallel runs),
- \* 3 FORTRAN routines *readcsk2ascii.f*, *readcsk2beok.f* and *readcsk2prtcls.f* to convert CORSIKA file into ASCII and a shell script *readcsk2prtcls.sh* to split a CORSIKA file into ASCII files each containing only one type of particle,
- \* two FORTRAN routines *showsimprods.f* and *showsimulist.f* to establish a tabular of several simulations giving characteristic values of each simulation (primary-id, energy, angles of incidence, ... ) from the binary particle output files,
- \* a FORTRAN routine *sumlistnkginfo.f* to sum up the NKG-output of .lst files to get XMAX averaged over the used showers,
- \* a FORTRAN routine *sumlongifiles.f* to sum up the content of several .longi files (also in case of parallel simulations),
- \* a PYTHON program *gdastool* to create some atmospheric profile readable by CORSIKA from GDAS database,
- \* several other shell scripts and data handling programs to treat the CORSIKA output files.

- a subdirectory `urqmd/` containing the source package of the UrQMD model<sup>10</sup>.

## 2.3 Program File Preparation

By unpacking the *corsika-78010.tar* file with the command

```
tar xvf corsika-78010.tar
```

the file structure of CORSIKA will be established. To install CORSIKA you type

```
./coconut
```

(no argument) and answer the questions ... it's done !

The main source file is *corsika.F* with all common blocks in *corsika.h*. It's a FORTRAN source file with some C-preprocessor commands which allows optional compilation of some part of the code. It uses standard *Makefile* and compiler options, and therefore doesn't require any special software to be installed on your machine<sup>11</sup>. Using the shell script *./coconut*, the user can interactively select the specific CORSIKA version for his application and compile it to get directly ready-to-run executable binary files in the `run/` subdirectory. Computer dependent options and *Makefiles* are prepared automatically by the shell script *configure* called by *./coconut*. The *configure* is a standard portable shell script used together with *make* by GNU-packages to be installed, but *configure* should **NOT** be used directly to get a proper installation of CORSIKA. The *./coconut* shell script has been designed to get a "user friendly" machine dependent installation, so that binary files can be compiled in parallel from the same source directory but on different systems<sup>12</sup> (in a large computer farm for instance). All object files and libraries are "hidden" in a `lib/` subdirectory. If you already used a former version of CORSIKA, and you want to use a *corsika\_compilefile.f* as before, you can optionally save this file during the installation process to compile it yourself.

If you extract executable binary files with different compilers which produce object files and libraries which are not compatible with each other (e.g. 'g77' and 'gfortran') then it is recommended to call

```
./coconut -d
```

between the extractions to erase all established libraries (see Sect. 11 page 154).

The *./coconut* script checks for all options if they can be used on your computer, so comments appearing during the installation should be read carefully.

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<sup>10</sup>*urqmd1.3\_cors.tar.gz* is automatically extracted if needed.

<sup>11</sup>A shell terminal, a C-compiler and a FORTRAN compiler are included in any UNIX based system.

<sup>12</sup>In principle, any system could be used, but only LINUX, Mac OSX, Dec OSF, SunOS, AIX have been tested.

### 2.3.1 Options

At present CORSIKA versions may be generated with the following hadronic interaction models with their cross-section (for determining the mean free path between the interactions) using the options:

**DPMJET** selects DPMJET III routines for the simulation of high energy hadron-nucleus and nucleus-nucleus collisions. Also the DPMJET cross-sections are selected. This option cannot be combined with the low-energy model FLUKA.

**EPOS** selects the EPOS LHC-R routines for the simulation of high energy hadron-nucleus and nucleus-nucleus collisions. Also the EPOS cross-sections are selected.

**NEXUS** selects the NEXUS 3.97 routines for the simulation of high energy hadron-nucleus and nucleus-nucleus collisions. Also the NEXUS 3.97 cross-sections are selected. Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

**QGSJET** selects QGSJET01d routines for the simulation of high energy hadron-nucleus and nucleus-nucleus collisions. Also the QGSJET01d cross-sections are selected. Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

**QGSII** selects QGSJET-II-04 routines for the simulation of high-energy hadron-nucleus and nucleus-nucleus collisions. Also the QGSJET-II-04 cross-sections are selected. Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

**QGSIII** selects QGSJET-III-01 routines for the simulation of high-energy hadron-nucleus and nucleus-nucleus collisions. Also the QGSJET-III-01 cross-sections are selected.

**SIBYLL** selects SIBYLL 2.3e routines for the simulation of high-energy hadron-nucleus and nucleus-nucleus collisions. Also the SIBYLL cross-sections are selected.

**VENUS** selects VENUS routines for the simulation of high-energy hadron-nucleus and nucleus-nucleus collisions. Also the VENUS cross-sections are selected. Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

**FLUKACERN** selects the FLUKA 2020 model<sup>13</sup> for the simulation of low-energy hadron-nucleus collisions. Always the appropriate FLUKA cross-sections are used. This option may be combined with all high energy interaction models excluding DPMJET III.

**FLUKA** selects the FLUKA 2024 model<sup>14</sup> for the simulation of low-energy hadron-nucleus collisions. Always the appropriate FLUKA cross-sections are used. This option may be combined with all high energy interaction models excluding DPMJET III.

**GHEISHA** selects the GHEISHA 2002d routines for the simulation of low-energy hadron-nucleus collisions. Always the appropriate GHEISHA cross-sections are used. This option may be combined with all high energy interaction models. Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

**URQMD** selects the UrQMD 1.3\_cors routines for the simulation of low-energy hadron-nucleus collisions. Always the appropriate UrQMD cross-sections are used. This option may be combined with all high energy interaction models.

The **default** setting obtained without specifying any option is EPOS. For low-energy hadronic interactions the URQMD model is taken.

In addition to these options you may select the following preprocessing options:

**ANAHIST** selects code to generate a histogram file *datnnnnnn.lhbook* for a short analysis of essential properties of the particles arriving at ground (as it is usually performed for the showers simulated at Lyon for the Auger experiment). This option needs the THIN option and linking with the CERN library<sup>15</sup> to get the routines of the HBOOK and ZEBRA packages [33].

**ATMEXT** selects code for treatment of the atmosphere according to MODTRAN model for various atmospheres by tabulated values. ATMEXT is recommended with the CERENKOV option for careful treatment of refractive index. This option links with (compiled) *atmo.c*, *fileopen.c*, and *straux.c* routines of the 'bernlshr' package.

**AUGCERLONG** selects code to fill the Cherenkov column in the table giving the longitudinal particle distribution.

**AUGERHIST** selects code to generate a histogram file *datnnnnnn.lhbook*, containing various histograms of different particle types at up to 20 vertical atmospheric depths to

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<sup>13</sup>As there exist two FLUKACERN libraries suited for different CPU/compiler versions (Linux 64-bit for gfortran 7 and Linux64-bit for gfortran 9), you should select the FLUKACERN library fitting with your installation.

<sup>14</sup>As there exist several FLUKA libraries suited for different CPU/compiler versions (Linux 32-bit and Linux 64-bit for gfortran >8.3), you should select the FLUKA library fitting with your installation.

<sup>15</sup>The CERN library is only available for 32bit computers.

follow the development of shower properties, which are of interest for the Auger experiment (Cherenkov and fluorescence photon production). This option needs the THIN + AUGERINFO options and linking with the CERN library<sup>15</sup> to get the routines of the HBOOK and ZEBRA packages [33].

**AUGERHIT** defines already during the simulation whether the particles arriving the observation level will hit an Auger detector station. Only those particles are written to output, thus saving a lot of disk space for the particle output file.

**AUGERINFO** selects code to writes a file *DATnnnnnnn.info* (rsp. *DATnnnnnnnnnn.info*) instead of *DATnnnnnnn.dbase* (rsp. *DATnnnnnnnnnn.dbase*) file.

**CEFFIC** selects code to respect the atmospheric absorption, mirror reflectivity, and photomultiplier quantum efficiency of Cherenkov light. This option is only available in connection with the CERENKOV option.

**CERENKOV** selects code for additional generation of Cherenkov light. It needs the simulation with EGS4.

**CERWLEN** selects Cherenkov code including the wavelength dependent generation of Cherenkov photons and respects their wavelength dependent refraction of the atmosphere. This option is only available in connection with the CERENKOV option.

**CHARM** selects code for the explicite treatment of charmed particles (only in connection with interaction codes which produce them) and the  $\tau$ -lepton rsp. the  $\tau$ -neutrinos (in connection with the NEUTRINO or NUPRIM option). CHARM excludes the TAULEP option.

**COASTUSERLIB** activates the use of a COAST user library (libCOAST.so) which has to be compiled separately and the parent directory of which must be specified in the environment variable \$COAST\_USER\_LIB. For this option *root* needs to be installed on your system.

**COMPACT** selects a compacted output format for the particle file (unit MPATAP) which is different from the standard output, but better suited for simulations of very large numbers of low energy showers, which produce mostly no or only very few particles.

**CONEX** selects code to solve numerically the cascade equations which describe the shower development after the very first interactions. This gives an enormous speed-up of the simulation especially at primary energies above  $10^{17}$  eV. This option works only in connection with the high-energy hadronic interaction models EPOS, QGSJET, QGSII, or SIBYLL in combination with all three low-energy interaction models. The CONEX option is only available in connection with the CURVED + SLANT + THIN + UPWARD



options.

**COREAS** activates the direct calculation of radio emission from extensive air showers on the basis of the "endpoint formalism". The enclosed COAST package is used to provide the interfacing with the radio emission code. Please refer to the separate CoREAS-manual (available after compiling with the COREAS option switched on) for further details.

**CURVED** selects special code to treat showers with large zenith angles  $70^\circ < \theta < 90^\circ$ .

**DYNSTACK** replaces the default storage for all intermediate particles with a customizable storage structure. This option links with the compiled *DYNSTACK* code in the 'baack' directory.

**EFIELD** selects special code to take into account the atmospheric electrical fields on the transport of charged particles.

**EHISTORY** selects additional code to bring (in combination with the MUADDI, NUADDI, or EMADDI keyword) additional information on the muon, neutrino, and/or electron and photon histories to the particle output file.

**IACT** selects code for simulation of **I**maging **A**tmospheric **C**herenkov **T**elescope arrays. This option is only available in connection with the CERENKOV option and links with (compiled) *iact.c* routines of the 'bernlshr' package.

**IACTEXT** selects code for extended interfacing with the *iact.c* routines of the 'bernlshr' package. This option is only available in connection with the CERENKOV + IACT option.

**ICECUBE1** suppresses uninteresting (low energy) portions of a shower and looks only for such portions which are able to produce neutrinos interesting for the IceCube experiment.

**ICECUBE2** enables options to write the output to a pipe buffer (to be used on-line by a detector simulation program) and/or to write the output as a gzip compressed output file.

**INCLINED** enables the output of ground particles on an inclined plane, which can be user-defined by the INCLIN keyword. The `InclinedPlane` routines used by this option are provided by the enclosed COAST package.

**INTCLONG** selects the longitudinal distribution of Cherenkov photons in longitudinally integrated form. (The default gives the generation of photons/step, so-called differential longitudinal distribution.) INTCLONG excludes the NOCLONG option. This option is only available in connection with the CERENKOV or AUGCERLONG option.

**INTTEST** selects special features for the test of the interaction models (needs routines of the HBOOK and ZEBRA packages [33]). This option is not suited for extensive air shower simulation.

**LPM** selects those EGS4 routines which include the LPM-effect, but without using the THIN option.

**MULTITHIN** selects code which gives the particle's weights for up to 6 different THIN modes of an unthinned shower. It activates automatically the LPM-effect.

**MUONHIST** selects code which generates a histogram file named *datnnnnnnn.lhbook* (rsp. *datnnnnnnnnnn.lhbook*) for the analysis of muon properties of showers. This option needs linking with the CERN library<sup>16</sup> to get the routines of the HBOOK and ZEBRA packages [33].

**MUPROD** selects additional code to bring (in combination with the MUADDI keyword) information to the particle output file on those muons which decay or interact before reaching the observation level. A combination with the EHISTORY option is recommended.

**NEUTRINO** selects code by which neutrinos emerging from pion, kaon, and muon decays are tracked explicitly but without interaction.

**NOCLONG** deselects the longitudinal Cherenkov photon distribution. NOCLONG excludes the INTCLONG option. This option is only available in connection with the CERENKOV or AUGCERLONG option.

**NRREXT** enables extended RUNNR numbers with 9 digits up to 999999999.

**NUPRIM** selects the HERWIG code to treat the first interaction of a primary neutrino.

**PARALLEL** selects code for calculation of a shower distributed on many cores in parallel using a set of shell scripts. As this selection is used at ultra-high energies without thinning, it activates the LPM-effect.

**PARALLELIB** selects code for calculation of a shower distributed on many cores in parallel with MPI (Message Passing Interface). This option is only available in connection with the PARALLEL option.

**PLOTSH** selects code which enables plotting the tracks of the electromagnetic, muonic, and hadronic particles to demonstrate the development of a shower optically.

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<sup>16</sup>The CERN library is only available for 32bit computers.

**PLOTSH2** selects code which enables plotting the tracks of the electromagnetic, muonic, and hadronic particles to demonstrate the development of a shower optically, avoiding the large output files of PLOTSH.

**PRESHOWER** selects routines for coupling with a C-program package to describe the interaction of primary gammas with the Earth's magnetic field at EeV energies and to treat the resulting swarm of em-particles falling onto the top of atmosphere as one shower.

**REMOTECONTROL** selects routines to activate a network layer to send and receive messages over ethernet. This allows to control the simulation and to monitor the simulation progress. This option links with the compiled *remote\_control* code in the 'baack' directory.

**RIGIDITY** rejects low-energy primary particles dependent on the incidence angle and Earth magnetic field which will not produce a shower.

**ROOTOUT** selects routines for writing the particle output in a *root* file e.g. for off-line analysis of the particle output by *root* routines (on the basis of the COAST package). This selection is attainable only by giving '*./coconut -e*' (expert mode of coconut) because it is compatible with *root* version 5 or earlier only.

**SLANT** selects a slant depth scale for the longitudinal distributions instead of the vertical depth scale used otherwise.

**STACKIN** selects code to read in the parameters (type, energy, momenta) of secondary particles resulting from the interaction of an exotic primary (WIMP) which has been treated off-line by a separate program.

**TAULEP** selects code for the explicit treatment of the  $\tau$ -lepton resp. the  $\tau$ -neutrinos (in connection with the NEUTRINO or NUPRIM option). TAULEP excludes the CHARM option.

**THIN** selects the thinning mechanism to reduce the computing time for full simulations, especially for EGS4, and activates the LPM-effect.

**THINMAX** selects a different thinning mechanism. In that case all thinned particles will reach the maximum weight reducing simulation time and artificial fluctuations but at the price of small violation of the energy conservation (both with THIN or MULTITHIN).

**TRAJECT** selects the zenith and azimuth angles in a manner that a (gamma or neutrino) source is followed along its trajectory in the sky.

**UPWARD** selects code to treat upward going particles. This option treats particles in the upward direction (zenith angle  $\geq 90^\circ$ ) with the usual lowest observation level below the starting point of the shower.

**UPWARDOLD** selects previous version of the code to treat upward going particles. In this option, the observation level which stops the particle propagation is above the starting point of the shower.

**VIEWCONE** selects the primary direction to come from a cone around a fixed zenith and azimuth angle.

**VOLUMECORR** selects the angular dependence of the zenith angle distribution as it is needed for a vertical long string detector e.g. AMANDA or ICECUBE (see page 72). (The default takes the zenith angle distribution as observed by a horizontal flat detector.)

**VOLUMEDET** selects the angular dependence of the zenith angle distribution as it will be observed by a volume detector e.g. Cherenkov telescope (see page 72). (The default takes the zenith angle distribution as observed by a horizontal flat detector.)

The major part of the program is machine independent due to the restriction to FORTRAN standards. Nevertheless, there are a few points where computer specific adaptations were necessary. There are prepared options for PCs and work stations running under UNIX (also LINUX) and for Apple Macintosh computers. The versions are automatically<sup>17</sup> activated for the following options:

**UNIX** selects code for calculation on UNIX systems including derivatives like LINUX (e.g. DEC-station under ULTRIX, ALPHA-station under DEC-UNIX [Tru64])<sup>18</sup>.

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<sup>17</sup>Thanks to *configure*, `./coconut` checks your machine configuration and sets the proper options.

<sup>18</sup>Some UNIX machines (HP, IBM RS6000, and installations with GNU g77 compiler [e.g. LINUX]) need the record length parameter RECL (used in the OPEN statement for the external stack MEXST) in bytes instead of 4-byte words. This holds also for DEC-UNIX running the f77 compiler with the option '-assume bytrecl'. For easy adaptation the additional preprocessor-option **BYTERECL** is available to enlarge the RECL-parameter by a factor 4.

Most UNIX systems offer a system routine *date\_and\_time* which overcomes the millenary border. In systems without this internal routine the preprocessor-option **OLDDATE** selects an older routine for date and time. The similar option **OLDDATE2** selects an alternative needed for PGF77 environment on LINUX platforms (IN2P3 Computing Center at Lyon) or HP-UX machines (with the +E1 option of the fort77 compiler). This selection is attainable only by giving '`./coconut -e`' (expert mode of coconut).

The preprocessor-option **IBMRISC** selects the routines *date* and *clock* available on IBM RS6000 machines.

If necessary you might select the option **TIMERC** or you should adapt subroutine *prtime* to call the routines of your system for date and time.

The options **BYTERECL**, **OLDDATE**, **OLDDATE2**, **IBMRISC**, and **TIMERC** are only available in connection with the UNIX option and can not be selected if it's not adapted to your system. They are attainable by calling '`./coconut -e`' (expert mode of coconut).

MAC selects code for calculation on a former 68k-APPLE Macintosh.

Within the *corsika.F* file the optional code for the various versions (to be unpacked by the preprocessor) is marked by

```
#if __flag__
    ...
#else
    ...
#endif
```

blocks. If you are using a sequential computer other than the ones listed here you should try to adapt the UNIX (or perhaps MAC) version to your machine as this is most straight forward to understand.

In the `src/` subdirectory, the *corsika.F* file contains some explanations, the main program with more than 330 subroutines and functions, the EGS4 routines, the NKG routines, the HDPM model routines, the interface routines to FLUKA, GHEISHA, UrQMD, DPMJET, EPOS NEXUS, QGSJET, SIBYLL, and VENUS, the interface routines to CONEX, the special routines for Cherenkov light generation, the routines for generation of Auger-oriented histograms, and the interaction test routines. It is about 84600 lines long. The *corsika.h* file (length about 5300 lines) contains the general common blocks with their explanations. These common blocks are included into the source file during preprocessing. The C-file *timerc.c* contains a 'date and time' routine *timerc* for those UNIX or LINUX systems where the more modern *date\_and\_time* system routine is not available. The C-file *tobuf.c* contains a routine to write out the particle data file independent of the structure (32 bit resp. 64 bit) of the used host. The subsubdirectory `src/preshower2.0-rev/` contains the C-routines *preshw.c*, *utils.c*, *utils.h*, *veto.c*, and *veto.h*.

The *gheisha\_2002d.f*, *qgsjet-II-04.f*, *qgsjet-III.f*, *qgsjet01d.f*, *sibyll2.3e.f*, *sibyll2.3e - star.f*, and *venus.f* routines are about 18200, 17700, 19000, 7500, 24000, 25300 and 18000 lines long, respectively.

The DPMJET III code consists of 4 portions *build.f*, *dpmjet.f*, *phojet.f*, and *pythia-6427\_cors.f* with about 500, 43900, 47200, and 89800 lines each in the `dpmjet/` subdirectory.

The 30 UrQMD 1.3\_cors program files with the 12 include files are collected within the `urqmd1.3_cors.tar.gz`<sup>19</sup> file in the `urqmd/` subdirectory.

The 24 EPOS program files with the 6 include files and the 11 data files are collected within the `epos.lhc-r_cors.tar.gz`<sup>19</sup> file in the `epos/` subdirectory.

The 21 NEXUS 3.97 program files with the 5 include files and the 4 data files are collected within the `nexus3.97_cors.tar.gz`<sup>19</sup> file in the `nexus/` subdirectory.

The CONEX program and data files are collected within the `conex-7.8010.tar.gz`<sup>19</sup> file in the `conex/` subdirectory.

---

<sup>19</sup>Automatically unpacked if selected.

The load modules of the FLUKACERN 2020 library routines and the needed FLUKACERN commons and data files may be downloaded (after being authorized) from the official FLUKACERN web page <https://www.fluka.cern/> and properly installed on your machine. Using the `./coconut` script you are asked for "*enter 'fluka' full path*" and you should enter the full path to the `fluka-4-1.1` files only if the `FLUKA bin` directory is not included in your `\$PATH` environment variable and the needed `\FLUKADATA` environment variable is not defined.

The load modules of the FLUKA 2024 library routines and the needed FLUKA commons and data files may be downloaded (after being authorized) from the official FLUKA web page <https://www.fluka.org/> and properly installed on your machine by setting the environment variable `FLUPRO`<sup>20</sup>. For further information on the FLUKA package you may consult the FLUKA web page or contact A. Ferrari<sup>21</sup>.

The HERWIG routines<sup>22</sup> needed for the NUPRIM version may be updated from <http://www.hep.phy.cam.ac.uk/theory/webber/Herwig/> and adapted according to the README file in the `herwig/` subdirectory (see Sect. A.1 page 162).

### 2.3.2 Example of Program Installation

As an example, if you have the file `corsika-78010.tar.gz` in the current directory, the installation on a LINUX system will look like (commands you are giving are underlined):

```
> tar -zxf corsika-78010.tar.gz
> cd corsika-78010/
> ./coconut
```

```
|=====
|
|           Welcome to COCONUT (v3.1)
|           -- the CORSIKA CONfiguration UTility --
|=====
|
|           create an executable of a specific CORSIKA version
|
|
|           Please read the documentation for a detailed description
|           of the options and how to use it.
|
|           Try './coconut -h' to get some help about COCONUT
|           Use './coconut --expert' to enable additional configuration steps.
|
| (press 'Enter' to select an option followed by "[DEFAULT]" or "[CACHED]")
```

<sup>20</sup>Assuming `csh` shell one uses: `setenv FLUPRO flukadirectory`. For `bash/sh` shells you give: `export FLUPRO=flukadirectory`.

<sup>21</sup><alfredo.ferrari@cern.ch> or <alfredo.ferrari@kit.edu>.

<sup>22</sup>Version `herwig6521` is currently used.

|  
|=====

\*\*\*\*\*

\* WARNING : File compilefile.f found in src/ !

\*

\* This temporary file is renamed "compilefile.f.bak"

\* to prevent conflict with the following installation...

\*\*\*\*\*

\*\*\*\*\*

\* INFO:

\* You are using the cached configuration from "include/config.h".

\* To turn off this you may use the --no-cache option.

\*\*\*\*\*

-----

Compile in 32 or 64bit mode ?

1 - Force 32bit mode

2 - Use compiler default ('-m64' on a 64bit machine) [CACHED]

r - restart (reset all options)

x - exit make

(only one choice possible):

SELECTED : NOM32

-----

Which high energy hadronic interaction model do you want to use ?

1 - DPMJET III (2017.1) with PHOJET 1.20.0

2 - EPOS LHC-R [CACHED]

3 - NEXUS 3.97

4 - QGSJET 01D (enlarged commons)

5 - QGSJETII-04

6 - QGSJETIII-01

7 - SIBYLL 2.3e

8 - VENUS 4.12

r - restart make

x - exit make

(only one choice possible): 4

SELECTED : QGSJET01

Answer the questions to select the options you want to use in CORSIKA (see Sect. 2.3.1 page 10 and Sect. 3 page 31 for details). Press "Enter" to select the options marked by "[CACHED]". After each question, you can choose to restart the installation at the first question ("r") or to stop now ("x").

If an option needs some other files, the installation program will test if they exist<sup>23</sup>. If a problem occurs<sup>24</sup>, the installation program will suggest the solution. Just follow what is written.

Using the default options on a LINUX system, you will get:

```
...
-----
Your final selection to build CORSIKA is:

options:  HORIZONTAL QGSJET01 GHEISHA TIMENEW

-----
Configuration is finished.  How do you want to proceed ?
f - Compiling and remove temporary files [DEFAULT]
k - Compile and keep extracted CORSIKA source code
n - Just extract source code.  Do not compile!

r - restart
x - exit make

(only one choice possible):
```

At this point, you should press "Enter" (rsp. "f"). If you want to see the code used to compile your CORSIKA program, you can answer "k", but this is not needed by the automatic compilation. Finally the installation program creates proper Makefiles and then runs `make install` to compile your CORSIKA program. Select "n" only if you want to modify the source code before compiling or if you want to use your former installation tools, you prepared for an older version of CORSIKA using a "compilefile.f". This option will preprocess the source code, save it in subdirectory `src/` as "corsika" followed by the version number "78010", an underscore with the selected high energy hadronic interaction model<sup>25</sup> and an underscore with the selected low energy hadronic interaction model<sup>26</sup> (and an optional "\_int" if you select the "INTTEST" option), and create the Makefiles. But it will not run `make`. See Sect. 11 (page 154) for more details.

---

<sup>23</sup>The bernlohr, epos, nexus, and urqmd packages will be automatically unpacked in their subdirectories.

<sup>24</sup>Trying to use the CERN library when not installed for instance.

<sup>25</sup>DPMJET, EPOS, NEXUS, QGSJET, QGSII, QGSIII, SIBYLL, SIBSTAR, or VENUS.

<sup>26</sup>flukacern, fluka, gheisha, or urqmd.



```
SELECTED      :  COMPILE
checking whether to enable maintainer-specific portions of Makefiles... no
checking build system type... x86_64-suse-linux
checking host system type... x86_64-suse-linux
checking for a BSD-compatible install... /usr/bin/install -c
checking whether build environment is sane... yes
checking for gawk... gawk
checking whether make sets $(MAKE)... yes
checking for pgf77... no
checking for ifc... no
checking for ifort... no
checking for g77... g77
checking whether the Fortran 77 compiler works... yes
checking for Fortran 77 compiler default output file name... a.out
checking for suffix of executables...
checking whether we are cross compiling... no
checking for suffix of object files... o
checking whether we are using the GNU Fortran 77 compiler... yes
checking whether g77 accepts -g... yes
checking for cc... cc
...
```

*It may take a while ...*

```
...
checking do not compile binaries, just extract CORSIKA compilefile... (cached) no
checking to keep the CORSIKA compilefile... (cached) no
configure: creating ./config.status
config.status: creating Makefile
config.status: creating bernlohr/Makefile
config.status: creating conex/Makefile
config.status: creating dpmjet/Makefile
config.status: creating epos/Makefile
config.status: creating pythia/Makefile
config.status: creating herwig/Makefile
config.status: creating nexus/Makefile
config.status: creating urqmd/Makefile
config.status: creating src/Makefile
config.status: creating run/Makefile
config.status: creating doc/Makefile
config.status: creating lib/Makefile
config.status: creating include/config.h
config.status: executing depfiles commands
config.status: executing libtool commands
```

-----  
configure: Optional packages not found or deactivated on this system: FLUKA COAST ROOT PYTHIA

Compile CORSIKA in "../../lib/Linux" and copy executable in "../../run"

Making install in src

```
make[1]: Entering directory `../../src'
cc -E corsika.F -DHAVE_CONFIG_H -I../include -traditional-cpp -o corsikacompfile.f
g77 -O0 -g -fbounds-check -c -o corsika-corsikacompfile.o `test -f
'corsikacompfile.f' || echo './'`corsikacompfile.f
if cc -DHAVE_CONFIG_H -I. -I. -I../include -g -MT trapfpe.o -MD -MP -MF ".deps/trapfpe.Tpo"
-c -o trapfpe.o trapfpe.c;
then mv -f ".deps/trapfpe.Tpo" ".deps/trapfpe.Po"; else rm -f ".deps/trapfpe.Tpo"; exit 1; fi
g77 -O0 -g -fbounds-check -c -o corsika-qgsjet01d.o `test -f 'qgsjet01d.f' || echo
'./'`qgsjet01d.f
g77 -O0 -g -fbounds-check -c -o corsika-gheisha_2002d.o `test -f
'gheisha_2002d.f' || echo './'`gheisha_2002d.f
/bin/sh ../libtool --mode=link g77 -g -fbounds-check -o corsika corsika-corsikacompfile.o
trapfpe.o corsika-qgsjet01d.o corsika-gheisha_2002d.o -L../../lib/Linux
mkdir .libs
g77 -g -fbounds-check -o corsika corsika-corsikacompfile.o trapfpe.o corsika-qgsjet01d.o
corsika-gheisha_2002d.o -L../../lib/Linux
g77 -g -fbounds-check -c -o plottracks3c.o plottracks3c.f
/bin/sh ../libtool --mode=link g77 -g -fbounds-check -o plottracks plottracks3c.o
g77 -g -fbounds-check -o plottracks plottracks3c.o
make[2]: Entering directory `../../src'
test -z "../../run" || mkdir -p -- . "../../run"
/bin/sh ../libtool --mode=install /usr/bin/install -c `corsika` `../../run/corsika`
/usr/bin/install -c corsika ../../run/corsika
/bin/sh ../libtool --mode=install /usr/bin/install -c `plottracks` `../../run/plottracks`
/usr/bin/install -c plottracks ../../run/plottracks
make[2]: Nothing to be done for `install-data-am'.
make[2]: Leaving directory `../../src'
make[1]: Leaving directory `../../src'
Making install in .
make[1]: Entering directory `../../..'
make[2]: Entering directory `../../..'
make install-exec-hook
make[3]: Entering directory `../../..'
--> "corsika78010Linux.QGSJET_gheisha" successfully installed in :
run/
```

```
--> You can run CORSIKA in run/ using for instance :  
./corsika78010Linux_QGSJET_gheisha < all-inputs > output.lst
```

```
make[3]: Leaving directory '././././.'  
make[2]: Nothing to be done for 'install-data-am'.  
make[2]: Leaving directory './././.'  
make[1]: Leaving directory './././.'
```

Now the installation is finished. As written, an executable binary file is copied into the `run/` subdirectory where all data files are placed. The name of this file is composed by "corsika" followed by the version number "78010", the system name as given by the `uname` UNIX command, an underscore with the selected high energy hadronic interaction model<sup>27</sup> and an underscore with the selected low energy hadronic interaction model<sup>28</sup> (and an optional "\_int" if you select the "INTTEST" option or an optional "\_conex" if you selected the "CONEX" option). As a consequence, you can select different model combinations on different systems without any conflict<sup>29</sup>.

## 2.4 Data Files

### 2.4.1 Input Files

To run a simulation one needs to read several input files. These are:

- The Glauber tables to derive nucleon-nucleus and nucleus-nucleus cross-sections from hadron-nucleon cross-sections are listed in file *NUCNUCCS* which is 2873 lines long. They are read via logical unit NUCNUC (NUCLeus-NUCLeus interactions, by default 11).
- The cross-sections and branching ratios for the EGS4 routines are contained in the data files *EGSDAT6\_x.x* with a length of 3021 lines each. These files differ in the lowest kinetic energy to be followed within a range between 3 MeV ( $x.x = 3.$ ) and 50 keV ( $x.x = .05$ ). Only one of the sets is selected by the program and connected to the logical unit KMPI (by default 12) in a manner appropriate to the lowest energy of em-particles specified by the user. As the muon nuclear interactions use the routines to treat photonuclear interactions, always an *EGSDAT6\_x.x* file must be read in.
- The DPMJET routine *DPMJIN* in *corsika.F* needs the Glauber tables named *glaubtar.glb* (rsp. *glaubint.glb* for the interaction test). The Glauber tables are read in via unit 9. Additionally some parameterized data contained within the data files *dpmCT14LL.pds* and *dpmjpar.dat* are read in by the coupled PHOJET routines via logical units 1 rsp. 12.

---

<sup>27</sup>HDPM, DPMJET, EPOS, NEXUS, QGSJET, QGSII, QGSII, SIBYLL, SIBSTAR, or VENUS.

<sup>28</sup>fluka, gheisha, or urqmd.

<sup>29</sup>In case of conflict, using different options but with the same models for instance, the program will ask you if you want to rename the binary file.

- The EPOS routines get user-specified parameters from a scratch file via logical unit EPOPRM (by default 97) and need some parameterized data contained within the data files *epos.inics*, *epos.iniev*, *epos.inirj* and *epos.initl*. If the latter files are not existent or do not fit with the user-specified parameters, they are established in a time consuming procedure (some 100 h on a DEC 3000/600 AXP with 175 MHz). Additionally 3 data files containing particle properties and decay channels are read by EPOS, namely *idresi.dt*, *idt.dt* and *idky5.dt*.
- The NEXUS routines get user-specified parameters from a scratch file via logical unit NEXPRM (by default 97) and need some parameterized data contained within the data files *nexus.inics*, *nexus.iniev*, *nexus.inirj*, and *nexus.initl*. If the latter files are not existent or do not fit with the user-specified parameters, they are established in a time consuming procedure (some 100 h on a DEC 3000/600 AXP with 175 MHz).
- The more recent QGSJET-III-01 routines need some parameterized data contained within the data files *qgsdat-III* and *sectnu-III*. The old QGSJET-II-04 routines need some parameterized data contained within the data files *qgsdat-II-04* and *sectnu-II-04*. (The QGSJET01d routines need some parameterized data contained within the data files *QGSDAT01* and *SECTNU*). These data files are written and read via logical units 1 and 2. **Attention: The *qgsdat-III* file has a size of  $\approx 1.4$  GB once uncompressed.**
- The STACKIN option needs the parameters (type, energy, momenta) of the secondary particles coming from the interaction of an exotic primary in a separate file. Its file name is specified by the keyword INFILE (page 80). The first line of the file contains (format free after a leading blank) the number of secondaries and the primary energy (GeV). The following lines contain the current particle number, the particle type, the energy (GeV), the longitudinal momentum, and the two transverse momenta (GeV/c) in the format (2I5,4(1X,E15.7)). The momenta are relative to the direction of the (exotic) primary defined by THETAP (page 77) and PHIP (page 77).
- The VENUS option reads the file *VENUSDAT* which contains parameterized structure function integrals and is 2051 lines long. The file is read via logical unit 14.
- The FLUKACERN option needs various data files from the FLUKACERN library, so you should set an environment variable pointing to the FLUKACERN library<sup>30</sup>
- The FLUKA option needs various data files from the FLUKA library, so you should set an environment variable pointing to the FLUKA library<sup>31</sup>.
- The URQMD option looks for the existence of the *tables.dat* decay width file. If this file exists, it is read in, otherwise the decay width tables are calculated and this file is created.

---

<sup>30</sup>Assuming *bash/sh* shell one uses: `export FLUKADATA=$(fluka-config --data) or export FLUKADATA=/...path.where.fluka.is.installed/data .`

<sup>31</sup>Assuming *csh* shell one uses: `setenv FLUPRO flukadirectory`. For *bash/sh* shells you give: `export FLUPRO=flukadirectory`.

- The ATMEXT option needs tabulated atmosphere input data of the MODTRAN model contained in the *atmprofi.dat* files which are read in by the *atmo.c* routines of the 'bernlrohr' package. Details are given in the comments at the beginning of *atmo.c* and in the documentation supplied with the 'bernlrohr' package.
- The CEFFIC option reads the atmospheric absorption table *atmabs.dat*, photomultiplier quantum efficiency table *quanteff.dat* and/or mirror reflectivity table *mirreff.dat* via logical units MCERABS (by default 20), MCERQEF (by default 21), and MCERMIR (by default 22) respectively.
- The CONEX routine *ConexInit* reads the steering commands for the CONEX calculations from the *conex.HIGHMODEL.lowmodel.param* file via the logical unit INLUN (by default 38). This file is copied from the appropriate *conex.HIGHMODEL.lowmodel.param.in* file contained within the *conex/* subdirectory. These steering commands give the files with the parameters for the calculations of the cascade equations.
- The RIGIDITY option reads the *gr3.txt* data file which contains the zenith and azimuth dependent rejection values (GRAPES-3 experiment).
- Besides these data files CORSIKA needs the input of steering keywords to select the subject and the parameters of the simulation. They have to be supplied by the user. They are read via logical unit MONIIN (MONItor INput, by default 5). The format of the steering keywords and their effect is described in detail in Sect. 4 (page 74 ff.). Examples are given on page 125 and in the file 'inputs'.

All these files are placed in the *run/* subdirectory, where the program can be executed. Using the DATDIR keyword (see page 109) the data input files may be placed in an arbitrary directory (with exception of FLUKA data files and the steering input file with the keywords).

Logical units for in- and output			
Unit name	default	I/O	File name and file
MONIIN	5	I	<i>input</i> , steering keywords
MONIOU	6	O	simulation control output on line printer
MDEBUG	6	O	debug output if DEBUG it selected
NUCNUC	11	I	<i>NUCNUCCS</i> , nucleus-nucleus cross-sections
MPATAP	90	O	<i>DATnnnnnn</i> , particle output and simulation results
MEXST	96	I/O	external particle stack, scratch file
	1	I	<i>dpmCT14LL.pds</i> , data table for DPMJET
	1	I	<i>elasct.bin</i> , table for FLUKA
	1	I	<i>sigmapl.bin</i> , table for FLUKA
	1	I/O	<i>epos.inics</i> , various tables for EPOS
	1	I/O	<i>epos.iniev</i> , various tables for EPOS
	1	I/O	<i>epos.inirj</i> , various tables for EPOS
	1	I/O	<i>epos.initl</i> , various tables for EPOS
	1	I/O	<i>idt.dt</i> , various tables for EPOS
	1	I/O	<i>iresi.dt</i> , various tables for EPOS
	1	I/O	<i>idky5.dt</i> , various tables for EPOS
	1	I/O	<i>nexus.inics</i> , various tables for NEXUS
	1	I/O	<i>nexus.iniev</i> , various tables for NEXUS
	1	I/O	<i>nexus.inirj</i> , various tables for NEXUS
	1	I/O	<i>nexus.initl</i> , various tables for NEXUS
	1	I/O	<i>QGSDAT01</i> , table for QGSDAT01d
	1	I/O	<i>qgsdat-II-04</i> , table for QGSJET-II-04
	1	I/O	<i>qgsdat-III</i> , table for QGSJET-III-01
	2	I/O	<i>SECTNU</i> , nucleus-nucleus cross-section for QGSDAT01d
	2	I/O	<i>sectnu-II-04</i> , nucleus-nucleus cross-section for QGSJET-II-04
	2	I/O	<i>sectnu-III</i> , nucleus-nucleus cross-section for QGSJET-III-01
LDAT	9	I	<i>glaubtar.glb</i> (rsp. <i>glaubint.glb</i> ), Glauber tables for DPMJET
LUNOUT	11	O	<i>fluka11.out</i> control output of FLUKA
KMPI	12	I	<i>EGSDAT6_x.x</i> , EGS4 cross-sections
	12	I	<i>dpmjpar.dat</i> , data tables for DPMJET
	14	I	<i>VENUSDAT</i> , structure function integrals for VENUS
	14	I	<i>nuclear.bin</i> , nuclear data for FLUKA
LUNERR	15	O	<i>fluka15.err</i> error output of FLUKA
MATMFI	19	I	<i>fileatm</i> input file for ATMFILE keyword
MCERABS	20	I	<i>atmabs.dat</i> for atmospheric absorption of Cherenkov light (CEFFIC option)
MCERQEF	21	I	<i>quanteff.dat</i> for photomultiplier quantum efficiency of Cherenkov light (CEFFIC option)
MCERMIR	22	I	<i>mirreff.dat</i> for mirror reflectivity of Cherenkov light (CEFFIC option)
LSTCK	23	I	STACKIN input data file
LSTCK2	24	O	output data file of particles from first interaction
ifcx & ifch	31	O	EPOS or NEXUS check file (not opened)
ifhi	35	O	EPOS or NEXUS histo file (not opened)
INLUN	38	I	<i>conex_HMODEL_1model.param</i> steering parameters for CONEX
MDBASE	45	O	<i>DATnnnnnn.dbase</i> (rsp. <i>DATnnnnnn.info</i> ), run summary file for use in an air shower library
MTABOUT	46	O	<i>DATnnnnnn.tab</i> , table output of $\gamma$ , $e^\pm$ , and $\mu^\pm$
MLONGOUT	48	O	<i>DATnnnnnn.long</i> , output of longitudinal particle numbers and energy deposit

Table 1: Logical units for in- and output with default values and file names (to be continued).

Logical units for in- and output (continued)			
Unit name	default	I/O	File name and file
ifdt	51	O	EPOS or NEXUS data file (not opened)
ifcp	52	O	EPOS or NEXUS copy file (not opened)
LUNPLT	54	O	histogram output file for INTTEST version
	55	O	<i>DATnnnnnnn.&lt;spec&gt;.&lt;proj&gt;.map</i> , output of PLOTSH2
	55	O	<i>DATnnnnnnn.track_em</i> , output of PLOTSH (em comp.)
	56	O	<i>DATnnnnnnn.track_mu</i> , output of PLOTSH (muon comp.)
	57	O	<i>DATnnnnnnn.track_hd</i> , output of PLOTSH (hadron comp.)
	75	I/O	( <i>tables.dat</i> ) decay widths tables for UrQMD
	76	I/O	( <i>UrQMD-&lt;VER&gt;-xs.dat</i> ) total cross-section table for UrQMD
	88	O	histogram output file for ANAHIST/AUGERHIST/MUONHIST vers.
MCETAP	91	O	<i>CERnnnnnnn</i> , Cherenkov photon output
MPACUT	92	I/O	cut parameters in/out for PARALLEL version
MPAINP	93	I	<i>steering</i> file for PARALLEL version
MPAJOB	94	O	parameter output for PARALLEL version
MEXSTJ	95	I/O	stack in/out for PARALLEL version
NEXPRM	97	I/O	EPOS or NEXUS parameters, scratch file
	99	I	parameter table input for RIGIDITY option

Table 1: (continued) Logical units for in- and output with their default values and file names.

## 2.4.2 Output Files

There are several streams of CORSIKA output:

- One is control information about the simulation run itself. This (standard) output<sup>32</sup> comes via the logical unit MONIOU (MONItor OUtput, by default 6). In case of a debugging run very much information is written to the logical unit MDEBUG (Monitor for DEBUGging, by default 6). Further details on this file are given in Sect. 10.1 (page 133).
- The second output stream contains the information about all the particles that reach the observation level. It gets the file name '*DATnnnnnnn*' (rsp. '*DATnnnnnnnnnn*' with option NRREXT, see Sect. 3.5.24 page 61) and is written onto the output directory DSN (keyword DIRECT page 108) via output unit MPATAP (PARTicle TAPE, by default 90) as a 'sequential' 'unformatted' FORTRAN file<sup>33</sup>. *nnnnnnn* (rsp. *nnnnnnnnnn* with option NRREXT) is the run number specified in the keyword RUNNR (page 74). This output may be suppressed (see keyword DIRECT and

<sup>32</sup>Renaming this output to '*DATnnnnnnn.lst*' and redirecting it to the directory specified by the keyword DIRECT is convenient as by the shell commands 'dir' or 'ls -l' all files belonging to one run are displayed consecutively which facilitates book-keeping.

<sup>33</sup>To read this data set the FORTRAN programs *corsikaread.f*, *corsikaread\_history.f* rsp. *corsikaread\_thin.f* available in the `src/utls/` subsubdirectory may be used.

keyword PAROUT, page 108). Further details on this file are given in Sect. 10.2 (page 136).

- For the INCLINED option with an inclined observation plane a binary data output is written to the file named '*DATnnnnnn.inclined*' (rsp. '*DATnnnnnnnnnn.inclined*' with option NRREXT, see Sect. 3.5.24 page 61). The format is identical to the normal binary output data, besides the fact that all coordinates are given within the inclined plane. The origin and the orientation depend on the parameters defining the new observation plane through the keyword INCLIN explained in Sec. 4.71 (page 104). Further details on the output file format are given in Sect. 10.3 (page 148).
- Optionally (steered by keyword OUTFILE page 80) a file is written out via the logical unit LSTCK2 (by default 24) containing the parameters of the secondary particles emerging from the first interaction of a hadronic primary. This file may be used later in a run with the STACKIN option (see Sect. 3.5.34 page 66) to maintain the first interaction even for different random number seeds.
- Optionally a table of the number of the binned  $\gamma$ 's,  $e^\pm$  and  $\mu^\pm$  particles might be written out to the file '*DATnnnnnn.tab*' (rsp. '*DATnnnnnnnnnn.tab*' with option NRREXT, see Sect. 3.5.24 page 61) onto the output directory DSN (keyword DIRECT page 108) via output unit MTABOUT (TABLE OUTput, by default 46). *nnnnnn* (rsp. *nnnnnnnnnn* with option NRREXT) is the run number specified in the keyword RUNNR (page 74). To be activated by the keyword PAROUT (page 108). Further details on this file are given in Sect. 10.6 (page 151).
- The longitudinal distribution of particle numbers and energy deposits can be written out to the file '*DATnnnnnn.long*' (rsp. '*DATnnnnnnnnnn.long*' with option NRREXT, see Sect. 3.5.24 page 61) onto the output directory DSN (keyword DIRECT page 108) via the unit MLONGOUT (LONGitudinal OUTput, by default 48). *nnnnnn* (rsp. *nnnnnnnnnn* with option NRREXT) is the run number specified in the keyword RUNNR (page 74). This output is activated by the FLONGOUT flag (see keyword LONGI page 101). Further details on this file are given in Sect. 10.5 (page 151).
- Another (optional) output file contains the compressed information of the Cherenkov photons. It gets the file name '*CERnnnnnn*' (rsp. '*CERnnnnnnnnnn*' with option NRREXT, see Sect. 3.5.24 page 61) and is written onto the output directory DSN (keyword DIRECT page 108) via output unit MCETAP (CErenkov TAPe, by default 91). *nnnnnn* (rsp. *nnnnnnnnnn* with option NRREXT) is the run number specified in the keyword RUNNR (page 74).
- To establish a summary file on the contents of an air shower library an optional output file containing a run summary is written to '*DATnnnnnn.dbase*' on the directory DSN (keyword DIRECT page 108) via output unit MDBASE (Data BASE, by default 45), activated by keyword DATBAS (page 118). If the NRREXT option (see



Sect. 3.5.24 page 61) is used the data set gets the name '*DATnnnnnnnnnn.dbase*'. *nnnnnnn* (rsp. *nnnnnnnnnn* with option NRREXT) is the run number specified in the keyword RUNNR (page 74). Using the AUGERINFO option the name of this file will be '*DATnnnnnnn.info*' (rsp. '*DATnnnnnnnnnn.info*' with option NRREXT, see Sect. 3.5.24 page 61) and the formats are adapted to the requirements of the Auger experiment. Further details on this file are given in Sect. 10.7 (page 151).

- The PLOTSH files contain all tracking steps for each particle with start and end point to produce demonstration plots of the development of showers. They are written onto the directory DSN via the units 55 (file '*DATnnnnnnn.track\_em*' for em-particles), 56 (file '*DATnnnnnnn.track\_mu*' for muons), and 57 (file '*DATnnnnnnn.track\_hd*' for hadrons) (see Sect. 3.5.27 page 62). *nnnnnnn* is the run number specified in the keyword RUNNR (page 74).
- The PLOTSH2 map files are, basically, two-dimensional histograms containing the number of tracks in each xy-/xz-/yz-bin. Details are given in Sect. 3.5.28, page 63. The map files are named '*DATnnnnnnn.<spec>.<proj>.map*', where *<spec>* stands for 'em', 'mu', or 'hd', and *<proj>* stands for 'xy', 'xz', or 'yz'. They are written onto the directory DSN via the unit 55. *nnnnnnn* is the run number specified in the keyword RUNNR (page 74).
- The interaction test option INTTEST generates histograms; their data are written to the file with the name defined by keyword HISTDS (page 123) via the output unit LUNPLT (Logical UNit for PLoT, by default 52).
- The ANAHIST, AUGERHIST, and MUONHIST options generate histograms; their data are written to the file named '*datnnnnnnn.lhbook*' via the output logical unit 88. *nnnnnnn* is the run number specified in the keyword RUNNR (page 74). Further details on this file are given in Sect. 10.4 (page 150).

During the calculation the program uses a temporary data set (scratch file) as an external particle stack if the internal one is over-full. This data set is connected to unit MEXST (EXternal STack, by default 96).

The names of EPOS data files *epos.inics*, *epos.iniev*, *epos.inirj*, and *epos.initl* may be changed using the keyword EOPAR (page 86). Compared to previous EPOS version, a new EOPAR line is needed to give the path to the *idt.dt*, *idresi.dt* and *idky5.dt* files. The parameters given in the keyword EOPAR are written to a scratch file connected to unit NEXPRM (by default 97). This file is read by the EPOS routines when initializing EPOS.

Similarly, the names of NEXUS data files:

*nexus.inics*, *nexus.iniev*, *nexus.inirj*, and *nexus.initl* may be changed using the keyword NEXPAR (page 87). The parameters given in the keyword NEXPAR are written to a scratch file connected to unit NEXPRM (**neXUS PaRaMeters**, by default 97). This file is read by the NEXUS routines when initializing NEXUS.

Some values of the in- and output units may be redefined by changing their values in the corresponding BLOCK DATA subprograms. Table 1 lists all units together with their default values and the corresponding file names.

## 3 Program Options

### 3.1 High-Energy Hadronic Interaction Models

#### 3.1.1 DPMJET Option

**DPMJET** [4] (**D**ual **P**arton **M**odel with **J**ETs) is a program developed to describe high-energy hadronic interactions of hadron-nucleus and nucleus-nucleus collisions using the two-component Dual Parton Model with soft chains and multiple mini-jets at each elementary interaction. For CORSIKA the version DPMJET III (2017.1) is available. It is linked automatically with PHOJET 1.20.0 and PYTHIA-6.4.27.

For using DPMJET you first have to select the DPMJET option when extracting the FORTRAN code from the source file. The *Makefile* will automatically compile the sources *build.f*, *dpmjet.f*, *phojet.f*, and *pythia-6427\_cors.f* and then link it with the compiled *gheisha\_2002d.f* code (rsp. UrQMD library) and with your CORSIKA program. In your input file you may supply the keyword (page 85):

```
DPMJET      T      0
```

Setting DPMJET to `.false.` (`F`) the simple HDPM routines are used (see also Sect. 3.1.3 page 34).

The DPMJET option needs about the same CPU-time as the VENUS option (NKG enabled, EGS4 disabled) at primary energies of  $10^{15}$  eV. Technically it is possible to use DPMJET up to the highest energies<sup>34</sup>.

DPMJET activates also the inelastic hadron-nucleus cross-sections at higher energies which are calculated by the subroutine *dpjsig*. Nucleus-nucleus cross-sections are derived from the DPMJET nucleon-nucleon cross-sections using the Glauber tables of CORSIKA [3]. The DPMJET cross-sections are selected automatically when the DPMJET option has been used for extracting the FORTRAN code from the source file.

In your input file you may supply the keyword (page 85):

```
DPJSIG      T
```

Setting DPJSIG to `.false.` (`F`) you will use the default cross-sections of CORSIKA as described in Ref. [3].

It should be emphasized, that the DPMJET option cannot be combined with any of the FLUKA options because of several identical common and subroutine names used within DPMJET and FLUKA.

#### 3.1.2 EPOS Option

**EPOS** [28] (**E**nergy conserving quantum mechanical multi-scattering approach, based on **P**artons, **O**ff-shell remnants and **S**plitting parton ladders) like NEXUS combines fea-

---

<sup>34</sup>At energies above  $\approx 1$  EeV some error messages may occur: *Error type m has occurred after nnnn PYEXEC calls....* These error messages may be ignored as after each error message internally a reject is performed to reestablish the failing interaction.

tures of the former VENUS [11] and QGSJET01 [6] with extensions enabling a safe extrapolation up to higher energies, using the universality hypothesis to treat the high energy interactions [12]. Compared to NEXUS, many technical problems have been solved and the screening effects have been simplified using a more phenomenological approach mainly based on the recent RHIC data. In addition, high density effects have been included. The most actual version is EPOS LHC-R [14], in which most recent LHC data are taken into account to constrain model parameters. EPOS LHC-R is a major update of EPOS LHC. Theoretical improvement based on EPOS 4 and new physics processes have been introduced (pion and Pomeron, exchange, soft (forward) and hard charm production, complete isospin symmetry (broken in EPOS LHC), improved nuclear fragmentation). But most important, for the first time, a global approach for tuning from e+e- to PbPb collisions has been used leading to a very nice, detailed and complete description of all accelerator data (including most recent from NA61, LHC and LHCf) AND to results compatible with UHECR. The number of muons is a bit larger (5% for proton primary) at all energies (with more high energy muons compared to EPOS LHC) and Xmax is larger by about 20g/cm<sup>2</sup> at all energies too.

For using EPOS you first have to select the EPOS option when extracting the FORTRAN code from the source file. The *Makefile* will compile<sup>35</sup> *epos\*.f* in the *libepos.a* library and then link it with the compiled *gheisha\_2002d.f* code (rsp. FLUKA library or UrQMD library) and with your CORSIKA program. In your input file you may supply the keyword (page 85):

```
EPOS      T      0
```

Setting EPOS to *.false.* (F) the simple HDPM routines are used (see also Sect. 3.1.3 below).

The standard parameters for EPOS are set in subroutine *aaset* of the *epos-bas-lhcr.f* file. If in your calling directory the data sets *epos.inics*, *epos.iniev*, *epos.inirj*, and *epos.initl* are not existent or not compatible with the selected parameters, they will be calculated at the first call of subroutine *psaini* of *epos-sem-xxx.f* (which takes some 100 h on a DEC 3000/600 AXP with 175 MHz).

Normally all parameters of EPOS are set by subroutine *aaset* (of the *epos-bas-lhcr.f* file) called from subroutine *nexini*. In special cases it may be necessary to overwrite one or more of these parameters or to rename the data files to identify *epos.inixx* files established for different parameter sets. This is performed using the keyword

```
EOPAR      aaaaaaaaaa
```

in the input file (see footnote to keyword EOPAR, Sect. 4.29 page 86). *aaaaaaaaaaaa* is a command line as described in the EPOS documentation. These *aaaaaaaaaaaa* commands are written onto the file connected with the logical unit NEXPRM (by default 97) and read by subroutine *aread* of the *epos-bas-lhc.f* file.

---

<sup>35</sup>The FORTRAN compiler options described in Sect. A.1 (page 160) are used.

Since EPOS reads particles properties and branching ratios for decay from external files now, the path where the files *idresi.dt*, *idt.dt* and *idky5.dt* can be found, HAS TO be defined using one EPOPAR command such as :

```
EPOPAR      fname pathnx   path-to/corsika-xxx/epos
```

The EPOS LHC-R version internally produces charmed particles. So the CHARM option (Sect. 3.5.7 page 50) is selected by default and these charmed particles are explicitly transported by CORSIKA and decay at the end of the transport step employing the PYTHIA decay routines. Also the charmed D-mesons and the charmed  $\Lambda_c$ ,  $\Xi_c$ , and  $\Omega_c$  baryons (including the anti-particles) may serve as projectiles. Moreover strange baryons and anti-baryons are accepted as projectiles in EPOS LHC-R. If EPOS is run without the CHARM option, the charmed particles are decayed in EPOS and not propagated in CORSIKA.

The EPOS option require significantly more time than the other hadronic interaction models. The good description of accelerator data in principle requires the use of URQMD to do hadronic rescattering of secondary particles which makes it incompatible with the URQMD model as low energy model in CORSIKA and multiply the simulation time by a factor 7. That is why the default version of EPOS in CORSIKA is NOT using this option. It gives almost the same Xmax and number of muons with or without hadronic rescattering because this process change the particles at mid-rapidity and not at large rapidities which are important for air shower measurements. The change are smaller than the intrinsic uncertainty of the model which is about  $\pm 5\text{g/cm}^2$  for Xmax and  $\pm 5\%$  for the number of muons (see below). But the hadronic rescattering should be taken into account when comparing to accelerator data (to tune the model parameters).

For people with large computer resources, the original version of EPOS with hadronic rescattering can be activated using the expert mode of coconut (`./coconut -e`). Then it can only be linked to FLUKA for the low energy model. This selection is NOT compatible with the CONEX option.

In this expert mode, 2 other options are available to study the systematic error related to the uncertainty of the pion-Air and kaon-Air multiplicity which is unconstrained by data at high energy. By default this multiplicity is set such that the ratio between these interactions and the multiplicity of proton-Air interaction is somewhat the same at low and high energy and “in the middle” of the 2 other options:

- **HIGHMU** Parameters are set such that this multiplicity ratio reach above 1. at the highest energies and the  $\eta$  particle yield is set as low as possible according to accelerator data. This gives the largest number of muons at ground ( $\sim + 5\%$  at high energy) we can get being compatible with particle detectors. Because of the higher multiplicity, Xmax is reduced by about  $\sim -5 \text{ g/cm}^2$  at  $10^{19}$  eV (no change at low energy).
- **LOWMU** Same parameters are used for p-Air,  $\pi$ -Air and K-Air multiplicity evolution (“theory option”) and the  $\eta$  particle yield is set as high as possible according to accelerator data. In this scenario the multiplicity ratio reach  $\sim 0.75$  at the highest

energies. This gives the lowest number of muons at ground ( $\sim -5\%$  at high energy) we can get being compatible with particle detectors. Xmax is increased by about  $\sim +5$  g/cm<sup>2</sup> due to the lower multiplicity.

Both options do not change the results at low energy where the multiplicity is constraint by accelerator measurements. These options can NOT be used with the CONEX option. Together with CONEX, the EPOS option allow to use primary mass up to 250 nucleons, but this has to be explicitly selected when extracting the FORTRAN code from the source file because of possible high memory consumption (<2GB).

EPOS activates also the inelastic hadron-nucleus interaction cross-sections at higher energies. They are calculated by the subroutine *nexsig*. The EPOS cross-sections are selected automatically when the EPOS option has been used for extracting the FORTRAN code from the source file. In your input file you may supply the keyword (page 86):

```
EPOSIG      T
```

Setting EPOSIG to .false. (F) you will use the default cross-sections of CORSIKA as described in Ref. [3].

The *all-inputs-epos* in `run/` subdirectory is an example input file to run CORSIKA with EPOS.

### 3.1.3 HDPM Routines

**HDPM** is a set of routines to simulate high-energy hadronic interactions. These routines are fast and adjusted to experimental data where available. Proton-proton interactions simulated with HDPM and other models agree fairly well with each other (see Ref. [29]). Experimental data are however rare for high energy nucleon-nucleus or nucleus-nucleus collisions and here the results start to disagree due to the simpler modeling in HDPM. If one is interested in differences of air showers induced by different nuclei one probably gets more realistic results by the detailed simulation with other models (DPMJET, EPOS, NEXUS, QGSJET, SIBYLL, VENUS) than with the HDPM routines.

As the HDPM routines are default you have nothing to specify when extracting the FORTRAN code from the source file. But the compiled *gheisha\_2002d.f* code (rsp. FLUKA library or UrQMD library) will be linked with your CORSIKA program.

### 3.1.4 NEXUS Option

**neXus** [12] (**NEX**t generation of **U**nified **S**cattering approach) combines features of the former VENUS [11] and QGSJET [6] with extensions enabling a safe extrapolation up to higher energies, using the universality hypothesis to treat the high energy interactions [12]. It handles nucleus-nucleus collisions with an up to date theoretical approach. The most actual version is NEXUS 3.97.

Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

For using NEXUS you first have to select the NEXUS option when extracting the FORTRAN

code from the source file. The *Makefile* will compile<sup>36</sup> *nexus-xxx.f* in the *libnexus.a* library and then link it with the compiled *gheisha\_2002d.f* code (rsp. FLUKA library or UrQMD library) and with your CORSIKA program. In your input file you may supply the keyword (page 87):

```
NEXUS      T      0
```

Setting NEXUS to *.false.* (**F**) the simple HDPM routines are used (see also Sect. 3.1.3 above).

The standard parameters for NEXUS are set in subroutine *aaset* of the *nexus-bas.f* file. If in your calling directory the data sets *nexus.inics*, *nexus.inicsei*, *nexus.inidi*, *nexus.iniev*, *nexus.inigrv*, *nexus.inirj*, *nexus.inirjei*, and *nexus.initl* are not existent or not compatible with the selected parameters, they will be calculated at the first call of subroutine *psaini* of *nexus-sem.f* (which takes some 10 h on a DEC 3000/600 AXP with 175 MHz).

Normally all parameters of NEXUS are set by subroutine *aaset* (of the *nexus-bas.f* file) called from subroutine *nexini*. In special cases it may be necessary to overwrite one or more of these parameters or to rename the data files to identify *nexus.inixx* files established for different parameter sets. This is performed using the keyword

```
NEXPAR      aaaaaaaaaa
```

in the input file (see footnote to keyword NEXPAR, page 87). *aaaaaaaaaa* is a command line as described in the NEXUS documentation. These *aaaaaaaaaa* commands are written onto the file connected with the logical unit NEXPRM (by default 97) and read by subroutine *aread* of the *nexus-bas.f* file.

The NEXUS option needs roughly 7.5 times more CPU-time than the VENUS option ( $E_0 = 10^{15}$  eV, NKG enabled, EGS4 disabled).

NEXUS activates also the inelastic hadron-nucleus interaction cross-sections at higher energies. They are calculated by the subroutine *nexsig*. The NEXUS cross-sections are selected automatically when the NEXUS option has been used for extracting the FORTRAN code from the source file. In your input file you may supply the keyword (page 88):

```
NEXSIG      T
```

Setting NEXSIG to *.false.* (**F**) you will use the default cross-sections of CORSIKA as described in Ref. [3].

The *all-inputs-nexus* in *run/* subdirectory is an example input file to run CORSIKA with NEXUS.

### 3.1.5 QGSJET, QGSII and QGSIII Options

**QGSJET** [6, 15, 16] (**Q**uark **G**luon **S**tring model with **J**ETs) is a program developed to describe high-energy hadronic interactions using the quasi-eikonal Pomeron parameterization for the elastic hadron-nucleon scattering amplitude. The hadronization process

---

<sup>36</sup>The FORTRAN compiler options described in Sect. A.1 (page 160) are used.

is treated in the quark gluon string model. The most actual version is QGSJET-III-01 [16] including higher twist effects and the cross-section is tuned to LHC data.

For using QGSJET you first have to select the **QGSJET** resp. **QGSII** or **QGSIII** options when extracting the FORTRAN code from the source file.

Since QGSJET01d is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

The *Makefile* will link the compiled *qgsjet-III.f* (resp. *qgsjet01d.f* or *qgsjet-II-04.f*) and *gheisha\_2002d.f* codes (resp. FLUKA library or UrQMD library) with your CORSIKA program. The *qgsjet-III.f* package will behave differently compared with the older *qgsjet-II-04.f* or *qgsjet01d.f*.

In your input file you may supply the keyword (page 88):

```
QGSJET      T      0
```

Setting QGSJET to .false. (F) the simple HDPM routines are used (see also Sect. 3.1.3 page 34).

If in your calling directory the data sets *qgsdat-III*<sup>37</sup> and *sectnu-III* (resp. *QGSDAT01* and *SECTNU* or *qgsdat-II-04* and *sectnu-II-04*) are not yet existent they will NOT be calculated because of excessively long time needed.

QGSJET activates also the inelastic hadron-nucleus and nucleus-nucleus interaction cross-sections at higher energies which are supplied in the *sectnu-III* file read in by the *qgsjet-III-01* [16] (resp. *sectnu-II-04* file read in by the *qgsjet-II-04* [15] or the *QGSDAT01 file* read in by the *qgsjet01d* [6]) program package. The *qgsjet-III-01* cross-sections (resp. *qgsjet-II-04* or *QGSJET01d* cross-sections<sup>38</sup>) are selected automatically when the QGSJET option has been used for extracting the FORTRAN code from the source file. In your input file you may supply the keyword (page 88):

```
QGSSIG      T
```

Setting QGSSIG to .false. you will use the default cross-sections of CORSIKA as described in Ref. [3].

### 3.1.6 SIBYLL Option

**SIBYLL** [7, 8, 9] is a program developed to simulate hadronic interactions at extreme high energies based on the QCD mini-jet model. The actual version [9] is SIBYLL 2.3e. For using SIBYLL you first have to select the SIBYLL option when extracting the FORTRAN code from the source file. The *Makefile* will link the (compiled) *sibyll2.3e.f* and *gheisha\_2002d.f* codes (resp. FLUKA library or UrQMD library) with your CORSIKA program. In your input file you may supply the keyword (page 88):

<sup>37</sup>Binary type file is available for LINUX systems on the gitlab server for downloading <https://gitlab.iap.kit.edu/AirShowerPhysics/corsika-data>.

<sup>38</sup>Omitting the default QGSJETOLD selection uses hadron-air cross-sections increased by 3 % to take into account the individual nuclear radii of <sup>14</sup>N and <sup>16</sup>O as stated in Ref. [35].



SIBYLL T 0

Setting FSIBYL to `.false.` (**F**) the simple HDPM routines are used (see also Sect. 3.1.3 page 34).

The SIBYLL version 2.3e internally produces charmed particles. So the CHARM option (Sect. 3.5.7 page 50) is selected by default and these charmed particles are explicitly transported by CORSIKA and decay at the end of the transport step employing the PYTHIA decay routines. Also the charmed D-mesons and the charmed  $\Lambda_c$ ,  $\Xi_c$ , and  $\Omega_c$  baryons (including the anti-particles) may serve as projectiles. Moreover strange baryons and anti-baryons (except the  $\Omega$ ) are accepted as projectiles in SIBYLL 2.3d. If SIBYLL is run without the CHARM option, the charmed particles are decayed in SIBYLL and not propagated in CORSIKA.

With or without the CHARM option, it is possible to switch off completely the charm production in SIBYLL using the keyword (page 89):

SIBCHM F

in the input file. It will suppress the high energy muons coming from charm decay. Note that it may change slightly the results of the model.

SIBYLL activates also the inelastic hadronic interaction cross-sections at higher energies which are supplied with the SIBYLL [8] program package. They are based on QCD calculations, details are given in [8]. SIBYLL also delivers nucleus-nucleus cross-sections. The SIBYLL cross-sections are selected automatically when the SIBYLL option has been used for extracting the FORTRAN code from the source file.

In your input file you may supply the keyword (page 89):

SIBSIG T

Setting SIBSIG to `.false.` (**F**) you will use the default cross-sections of CORSIKA as described in Ref. [3].

For people who wants to have a model producing artificially a lot more muons (up to +40% at the highest energy starting from no increase at low energy and a reduction of Xmax by about -5 to -10 g/cm<sup>2</sup>), the original version of SIBYLL can be replaced by SIBYLL\* (SIBSTAR) [8] using the expert mode of coconut (`./coconut -e`). Various options are then available for SIBYLL\*:

- **IMODO** Same as default Sibyll2.3e.
- **IMOD1** rho-meson enhancement
- **IMOD2** baryon pair enhancement
- **IMOD3** kaon enhancement
- **IMOD4** mixed enhancement (rho and baryon), default

- **IMOD5** rho-mix (rho component of mixed model)
- **IMOD6** baryon-mix (baryon component of mixed model)

This is NOT based on any real physics processes and should **NOT** be used to analyse and interpret data. It is only meant to test detector response for instance.

This selection is compatible with the CONEX option only for the default mode IMOD4.

### 3.1.7 VENUS Option

**VENUS** [11] (**V**ery **E**nergetic **N**Uclear **S**cattering) is a program developed to simulate ultra-relativistic heavy ion collisions. The actual version is VENUS 4.12.

Since this model is deprecated, this selection is attainable only by using the expert mode of coconut (page 154).

For using VENUS you first have to select the VENUS option when extracting the FORTRAN code from the source file. The *Makefile* will link the compiled *venus.f* and *gheisha\_2002d.f* codes (rsp. FLUKA library or UrQMD library) with your CORSIKA program. In your input file you may supply the keyword (page 89):

```
VENUS      T      0
```

Setting VENUS to .false. (F) the simple HDPM routines are used (see also Sect. 3.1.3 page 34).

Normally all parameters for VENUS are supplied by the routine *venini*. In special cases it may be necessary to overwrite one or more of these parameters specified by its name PARCHA and its new value PARVAL. This is performed using the keyword (page 90)

```
VENPAR     PARCHA   PARVAL
```

in the input file (page 90).

The VENUS option needs roughly 15 times more CPU-time than the HDPM option (NKG enabled, EGS4 disabled).

VENUS activates also the inelastic hadron-nucleus interaction cross-sections at higher energies which are calculated by the subroutine *vensig*. Nucleus-nucleus cross-sections are derived from the VENUS nucleon-nucleon cross-sections using the Glauber tables of CORSIKA [3]. The VENUS cross-sections are selected automatically when the VENUS option has been used for extracting the FORTRAN code from the source file. In your input file you may supply the keyword (page 90):

```
VENSIG     T
```

Setting VENSIG to .false. you will use the default cross-sections of CORSIKA as described in Ref. [3].

## 3.2 Low-Energy Hadronic Interaction Models

### 3.2.1 FLUKA Options

**FLUKA** (**FL**Uctuating **KA**scade) [17] is a package of routines to follow energetic particles through matter by the Monte Carlo method. In combination with CORSIKA only that part is used which describes the low-energy hadronic interactions. The FLUKACERN versions provided by CERN may be found on the web page : <https://www.fluka.cern/> . A detailed description of the processes simulated by the FLUKA versions may be found on the FLUKA web page : <https://www.fluka.org/> .

Both FLUKA versions are used within CORSIKA to calculate the inelastic hadron cross-sections with the components of air and to perform their interaction and secondary particle production, including many details of the de-excitation of the target nucleus.

If you have selected the FLUKACERN option, the *Makefile* will link the FLUKACERN library<sup>39</sup> with your CORSIKA program (pages 161 and 163). To run the FLUKACERN version, a mandatory environment variable<sup>40</sup> tells the system where to find the binary data files (page 24) needed by the FLUKACERN routines. If this variable is not defined, CORSIKA will crash during initialization. The version number saved in the the event header block (page 139) for FLUKACERN is -1.

If you have selected the FLUKA option, the *Makefile* will link the FLUKA library<sup>41</sup> with your CORSIKA program (pages 161 and 163). Please verify that the compiler (g77 resp. gfortran in 32 or 64 bit mode) which your system uses for the CORSIKA program fits with that employed in the compilation of the available FLUKA package. To run the FLUKA version, an environment variable<sup>42</sup> tells the system where to find the binary data files (page 24) needed by the FLUKA routines. The actual version number is saved in the the event header block (page 139) as (Major Fluka version number\*100000 + (Minor Fluka version number)\*100 + (Respin Fluka version number).

In case of problems with any of the FLUKA versions you should consult the files *fluka11.out* and *fluka15.err* which might give hints on the reason for a crash (e.g. expiration date of FLUKA exhausted).

It should be emphasized, that none of the FLUKA options can be combined with the DPMJET option because of several identical common and subroutine names used within FLUKA and DPMJET.

### 3.2.2 GHEISHA Option

**GHEISHA** (**G**amma **H**adron **E**lectron Interaction **S**hower code) is an interaction package widely used in the detector Monte Carlo program GEANT [36] that has proven its qualities in describing hadronic collisions up to some 100 GeV in many experiments.

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<sup>39</sup>See footnote page 11 for the two available FLUKACERN libraries.

<sup>40</sup>Assuming *bash/sh* shell one uses: `export FLUKADATA=$(fluka-config --data) or export FLUKADATA=/...path.where.fluka.is.installed/data .`

<sup>41</sup>See footnote page 11 for the different available FLUKA libraries.

<sup>42</sup>Assuming *csh* shell one uses: `setenv FLUPRO flukadirectory`. For *bash/sh* shells you give: `export FLUPRO=flukadirectory`.

Since this model is deprecated, this selection is attainable only by using the expert mode of *coconut* (page 154).

A detailed description of the physics processes covered by GHEISHA may be found in Ref. [18]. The GHEISHA version is taken as distributed in October 17, 1994 with the GEANT package [36] version 3.21/03 by CERN. Recently some errors were eliminated using fixes obtained from SLAC [37] and now all variables are used in double precision. To discriminate against the uncorrected single-precision GHEISHA version it is renamed to *gheisha\_2002d.f*. GHEISHA is used in CORSIKA to calculate the elastic and inelastic cross-sections of hadrons below 80 GeV in air and their interaction and particle production.

The *Makefile* will link the compiled *gheisha\_2002d.f* code with your CORSIKA program, if you have selected the GHEISHA option.

### 3.2.3 URQMD Option

**UrQMD** (Ultra-relativistic Quantum Molecular Dynamics) is an interaction package designed to treat low energy hadron-nucleus interactions. A detailed description of this model may be found in Ref. [19]. UrQMD 1.3.cors is used in CORSIKA to perform the elastic and inelastic interactions of hadrons below 80 GeV in air. The actual linking routines operate with the special UrQMD 1.3.cors version adapted to CORSIKA.

For using UrQMD you first have to select the URQMD option when extracting the FORTRAN code from the source file. For compilation of the *compilefile.f* the UrQMD 1.3.cors include files *boxinc.f*, *colltab.f*, *comres.f*, *coms.f*, *inputs.f*, *newpart.f*, and *options.f* must be available. So the *./coconut* script will install these files if needed before doing a *Makefile* to make a UrQMD library *liburqmd.a* from the UrQMD 1.3.cors FORTRAN files and then link it with your CORSIKA program (pages 161 and 164). In your input file you may supply the keyword (page 120):

```
URQMD    T    0
```

Setting the first parameter FURQMD to *.false.* (F) the program will stop.

### 3.3 Electromagnetic Interactions (NKG/EGS4 Option)

The NKG and EGS4 options are selected by flags of the input file keyword ELMFLG only. A detailed description of the EGS4 program can be found in Ref. [20], and the modifications applied to it are published in Ref. [3].

For using NKG and/or EGS4 you have to activate the flags of the keyword (page 92)

```
ELMFLG   T    T
```

in the input file.

It must be emphasized that at the highest electron and  $\gamma$ -energies above  $10^{17}$  eV the NKG option does not contain the Landau-Pomeranchuk-Migdal effect (which is added to EGS4) which may alter the shower development by the decrease of the pair formation

and bremsstrahlung cross-sections with increasing energy. Therefore the analytical NKG treatment deviates more and more from results gained by the EGS4 option. For example in a  $\gamma$ -induced shower ( $10^{19}$  eV energy, inclined with  $\theta = 60^\circ$ , without PRESHOWER option) the maximum of the electron longitudinal development simulated by EGS4 is reached deeper in the atmosphere by  $\approx 100$  g/cm<sup>2</sup> slant depth than predicted by NKG.

### 3.3.1 NKG Treatment

The first flag activates the analytic NKG treatment of the electromagnetic component. The longitudinal electron numbers and pseudo-age parameters<sup>43</sup> are calculated every 100 g/cm<sup>2</sup> above the lowest observation level and the lateral electron densities are calculated for a radial grid of 80 points at the two lowest observation levels. The inner grid radius is fixed at 100 cm, while the outer radius RADNKG (in cm) is selected by the keyword (page 93):

```
RADNKG    200.E2
```

Also lateral pseudo-age parameters<sup>43</sup> are calculated.

As the NKG formulas do not take into account the curvature of the Earth's surface, for the CURVED option the NKG option is suppressed. As with the COMPACT option the NKG parameters cannot be written out onto the particle file, the NKG flag must be disabled in this case.

### 3.3.2 EGS4 Treatment

The second flag of the keyword ELMFLG activates the full Monte Carlo treatment of the electromagnetic component by the EGS4 package. Both flags may be activated or deactivated independently. No special option for extracting the FORTRAN code from the source file is required. By selecting the CERENKOV option the EGS4 routines are activated automatically.

In most applications (especially Cherenkov radiation from showers induced by primaries with energies in the GeV range) an abbreviated treatment of the multiple scattering of electrons within the EGS4 code is not recommended. If you can afford a lower quality of your simulations but no long computing times, you may specify an enlarged step length factor STEPFC by the keyword (page 93):

```
STEPFC    1.0
```

(See also the comments in Sect. 4.48 page 93.) A detailed discussion on the use of this step length is given in [38].

In the standard version treating pair production and bremsstrahlung, the EGS4 routines do not regard the Landau-Pomeranchuk-Migdal (LPM) effect which should be applied at

---

<sup>43</sup>These pseudo-age parameters should only be used qualitatively. For scientific applications you extract better age parameters from a fit to the lateral distribution of the electrons as simulated by the EGS4 option.

energies above  $E_{lab} > 10^{16}$  eV. The LPM-effect is switched on automatically using the THIN option (see Sect. 4.55 page 96), the PARALLEL option (see Sect. 3.5.26 page 61), the MULTITHIN option (see Sect. 3.5.20 page 59), or the LPM option (see Sect. 3.5.19 page 59).

The files named *EGSDAT6\_x.x* replace the files *EGSDAT5\_x.x*, *EGSDAT4\_x.x*, *EGSDAT3\_x.x*, *EGSDAT2\_x.x*, or *EGSDATA* used in older CORSIKA versions. For the extrapolation to the highest energies the photo-nuclear cross-section is extrapolated according to Cudell et al. [39] published by the Particle Data Group. The low energy threshold of the *EGSDAT6\_x.x* files ranges from 0.05 MeV to 3 MeV. These files differ from the older data sets by the arrangement of the tables containing the  $e^\pm$ -branching ratios and  $\gamma$ -branching ratios, thus giving a more smooth branching ratio for the rare processes of electro-nuclear and photo-nuclear interactions resp. of  $\mu^+\mu^-$  pair formation [40]. A data set with an energy threshold far below ELCUT(4) implies the explicit, but unnecessary production of many bremsstrahlung photons above threshold but below ELCUT(4), resulting in a considerable prolongation of wasted CPU-time. Therefore CORSIKA automatically selects the *EGSDAT6\_x.x* set best suited for the user's specification of the ELCUT(3) and ELCUT(4), thus saving CPU-time.

## 3.4 Cherenkov Options

### 3.4.1 Cherenkov Standard Option

The routines treating the Cherenkov radiation have been supplied by the HEGRA Collaboration [41] and considerably improved by K. Bernlöhner [42]. The Cherenkov light production by electrons, positrons, muons, and charged hadrons is considered in the subroutine `cerenk`. The Cherenkov photons are considered within a wavelength band which may be specified by the lower and upper limits WAVLGL and WAVLGU. Atmospheric absorption of the Cherenkov photons is not taken into account by default, but might be added by the CEFFIC option (see Sect. 3.4.5 page 46). Only Cherenkov photons from downward going particles and arriving at the lowest observation level are recorded.

Charged particles create Cherenkov photons at each tracking step when the condition  $\beta > 1/n$  ( $\beta = v/c$  and  $n =$  refractive index) is fulfilled. The step is subdivided into smaller sub-steps such that the number of Cherenkov photons per sub-step is less than the fixed number CERSIZ, predefined by an input keyword. In such a sub-step all the photons are sent in a compact bunch along a straight line, defined by the emission angle  $\theta_C$  relative to the electron or hadron direction and a random value for the angle  $\phi$  around this direction.

As the major part of the Cherenkov light is produced by electrons it makes no sense to simulate showers with Cherenkov light production unless using the EGS4 option. Therefore the CERENKOV option automatically activates the EGS4 option, too. The CERENKOV option reduces the step length factor STEPCF to 1 by default (page 93).

The generation of Cherenkov photons by charged primary particles is always considered and therefore the internal clock is always started by the entrance into the atmosphere

(and not by the first interaction). By default the deflection in the Earth magnetic field is considered for the charged primary particles, see keyword TSTART (Sect. 4.12 page 78). For higher primary energies it is impossible to write all the photon bunches of one shower to the output file. Therefore, only those bunches are recorded which hit an array at the lowest observation level consisting of  $\text{NCERX} \times \text{NCERY}$  photon detectors arranged with a grid spacing of  $\text{DCERX}$  and  $\text{DCERY}$  cm in  $x$  and  $y$  direction respectively and with  $\text{ACERX} \times \text{ACERY}$  cm<sup>2</sup> area each. Each bunch is represented by 7 words which are the number of Cherenkov photons, the  $x$  and  $y$  position coordinates at the observation level, direction cosines  $u$  and  $v$ , arrival time, and height of production above sea level.

To obtain this program version the CERENKOV option has to be selected when extracting the FORTRAN code from the source file. Via the keyword (page 110)

```
CERARY    27    27   1500.   1500.   100.  100.
```

the geometry of your Cherenkov array may be defined. In the CURVED option even for large arrays it is always assumed that the detector array is positioned in a plane which is not bent by the curvature of the Earth surface (see Sect. 4.72 page 105). A rotation of the Cherenkov array  $x$ -axis relative to (magnetic) North may be respected by the keyword (page 106):

```
ARRANG    0.
```

The bunch size may be selected by the keyword (page 111):

```
CERSIZ    0.
```

The optimal choice of the bunch size depends on the employment of the atmospheric absorption, mirror reflectivity, and photomultiplier quantum efficiency (CEFFIC option, see Sect. 3.4.5 below). Without the CEFFIC option a  $\text{CERSIZ} = 5$  is reasonable, as about one photon of such a bunch survives in an off-line treatment of these effects.

By the keyword (page 111)

```
CERFIL    1
```

the Cherenkov output is directed to the separate Cherenkov output file MCETAP or ( $\text{CERFIL} = 0$ ) to the particle output file MPATAP. In case of a separate output file the Cherenkov output is structured as the particle output file. It contains the event header and the event end block and in between the data blocks. The data structure of the Cherenkov output data set is given in Table 11 (page 141). In the case the Cherenkov bunches are stored in the MPATAP particle output file together with the other shower particles on the same particle output file, a Cherenkov bunch is treated like a particle.

The definition of an array of Cherenkov detectors serves to reduce the required disk space for Cherenkov shower. On the other hand one loses the possibility of using an air shower several times during the analysis with different core locations with respect to the detector. Keeping in mind the excessive computation time for Cherenkov showers a possibility is introduced to use Cherenkov showers multiple times with only a tolerable increase of

storage space. Therefore, already during the simulation it is defined how often a single shower should be used and where in the array the core locations should be. The core locations for each event are chosen with the Sobol quasi-random number generator [43] and are stored in the event header. Correspondingly, the array of Cherenkov detectors is placed several times in the observation plane and store all Cherenkov bunches that hit one of the detectors. This possibility is selected by the keyword (page 113):

CSCAT    ICERML    XSCATT    YSCATT

An event is used ICERML times and the core is scattered in the range  $-XSCATT \leq x_{core} \leq XSCATT$  and  $-YSCATT \leq y_{core} \leq YSCATT$ . Another possibility to define the positions of core locations (without the Sobol quasi-random generator) uses the keyword COREPOS (see Sect. 4.94 page 113). For the analysis of such CORSIKA events the user has to specify the same core locations in the analysis that have been determined during the simulation. The output file size will basically scale with the number of times each event is used, but it is still considerably smaller than the output of the complete Cherenkov component would be.

To obtain this program version the CERENKOV option has to be selected when extracting the FORTRAN code from the source file.

It should be noted, that it is not foreseen to combine the CERENKOV option with the PARALLEL or PARALLELIB options.

### 3.4.2 Cherenkov Wavelength Option

In the CERWLEN option the index of refraction is made wavelength dependent. As a consequence, photon bunches will carry a specific wavelength. Photons of shorter wavelengths (with larger index of refraction) will result in larger Cherenkov cone opening angles and larger bunch sizes. For very fast particles this will generally have a small effect (less than  $0.03^\circ$  in the opening angle, for example) but near the Cherenkov threshold the effect can be larger.

This option may also require to use a smaller maximum bunch size (see keyword CERSIZ page 111) since all photons in a bunch are of the same wavelength and, therefore, the peak quantum efficiency rather than the average quantum efficiency determines the maximum acceptable bunch size. (In combination with the CEFFIC option (see Sect. 3.4.5 page 46) you should use a maximum bunch size of 1, as usual.)

In the THIN option the wavelength of each photon bunch may be included as last item (see Table 11 page 141) replacing the weight parameter. If the output is directed to the output file MCETAP by setting MCERFI  $> 0$  the CERFIL keyword (see Sect. 4.90 page 111) enables to specify the last item as thinning weight or as wavelength.

To obtain this program version the CERWLEN option has to be selected in combination with the CERENKOV option when extracting the FORTRAN code from the source file.



### 3.4.3 Imaging Atmospheric Cherenkov Telescope Option

The routines treating the Cherenkov radiation for **Imaging Atmospheric Cherenkov Telescopes** (IACT option) have been supplied by K. Bernlöhr [42, 44]. The Cherenkov light production by electrons, positrons, muons, and charged hadrons is considered in the subroutine *cerenk*. The positions of the telescopes are defined by the keyword (page 113)

```
TELESCOPE    0.  0.  0.  0.
```

giving the coordinates relative to the center of the lowest observation level (see Sect. 4.95 page 113). So at (0,0,0) means at the core of the shower at the observation level. The data set name for the telescope-specific data output is defined by the keyword (page 114):

```
TELFIL      filename
```

For further details of the IACT option see Ref. [42], the comments at the beginning of the *iact.c* routines, and the documentation supplied with the 'bernlöhr' package.

With the IACT option by default the TMARGIN flag (keyword TSTART) is set to .true. (see Sect. 4.12 page 78), but it may be overridden. This affects the x and y coordinates of the Cherenkov photons arriving at the observation level if the altitude of the first interaction is fixed by the keyword FIXHEI (page 79).

To obtain this program version the IACT option has to be selected together with the CERENKOV option when extracting the FORTRAN code from the source file. The *Makefile* will compile the needed C files of the 'bernlöhr' package in the *libbern.a* library and then link it with your CORSIKA program.

Imaging atmospheric Cherenkov telescopes may be specified **without** the IACT option. In this case the output is directed to the MCETAP output files and a maximum number of 99 telescopes may be specified with the TELESCOPE keyword (see Sect. 4.95 page 113). Now the 5<sup>th</sup> argument ID of this keyword:

```
TELESCOPE    0.  0.  0.  0.  ID
```

specifies the telescope identification number, which is used for the MCETAP output files. For each telescope you will get an individual Cherenkov output file named CERnnnnnnn-teliii (rsp. CERnnnnnnnnnn-teliii with option NRREXT, see Sect. 3.5.24 page 61) where *iii* gives the telescope ID. For MCERFI at least 3 or larger (see CERFIL keyword Sect. 4.90 page 111) the filename convention is changed to DATnnnnnnn.cher-teliii (rsp. DATnnnnnnnnnn.cher-teliii).

### 3.4.4 Imaging Atmospheric Cherenkov Telescope Extension Option

With the IACTEXT option the interface to the *telout* function (*iact.c* routines) is extended by parameters describing the emitting particle. This extended information is stored as an additional photon bunch (after the normal one) with mass, charge, energy, and emission time replacing the *cx*, *cy*, *photons*, and *zem* fields, respectively, and are identified by a

wavelength of 9999. The compact output format is disabled for making that possible. In addition, all particles arriving at the observation level are included in the `eventio` format output file, in a photon-bunch like block identified by array and detector numbers 999.

The `x`, `y`, `cx`, `cy`, and `ctime` fields keep the normal sense, with coordinates, directions and time counted in the detection level reference frame. The particle momentum is filled into the `zcm` field (negative for upward-moving particles) and the particle ID is filled into the `lambda` field. If thinning is used, the particle weight is in the `photons` field.

When compiling `iact.c` manually (instead of taking advantage of the `./coconut` script or the GNU-makefile supplied with the 'bernlshr' package), an additional option `-DIACTEXT` is required to have a consistent interface on both sides.

To obtain this program version the `IACTEXT` option has to be selected together with the `CERENKOV` and `IAC` options when extracting the FORTRAN code from the source file.

### 3.4.5 Cherenkov Light Reduction Option

The standard simulation of Cherenkov photons does not regard light absorption within the atmosphere, telescope mirror reflectivity, or quantum efficiency of the detecting photomultiplier tubes. In the `CEFFIC` option these effects are taken into account at an early stage of the Cherenkov photon simulation, thus the computing time is shortened considerably and the requirements for storage of Cherenkov output are reduced additionally.

Data tables containing information on these three effects as function of photon wavelength are needed in this option depending on the status of the flags specified by the keyword (page 112):

```
CERQEF      F      F      F
```

Standard tables for atmospheric absorption (`atmabs.dat`), mirror reflectivity (`mirreff.dat`) (measured for the re-coated mirrors of the Whipple telescope, Sept. 1993), and quantum efficiency (`quanteff.dat`) (measured for Hamatsu R1398HA photo-multipliers with UV-window and 1.125" tube) are supplied with `CORSIKA`<sup>44</sup>. For applications which are not covered by these standard tables corresponding user supplied tables should be available. The wavelength range, if differing from the standard tables, might be specified by the keyword:

```
CWAVRANG    WLMAX    WLMIN
```

The wavelength spacing is fixed to 5 nm and the number of wavelength bins is derived from  $1 + (WLMAX - WLMIN) / 5.0$  for all three tables. For the `atmabs.dat` table the 51

---

<sup>44</sup>The `atmabs.dat` table is composed of 105 wavelength values between 180 and 700 nm in steps of 5 nm; one line for each wavelength, beginning with the wavelength value [nm] as integer, followed by 51 extinction values, starting at sea level up to 50 km height in steps of 1 km. The data format is (105(I4,5(10F10.3),F10.3)).

For the same 105 wavelengths the `mirreff.dat` and `quanteff.dat` tables contain reflectivity resp. quantum efficiency values written in the format (8F6.3).

extinction values start at sea level and go up to 50 km height in steps of 1 km. The data format for reading is (I4,5(10F10.3),F10.3). For the *mirreff.dat* and *quanteff.dat* tables the format (8F6.3) is used. Further details may be taken from the comments in the employed subroutine `tpdini`.

By early eliminating those Cherenkov photons which are absorbed within the atmosphere, not reflected by the mirror, or not producing photo electrons within the photomultiplier, those suppressed photons are also not counted in the various forms of the longitudinal distributions (see Sect. 3.4.6 below).

As in the CEFFIC option the atmospheric absorption is treated only in a manner suited for planar atmospheres, you should not combine the CEFFIC option with CURVED.

To obtain this program version the CEFFIC option has to be selected together with the CERENKOV option when extracting the FORTRAN code from the source file.

### 3.4.6 INTCLONG and NOCLONG Options

In the Cherenkov version the longitudinal distribution of photons is given in differential mode (i.e. the number of photons generated within each step) by default. By the preprocessor option INTCLONG the integral mode is selected (i.e. accumulated number of generated Cherenkov photons for each step) which needs additional computing time. If both kinds of longitudinal distribution are of no interest, you may deselect the Cherenkov photon distribution completely by the preprocessor option NOCLONG thus saving computing time.

The option INTCLONG is effective also with the AUGCERLONG option (see Sect. 3.5.3 page 48).

To obtain these program versions the INTCLONG resp. NOCLONG option has to be selected together with the CERENKOV resp. AUGCERLONG option when extracting the FORTRAN code from the source file.

## 3.5 Other Non-standard Options

### 3.5.1 ANAHIST Option

The ANAHIST option produces a series of histograms generated with HBOOK routines [33]. The histograms are written into the file named '*datnnnnnn.lhbook*' (page 150) onto the output directory DSN specified by the keyword DIRECT (page 108). To suppress the ordinary particle output file the keyword PAROUT (page 108) might be used. This analysis gives a short overview on various shower properties of the particles arriving at the observation level. The histograms are only established for the lowest observation level. Radial thinning is not applied to the particles sorted into the histograms.

Because of its permanent modifications a comprehensive description of the ANAHIST option is not available.

To obtain this program version the ANAHIST option has to be selected when extracting the FORTRAN code from the source file. The HBOOK routines require linking of the

CERN library<sup>45</sup> with the program; they are not supplied with the CORSIKA package.

### 3.5.2 ATMEXT Option with External Atmospheres

The ATMEXT option allows to use external tabulated atmospheres of the MODTRAN model documented in Ref. [45]. They are provided together with the 'bernlrohr' package [44] as files *atmprofi.dat* and read in with special routines<sup>46</sup> written in C. Further details may be found in Ref. [42], at the beginning of the *atmo.c* routines and in the documentation supplied with the 'bernlrohr' package. See also Sect. 4.22 page 83.

The CURVED option needs the atmospheres in the 5-layer model rather than interpolated from ATMEXT tables. Therefore in the CURVED option AATM, BATM, CATM values are fitted to the tabulated atmospheres and the table interpolation is disabled.

To obtain this program version the ATMEXT option has to be selected when extracting the FORTRAN code from the source file. A linking with the (compiled) *atmo.c* routines of the 'bernlrohr' package is done by the *Makefile*.

### 3.5.3 AUGCERLONG Option

With the AUGCERLONG option it is possible to fill the Cherenkov column in the table of the longitudinal particle distribution without full simulation of the Cherenkov radiation. Because of an elongation of the CPU time the AUGCERLONG option should only be used if the longitudinal Cherenkov distribution is really needed. By the INTCLONG option (page 47) the integrated instead of the differential longitudinal Cherenkov intensity may be selected. The keywords CERSIZ (page 111), CWAVLG (page 111) and CWAVRANG (page 112) are activated with the AUGCERLONG option.

To obtain this program version the AUGCERLONG option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.4 AUGERHIST Option

The AUGERHIST option produces a series of histograms generated with HBOOK routines [33]. The histograms are written into the file named '*datnnnnnn.lhbook*' (page 150) onto the output directory DSN specified by the keyword DIRECT (page 108). As the HBOOK routines use only lower case characters, the DSN directory name should not contain capitals. These histograms show properties of different particle types at up to 20 horizontal levels (to be defined by keyword OBSLEV, page 103). They enable to study the longitudinal development of various shower parameters. At each defined level a series of histograms is generated e.g. for radial distances of different particle species, for the energy deposit by different particle species as function of distance from the shower axis, for energy spectra (as function of distance) and much more. Examples are given in Ref. [46]. For histograms relating to the emission of Cherenkov radiation the bunch

---

<sup>45</sup>The CERN library is only available for 32bit computers.

<sup>46</sup>The reading C-program expects the *atmprofi.dat* files in the directory where you start the CORSIKA executable. Therefore install a symbolic link to the atmosphere files or copy them to the run directory.

size and wavelength band has to be specified using the keywords CERSIZ (page 111) and CWAVLG (page 111).

Only particles arriving at the lowest observation level are directed to the particle output file (page 136) and/or the table output file (keyword PAROUT page 108) and/or binned into the ANAHIST histograms.

Because of its permanent modifications a comprehensive description of the AUGERHIST option is not available.

The AUGERHIST option works only in combination with the THIN option. To obtain this program version the AUGERHIST option has to be selected when extracting the FORTRAN code from the source file. The HBOOK routines require linking of the CERN library<sup>46</sup> with the program; they are not supplied with the CORSIKA package.

### 3.5.5 AUGERHIT Option

With the AUGERHIT option only those particles are written to the particle data output file which hit a detector or its direct neighbourhood thus respecting also particles with inclined arrival directions. This option enormously reduces the disk space required by the particle data output file [47].

The Auger detectors [24] are aligned in a triangular grid with 1500 m distance between the detectors extending in West-East direction and at angles  $\pm 60^\circ$  wrt. the West-East direction. The detector distance may be specified by the keyword AUGSCT (see Sect. 4.74 page 106). To select particles for writing out it is sufficient to examine whether a particle falls on the parallel stripes covering the detector rows in all three directions. The strip half width DRADIUS has to be selected large enough to neglect small deviations of the detectors from the exact grid positions. Also the value of DRADIUS is selected by default large enough to retain particles coming from inclined showers with zenith angles up to  $88^\circ$  which just skim the cylindrical detectors at the upper edge. A more stringent selection and thus a reduced output particle data file size can be achieved by the FTANKSHADOW flag (see Sect. 4.74 page 106) to consider only particles which hit the shadow of an Auger tank. By using the flag FANYMODE only those particles are brought to the output which survive all thinning modes in the MULTITHIN option (Sect. 3.5.20 page 59).

Using the flag FANYMODE of the keyword AUGSCT (Sect. 4.74 page 106) the size of the particle data output file may be reduced respecting only those particles which survive in all thinning modes of the MULTITHIN option (see Sect. 3.5.20 page 59).

To use a shower several times with core locations scattered relative to the detector grid up to 20 core locations can be chosen. This selection is performed with the Sobol quasi-random number generator [43] setting the parameter MAUGPOS of the keyword AUGSCT (see Sect. 4.74 page 106) to a positive value. Setting the parameter MAUGPOS to zero the selection is performed manually with one or several input steering lines with the keyword AUGHIT (see Sect. 4.75 page 107). In both cases the core locations are stored in the event header. Correspondingly, the array of Auger detectors is placed several times in the observation plane, and all those particles which hit one of the detectors are

---

<sup>46</sup>The CERN library is only available for 32bit computers.

stored. Moreover, the detectors around the shower core will reach saturation and therefore particles hitting closer to the shower core than the selectable distance CORECUT (see Sect. 4.62 page 99) can be discarded.

In the AUGERHIT option by default the orientation of the Auger detector array is set by the angle ARRANG (see Sect. 4.73 page 106) which is chosen to  $-92.08^\circ$ . Therefore the positive X-axis is pointing to East. Also the magnetic field components at the Malargue site ( $B_x = 19.52\mu\text{T}$  and  $B_z = -14.17\mu\text{T}$ ) are taken as default values. All other parameters to be used with the AUGERHIT option are read in using the keyword AUGSCT (see Sect. 4.74 page 106).

To obtain this program version the AUGERHIT option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.6 AUGERINFO Option

The AUGERINFO option produces an output file named '*DATnnnnnnn.info*' replacing the '*DATnnnnnnn.dbase*' file, when activated by the DATBAS keyword (page 118). Using the NRREXT option (see Sect. 3.5.24 page 61) the output file gets the name '*DATnnnnnnnnnn.info*'. The output format of the .info file differs from that of the .dbase file to enable the automatic production of a data base for the Auger experiment showing the content of the CORSIKA shower library at the IN2P3 computing center Lyon.

To obtain this program version the AUGERINFO option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.7 CHARM Option

In the standard CORSIKA program the charmed particles are treated implicitly in the hadronic interaction codes. But several codes (at present DPMJET, QGSJET01d, and SIBYLL 2.3d) produce them in a manner that they can be treated with explicitly transporting them respecting their interaction cross-sections calculated by reference [48] and their lifetimes using the CHARM option [49].

With the SIGMAQ keyword (page 91) the cross-sections for the interaction of the charmed (rsp. bottom) particles are determined. The PROPAQ keyword (page 92) determines whether the interactions are handled by routines [48] which extend the PYTHIA 6.411 (rsp. PYTHIA 6.427) package<sup>47</sup>, or in combination with QGSJET01d by the routines of *qgsjet01d.f*<sup>48</sup>. The decay of charmed particles is treated by the PYTHIA package [34]. Therefore in combination with QGSJET01d and SIBYLL 2.3d the modified and extended PYTHIA 6.411 package is installed in the *pythia/* subdirectory for automatic compilation and linking. In combination with DPMJET the implicit PYTHIA 6.427 package with extensions [48] is used and needs no extra linking.

---

<sup>47</sup>The PYTHIA routines *pyspli.f* and *pypdfu.f* are modified and the routines *chabadif.f*, *chabapar.f*, *chamedif.f*, *chamepar.f*, *bobadif.f*, *bobapar.f*, *bomedfif.f*, and *bomepar.f* are added to the PYTHIA 6.411 package.

<sup>48</sup>Before compiling *qgsjet01d.f* the probability of charmed quark rsp. di-quark formation should be enabled by setting the appropriate values of DC(3) and DC(5) in subroutine XXASET.

The CHARM option [49] also enables the treatment of the  $\tau$ -leptons. All possible interactions (bremsstrahlung, pair production, nuclear interactions) of the  $\tau$ -leptons are respected in a manner in analogy with the  $\mu$ -leptons. For the treatment of primary  $\nu_\tau$  and  $\bar{\nu}_\tau$  neutrinos the CHARM option has to be combined with the NUPRIM option (page 61). For tracing back PYTHIA problems in combination with QGSJET01d or SIBYLL 2.3d the PYTHIA keyword is available (page 121).

To obtain this program version the CHARM option has to be selected when extracting the FORTRAN code from the source file.

**Please verify** that in combination with QGSJET01d:

- you have activated the charm production in the *qgsjet01d.f* file<sup>48</sup>;
- you have available the PYTHIA 6.411 package with all modifications and extensions<sup>47</sup>.

### 3.5.8 COASTUSERLIB Option

By the COASTUSERLIB option one gets access to each path element of all particles produced during the simulation process. Therefore this option opens a variety of powerful analysis tools by user-written C++-routines through the enclosed COAST package. The user-defined code can be applied for many purposes, such as the visualization of particle tracks or the generation of histograms.

Before starting *./coconut* the user has to define the environment variable `$COAST_USER_LIB` pointing to an existing user-defined libCOAST.so library. Without the environment variable this option is not visible in the *./coconut* selection menu. By selecting the COASTUSERLIB option CORSIKA will transmit the path elements of all tracked particles to the user-defined code.

For code development based on COAST, a doxygen documentation is available in `coast/Documentation`

After at least one use of *./coconut*, please use the *make* command in this subdirectory. Then you can find all details using

```
firefox coast/Documentation/COASTDocumentation/html/index.html
```

The COASTUSERLIB option<sup>49</sup> cannot be combined with any other option making use of the COAST interface, i.e. COREAS, INCLINED, or ROOTOUT options, nor with histogram producing options ANAHIST, AUGERHIST, or MUONHIST and also the COMPACT option..

To obtain this program version the COASTUSERLIB option has to be selected when extracting the FORTRAN code from the source file, and also *root* must be installed on your system.

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<sup>49</sup>This option compiles only on LINUX and Mac-OSX machines.

### 3.5.9 COMPACT Output Option

The standard output of CORSIKA is not adequate when simulating a large number of showers initiated by primaries of so low energies, that only a small percentage of them produces particles arriving at the detector level. As most Data blocks of the MPATAP file would be filled up with zeros, a large amount of useless information for the Data blocks and the unnecessary overhead of the Event Header and Event End blocks would be written in this case. This is avoided in the COMPACT option, which writes out only the Run Header and the full Event Header for the first event. For subsequent events only shortened Event Headers (the first 12 parameters) are written. Event End blocks are omitted completely, the Run End block is written as usual. The Data blocks have a maximal length of 39 particles, trailing zeros are suppressed. Further details are given in Sect. 10.2.4 (page 144).

As the NKG parameters are not written out in the COMPACT version, the NKG flag (keyword ELMFLG page 92) should be disabled. The COMPACT option cannot be combined with the COASTUSERLIB, COREAS, INCLINED, or ROOTOUT options.

To obtain this program version the COMPACT option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.10 CONEX Option for Cascade Equations

**Warning:** Despite a lot of successful tests, the authors couldn't test all possible combinations of models and observable types. For the moment the precision expected by the use of the CONEX option should be better than 10% compared to usual CORSIKA for a factor of 10 gain in CPU time but it may depend on the initial conditions and observables. As a consequence they decline any responsibility if an analysis is biased by the use of this option. The user should validate his analysis by the use of CORSIKA without CONEX option for particle based analysis. 1D simulations based on CONEX are already used intensively since many years. The original CONEX program [25, 26] treats the development of showers by numerically solving cascade equations (CE) for the different particle species (hadrons, muons, and electromagnetic particles).

In the CONEX-CORSIKA coupling [50, 51] which resembles the procedure of the SENECA code [52] the shower development starts with a Monte Carlo treatment of the primary particle and also the secondary particles are handled by Monte Carlo techniques as long as their energy  $E$  exceeds the preselected value  $E_{thr}$ . The secondaries with energies  $E < E_{thr}$  are binned into the energy-depth tables which define the "source terms" of the CE. From these source terms the further shower evolution is calculated by solving the CE down to the energy  $E_{low}$  which marks the end of the CE treatment. The parameters needed for solving these CE are transmitted from the parameter sets which are selected by the steering file *conex.HIGHMODEL\_lowmodel.param* (see page 25) in dependence of the selected hadronic interaction model. The solution of the CE can be sampled into individual particles which are saved on the CORSIKA particle stack (hadrons and muons) resp. on the CORSIKA-EGS4 stack (electromagnetic particles). At this point to each particle with  $E < E_{low}$  a weight can be attributed which is used in the THIN option (see Sect. 4.55 page



96) to treat the bulk of the low-energy particles by the standard Monte Carlo method (with thinning) until these particles reach the observation level. For an equivalent precision level a minimum speed-up by a factor 5 can be expected using CONEX instead of standard Monte Carlo methods (including optimized thinning).

The parameters needed for the CONEX code are derived from the settings for the standard CORSIKA Monte Carlo procedure (including the THIN option) and are completed by default values which cover all normal applications. Only in special cases these standard settings may be modified by experts using the keywords CASCADE, CONEX, CX2COR, and CXWMX (Sects. 4.50 - 4.53, pages 93 - 95).

The CONEX option is only available for the high-energy hadronic interaction models EPOS, QGSJET, QGSII, or SIBYLL 2.3d in combination with all low-energy interaction models. If the CONEX option is selected and none of the special keywords are used, the simulation will include the use of cascade equations at intermediate energy and the weights are defined by the usual THIN, THINH and THINEM keywords.

Together with EPOS, the CONEX option allow to use primary mass up to 250 nucleons, but this has to be explicitly selected when extracting the FORTRAN code from the source file because of possible high memory consumption ( $\geq 2$ GB).

The CONEX option is compatible with COAST (or COASTUSERLIB). CoREAS should be used with care since the cascade equations are not used in the radio signal emission.

The CONEX option cannot be combined with any of the CERENKOV options (Sect. 3.4.1 to Sects. 3.4.6, pages 42 - 47) or UPWARDOLD.

To obtain this program version the CONEX option has to be selected together with the CURVED+SLANT+THIN+UPWARD options when extracting the FORTRAN code from the source file.

### 3.5.11 COREAS Option

The COREAS option activates the CoREAS code [53] for the calculation of radio emission from extensive air showers. The emission is calculated by applying the "endpoint formalism" [54] to each individual path element tracked in CORSIKA. CoREAS is the successor of REAS [55] and is not based on any particular model for the generation mechanism of the radio emission (except of course classical electrodynamics). Please refer to the separate user manual for CoREAS, which is automatically generated when compiling CORSIKA with the COREAS option active.

For high-precision work, it is possible to use an altitude dependent refractive index. The refractive index profile is used in CORSIKA as well as in the CoREAS simulations of the radio emission, where it was found that slight variations in refractive index (at the level of  $4\% \times \ln(n - 1)$ ) lead to noticeable changes in the radio signal. This is a systematic effect that is taken into using the profiles from `gdastool` see Sect. 4.23 page 83 and Sect. G page 185.

The COREAS option<sup>50</sup> cannot be combined with any other option making use of the COAST interface, i.e. COASTUSERLIB, INCLINED, or ROOTOUT options, nor with

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<sup>50</sup>This option compiles only on LINUX and Mac-OSX machines.

histogram producing options ANAHIST, AUGERHIST, or MUONHIST and also the COMPACT option..

To obtain this program version the COREAS option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.12 CURVED Atmosphere Option

The standard CORSIKA program models the Earth's atmosphere as a flat disk where the density of the air decreases with the height. The shower calculations and tracking are using Cartesian coordinates. In a flat atmosphere the thickness increases with  $1/\cos\theta$ . This is a good approximation for inclined showers if their zenith angles are below  $\approx 70^\circ$ . Above this value the differences between a flat and a curved atmosphere become more and more important. At  $90^\circ$  eventually the thickness of the flat atmosphere becomes infinite whereas the correct thickness is  $\approx 37000 \text{ g/cm}^2$ .

Within the CURVED atmosphere option for large zenith angles above  $\approx 70^\circ$  the Earth's atmosphere is no longer assumed to be completely flat (as in the standard version for smaller zenith angles). Rather the atmosphere is replaced by a 'sliding plane atmosphere'. Each time the horizontal displacement of a particle exceeds a limit of 6 to 20 km (depending on altitude), a transition to a new, locally plane atmosphere is performed. By these means the advantages of the simpler transport formulas within a planar atmosphere are combined with the faster simulation speed by avoiding the lengthy and more complicated treatment by using a true spherical system. Because of technical reasons only one observation level may be specified in the CURVED option and the X and Y coordinates have a special definition in order to be able to reconstruct the 3 coordinates (x,y,z) in a cartesian frame (see Sect. 4.72 page 105).

In the CURVED option the ionization energy loss, deflection within the Earth's magnetic field, and the generation of Cherenkov photons is enabled for charged hadronic primaries on their path between entering the atmosphere and the first interaction. (The deflection is disabled in the standard version without using keyword TSTART, page 78.) The arrival time refers to the start at the margin of the atmosphere, which is indicated by a negative value of element 7 of the event header block (page 139).

The NKG formulas do not take into account the curvature of the Earth's surface. Therefore the NKG output is suppressed in the CURVED option.

The combination of the CURVED option with the UPWARD option is described in Sect. 3.5.38 (page 70).

As in the CEFFIC option the atmospheric absorption is treated only in a manner suited for planar atmospheres, you should not combine the CURVED option with CEFFIC.

The CURVED option needs the atmospheres in the 5-layer model rather than interpolated from ATMEXT tables. Therefore in the CURVED option AATM, BATM, CATM values are fitted to the tabulated atmospheres and the table interpolation is disabled.

To obtain this program version the CURVED option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.13 DYNSTACK

The DYNSTACK option enables the user to change and modify the internal particle stack with a prebuild API (Application programming interface). For experiments where the observed area is much smaller than the area covered by the shower or tight boundaries are present it is possible to define constraints on the calculated particle which can lead to a large runtime reduction. Two possible implementations are already delivered with CORSIKA. This includes an output routine for every particle into a file and a storage that reproduces the default CORSIKA behavior with the possibility to set its size in the steering card to reduce HDD access.

For a full manual/documentation and continuous updates visit the Github web page

<https://github.com/tudo-astroparticlephysics/Cor-PlusPlus>\\

or contact D. Baack<sup>51</sup>.

Any removal or modification of particles during the simulation can impact physical correctness. Select all changes with care and test your results.

To obtain this program version, the DYNSTACK option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.14 EFIELD Option

In the EFIELD option the presence of atmospheric electrical fields acts on the movement of the charged particles of the shower. Depending on the direction of the electrostatic field relative to the movement of the charged particles the positively (resp. negatively) charged particles are deflected, accelerated, or slowed down. In CORSIKA versions > 7.56 these effects are considered not only for positrons/electrons, but also for muonic and charged hadronic particles. The shape and functional form of the electrostatic field in the atmosphere must be specified by the user in the C-subroutine *elfield.c* (resp. in the FORTRAN subroutine *elfield.f*). An *elfield.c* template (with zero electrical field) (resp. *elfield.f*) is available in the `src/` subdirectory. The used variables for input and output with their definitions are explained in the comment lines of this subroutine. Further details on the EFIELD option may be found in Ref. [56].

The EFIELD option should only be used if the electromagnetic shower component is treated with the EGS4 routines. The EFIELD option has no effect on the calculation of the electromagnetic shower component with the analytic NKG formulas.

To obtain this program version the EFIELD option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.15 EHISTORY Option

With this option additional information on the prehistory of muons, neutrinos, and electromagnetic (EM) particles is given out to the particle output file. This extended addi-

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<sup>51</sup><dominik.baack@tu-dortmund.de>.

tional information is activated for the muons by the keyword MUADDI (see page 102) and for the neutrinos by the keyword NUADDI (see page 103). For EM particles it is activated with the keyword EMADDI (see page 102). The precursor<sup>52</sup> hadrons 'grandmother' and 'mother' of the particles at ground are specified with the penetrated matter between start and end of grandmother's track, the position of mother's birth, and the momenta of grandmother and mother at that point. The full generation counter  $G$  of the particle at ground is also given.  $G$  can be decomposed into different components  $G = \mu \times 10^9 + \nu \times 10^8 + \gamma \times 10^6 + e \times 10^3 + h$  from which we get the history of the recorded particles:

- $\mu$  is the number of  $\gamma \rightarrow \mu^+ + \mu^-$  interactions,
- $\nu$  is the number of photonuclear interactions,
- $\gamma$  is the number of photon interactions,
- $e$  is the number of electron or positron interactions (including the emission of low energy particles),
- $h$  is the number of hadronic interactions/decays. Usually +1 per interactions but +31 is used for charmed particle decay, +51 in case of pion decay, +100 for muon decay and +200 for muon bremsstrahlung or pair production such that only EM coming from muons have  $h > 100$ .

counters of the grandmother particle are also given, which might be compared to the generation counters of the particle at ground, thus indicating additional decays between mother particle and the final particle, which otherwise are not given. Further details are described in [57].

The FORTRAN program *corsikaread\_history.f* available in the subdirectory `src/utills/` is designed to read binary particle output files produced with the EHISTORY option.

The EHISTORY option cannot be combined with any other option which produces histograms (ANAHIST, AUGERHIST, or MUONHIST) or plots (PLOTSH, PLOTSH2), nor with PARALLEL, CONEX, or CERENKOV options.

To obtain this program version the EHISTORY option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.16 ICECUBE1 and ICECUBE2 Options

The options ICECUBE1 and ICECUBE2 have been developed to meet the special requirements of the IceCube experiment [58]. The discrimination of high-energy extra-terrestrial neutrinos from the steeply falling background of atmospheric neutrinos is a challenging task. Neutrinos which interact within the fiducial volume of the IceCube detector are assigned to be atmospheric, if they are accompanied by a detectable muon from the same

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<sup>52</sup>For dpmjet III which handles internal decays by PYTHIA the direct parent particle of the decay leading to the final muon resp. neutrino is taken as mother particle.

cosmic ray cascade which produced the neutrino. The probability for such events may be calculated analytically [59], but should be confirmed by Monte Carlo methods.

In the ICECUBE1 option the stack of secondary particles (which usually is handled in the last-in-first-out mode) is replaced by the first-in-first-out sequence which postpones those particles exciting mostly low-energy subshowers not leading to the emission of a high-energy neutrino, while the high-energy particles are treated preferentially. To discard a non-interesting shower because of missing high-energy neutrinos can be decided much earlier, thus saving computing resources. The necessary input parameter of the threshold energy above which particles are interesting to produce a high-energy neutrino may be adjusted using the keyword EINTER (see Sect. 4.104 page 117).

The ICECUBE2 option permits a compression of the particle data output file employing the gzip procedure (keyword COMPRESS see Sect. 4.106 page 117). Alternatively the output file is created as a pipe buffer instead of a regular file, and a detector simulation program can read from this pipe on-line saving useless disk files. This pipe is activated using the keyword PIPE (see Sect. 4.105 page 117).

Both options ICECUBE1 and ICECUBE2 may be activated simultaneously or independently.

To obtain these program versions the ICECUBE1 resp. ICECUBE2 options have to be selected in *./coconut* when extracting the FORTRAN code from the source file.

### 3.5.17 INCLINED Observation Plane Option

In the INCLINED option, in addition to the horizontal observation plane (by default) an inclined observation plane is used to produce a second output. The keyword INCLIN explained in Sec. 4.71 (page 104) is used to specify the geometry of the inclined sampling plane. It can be defined anywhere compared to the shower core on the lowest observation plane, with any orientation, to get particles on a mountain slope for instance, but also very conveniently, as the shower plane perpendicular to the shower axis at a given depth (either slant (along the shower axis) or vertical (fixed height)). Particles are counted on both sides of the plane (longitudinal momentum is positive when the particle goes in the direction of the normal, negative otherwise), allowing a longitudinal view of the shower of the shower axis is “in” the plane. This means that this inclined observation plane do not stop the shower development. Only the normal lowest observation level (see UPWARD and UPWARDOLD page 70), the border of the atmosphere, or time limitation (see TIMLIM 100 can stop the particle tracking).

The output is written as a binary data file to the file named *'DATnnnnnnn.inclined'* (see Sect 2.4.2 page 28) (resp. *'DATnnnnnnnnn.inclined'* with option NRREXT, see Sect. 3.5.24 page 61).

The format is identical to the normal binary output data, besides the fact that all coordinates are given within the inclined plane. Further details on the output file format are given in Sect. 10.3 (page 148).

Technically, the INCLINED option uses a particular COASTUSERLIB library (Sec. 3.5.8 page 51), namely *InclinedPlane*, which is shipped with COAST in the *CorsikaOptions*

subdirectory by default.

The INCLINED option<sup>53</sup> cannot be combined with the COASTUSERLIB, COMPACT, COREAS, or ROOTOUT options.

To obtain this program version the INCLINED option has to be selected in *./coconut* when extracting the FORTRAN code from the source file.

### 3.5.18 INTTEST Interaction Test Option

With this option the interaction model routines can be tested which are used to describe the hadronic collisions. Only the first interactions are simulated and no air showers are developed. Various projectiles (keyword PRMPAR) and targets (keyword INTTST) may be selected. The result of the test is a series of histograms generated with the HBOOK routines [33] and written to the unit HISTDS. The histograms show properties of the secondary particles produced in the first interaction e.g. distributions of transverse momenta, of squared transverse momenta, of longitudinal momenta, of Feynman x-distributions, of rapidity and pseudo-rapidity distributions, and various particle multiplicity distributions. To obtain this program version the INTTEST option has to be selected when extracting the FORTRAN code from the source file. The HBOOK routines require linking of the CERN library<sup>54</sup> with the program; they are not supplied with the CORSIKA package. Via the keywords with their parameters (pages 121 - 123)

INTTST	ITTAR	MCM		
INTDEC	LPIO	LETA	LHYP	LKOS
INTSPC	LSPEC			
DIFOFF	NDIF			
TRIGGER	NTRIG			
HISTDS	HISTDS			

the conditions of the interaction test run have to be specified.

Because of its permanent modification a comprehensive description of the INTTEST option is not available. It should be noted that the combinations of allowed parameters for the INTTEST option differ from the standard version without explicitly noting this or checking this during the program run.

To obtain this program version the INTTEST option has to be selected when extracting the FORTRAN code from the source file.

---

<sup>53</sup>This option compiles only on LINUX and Mac-OSX machines.

<sup>54</sup>The CERN library is only available for 32bit computers.

### 3.5.19 LPM Option

The LPM option switches on the **Landau-Pomeranchuk-Migdal** effect (which is added to EGS4) causing an effective reduction of the pair production and bremsstrahlung cross-sections [60] at the highest energies. With this option it is possible to include the LPM-effect without selecting the THIN option (see Sect. 3.5.36 page 68), the MULTITHIN option (see Sect. 3.5.20 page 59), or the PARALLEL option (see Sect. 3.5.26 page 61).

### 3.5.20 MULTITHIN Option

In the MULTITHIN option showers are simulated unthinned in the standard manner, but for each particle additionally 7 words (one data sub-block) are written which contain the weights of this particle for up to 6 different thinning modes (see Table 14 page 143 and Ref. [61]). A negative weight indicates that the particle did not survive in that mode, its amount gives the weight at the moment of elimination. The different thinning modes are steered by the input parameters of the MTHINH keyword (Sect. 4.58 page 98). For each thinning mode a dedicated random generator sequence is used. The different seeds for the thinning modes are initialized by the MSEED keyword (see Sect. 4.59 page 98). To have all the thinned particles at the maximum weight and reduce further the number of particles or reduce the maximum weight for a given simulation time, it is possible to activate the “THINning with all particles at MAXimum weight” (THINMAX) option. This is possible by allowing a small energy conservation violation while the “normal” hillas algorithm conserve energy perfectly. To avoid large particle output files the particles close to the shower axis may be suppressed using the keyword MTHINR (which acts identically as keyword CORECUT, see Sect. 4.62 page 99). With this keyword a maximum radius RCUT may be defined within which the output particles are discarded (so it cuts the unthinned shower around the core) (see Sect. 4.62 page 99).

The MULTITHIN option may be combined with all other options except those which demand the normal thinning (i.e. ANAHIST, AUGERHIST, CONEX, INTTEST, or THIN). In the MULTITHIN option the LPM-effect is activated automatically. In the particle output file each MULTITHIN sub-block (see Table 14 page 143) will be preceded by a particle data sub-block with the particle arriving at the detector level or - in case of the MUPROD option - with a decaying muon at track end [57].

If the COASTUSERLIB option is selected, one can choose which thinning selection should be used in COAST with the keyword MWEIC and if the weight sub-block should be saved in the final output (thinned shower used for COAST only with a normal DAT output file without additional weight sub-block) (Sect. 4.61 page 99).

To obtain this program version the MULTITHIN option has to be selected when extracting the FORTRAN code from the source file. The THINMAX option should be selected together with the MULTITHIN option to be active.

### 3.5.21 MUONHIST Option

The MUONHIST option produces a series of histograms generated with HBOOK routines [33]. The histograms give various properties (energy, transverse momentum, slant depth) of the muons at the position of their origin. An example of the application of this option is given in Ref. [62].

The histograms are written into the file named '*datnnnnnnn.lhbook*' (page 150) onto the output directory DSN specified by the keyword DIRECT (page 108). Using the option NRREXT (see Sect. 3.5.24 page 61) the file is named '*datnnnnnnnnnn.lhbook*'. To suppress the ordinary particle output file the keyword PAROUT (page 108) might be used. To obtain this program version the MUONHIST option has to be selected when extracting the FORTRAN code from the source file. The HBOOK routines require linking of the CERN library<sup>54</sup> with the program; they are not supplied with the CORSIKA package.

### 3.5.22 MUPROD Option

With this option additional information on the prehistory of muons is given out to the particle output file even if those muons don't reach the observation level. To get this information the MUPROD option has to be selected and activated by the MUADDI keyword (see Sect. 4.68 page 102). This information is written to the particle output file, where the first entry with particle code 85 (for  $\mu^+$ ) resp. 86 (for  $\mu^-$ ) represents the muon with its momentum and position components at the production point. To get more information the EHISTORY option (see Sect. 3.5.15 page 55) can be combined with the MUPROD option. In the following entries the precursor hadrons 'grandmother' and 'mother' of the decaying muon are specified with the penetrated matter between start and end of grandmother's track, the position of mother's birth, and the momenta of grandmother and mother at that point. The complete generation counter of the grandmother is also given, which might be compared to the generation counter of the muon thus indicating additional decays between mother particle and the muon, which otherwise are not given. The fourth entry with particle code 95 (for  $\mu^+$ ) resp. 96 (for  $\mu^-$ ) represents the muon at its decay resp. interaction point with its momentum and position components. Further details are described in [57].

### 3.5.23 NEUTRINO Option

Muonic and electronic neutrinos and anti-neutrinos originate from the decays of  $\pi^\pm$ , the leptonic decays of  $K^\pm$  and  $K_L^0$ , and the decays of  $\mu^\pm$ . The neutrino formation is simulated with exact kinematics of all two and three body decays taking into account the polarization of the muons. The neutrino trajectories are followed down through the atmosphere disregarding any interaction with the target nuclei of the air. The neutrinos are written to the particle output file using the particle type numbers 66 to 69, if their energy exceeds the cutoff energy ELCUT(1).

To obtain this program version the NEUTRINO option has to be selected when extracting the FORTRAN code from the source file.



### 3.5.24 NRREXT Option

The NRREXT option enables to use run numbers RUNNR (see Sect. 4.1 page 74) with 9 instead of 6 digits and the file names of the various output files get elongated from e.g. 'DATnnnnnnn' to 'DATnnnnnnnnnn' where nnnnnnnnnn gives the extended run number.

### 3.5.25 NUPRIM Option for Primary Neutrinos

With this program version showers induced by primary neutrinos can be simulated. Possible primary particles are (up to now)  $\nu_e$ ,  $\bar{\nu}_e$ ,  $\nu_\mu$ , and  $\bar{\nu}_\mu$ . If combining the NUPRIM option with the CHARM or TAULEP option also the  $\nu_\tau$  and the  $\bar{\nu}_\tau$  neutrinos can be handled.

Because of the low cross-sections of neutrino-induced interactions it is recommended to fix the height of the primary interaction using the keyword FIXHEI (page 79). As the first (neutrino-induced) interaction is handled by the HERWIG code [22], a link [23] with the HERWIG routines is necessary<sup>55</sup>. The type of interaction (charged current or neutral current) may be selected by the keyword NUSLCT (Sect. 4.43 page 91). HERWIG produces - besides others - secondary particles (with charm). Those particles can be treated by CORSIKA only if you combine the NUPRIM option with the CHARM option (see page 50). Without specifying the CHARM option those charmed secondaries decay at the vertex and are not explicitly transported in CORSIKA. Specifying the TAULEP option the  $\tau$ -leptons emerging from CC-interactions of  $\nu_\tau$  or  $\bar{\nu}_\tau$  neutrinos are treated explicitly in CORSIKA. The interactions of secondary particles coming from the primary neutrino reaction are treated by the selected high-energy resp. low-energy hadronic interaction model or by NKG/EGS4. Secondary neutrinos are explicitly generated only, if the NEUTRINO option (see Sect. 3.5.23 page 60) is selected, but they are not treated by HERWIG. Initiating a shower with non-neutrino primary particles will not call the HERWIG routines.

Sequence 5 of the CORSIKA random generator is foreseen for HERWIG, therefore you should initialize it (see Sect. 4.3 page 75).

For using the NUPRIM version you first have to select the NUPRIM option when extracting the FORTRAN code from the source file. The *Makefile* will compile the HERWIG FORTRAN routines and link them with your CORSIKA program.

A combination of the NUPRIM option with the INTTEST option is not possible.

### 3.5.26 PARALLEL and PARALLELIB Options

The PARALLEL option enables the CORSIKA program to run a single shower on several cores in parallel to reduce the time to complete the shower [27]. This enables the simulation at highest energies ( $> 10^{17}$ eV) without the THIN option. Therefore the LPM-effect is activated automatically.

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<sup>55</sup>The HERWIG 6521 package has been adapted as described in the README file which you find in the `herwig/` subdirectory (see Sect. A.1 page 162).

There exist two modes of parallelization:

- a) by shell scripts or
- b) by **M**essage **P**assing **I**nterface (MPI) system.

They are available after having selected the option "**b - PARALLEL treatment of subshowers (includes LPM)**" in the additional CORSIKA program options of `./coconut`. Then you have to select the mode of parallelization by answering the following question:

```
Compile CORSIKA as subroutine for parallelization with MPI ?
1 - Special stack for shell scripts without MPI [DEFAULT]
2 - Library to be used with MPI system
```

In the first case a set of shell scripts handles the external stack with the secondary particles and starts the subshower processes. These shell scripts are available in the subdirectory `src/parallel/` together with a short description (*userguide-pll-\*.txt*).

To obtain this program version the **PARALLEL** option has to be selected when extracting the FORTRAN code from the source file.

In the second case the **PARALLELIB** option is activated and the parallelization is handled by MPI. In this case the program *mpi\_runner.c* is compiled and linked with the subroutine *corsikacompfile.f*. The first interactions are simulated on a core (MASTER) and produce secondary particles which are used to start subshowers in parallel on many cores (SLAVES) to be run simultaneously. The communication between the MASTER and the SLAVES is established by MPI. Further technical details are explained in the MPI-Runner\_GUIDE [63] available in the subdirectory `doc/`.

To obtain this program version the **PARALLEL** and **PARALLELIB** options have to be selected when extracting the FORTRAN code from the source file.

It should be noted, that it is not foreseen to combine the **PARALLEL** or **PARALLELIB** options with the **CERENKOV** option.

### 3.5.27 PLOTSH Shower Plot Production Option

In the **PLOTSH** option the start and end points of each particle track are written to extra files (see Table 1 page 26); separate files are used for the electromagnetic, muonic, and hadronic component. Each track is written unformatted as one block consisting of 10 real\*4 numbers: Particle identification, Energy (in GeV), Xstart, Ystart, Zstart (all in cm), Tstart (in sec), Xend, Yend, Zend (all in cm), and Tend (in sec). In case of **THINning** the particle Weight is added at the end of each block (i.e. 11 real numbers).

Plots may be constructed from the content of these files by using the *plottracks* program. The program *plottracks3c.f* (available in `src/utils/`) converts the particle track maps generated by the **PLOTSH** option into PPM image files (readable by *xv*). *plottracks* is automatically compiled<sup>56</sup> when the **PLOTSH** option is selected.

The program *plottracks* reads in the '*DATnnnnnn.track\_em*', '*DATnnnnnn.track\_mu*', and '*DATnnnnnn.track\_hd*' files specified on the command line by the parameter

---

<sup>56</sup>`g77 plottracks3c.f -o plottracks`

run=nnnnnn (integer between 0 and 999999). Then it converts and combines them into 4 PPM images 'tracknnnnnn.em.ppm', 'tracknnnnnn.mu.ppm', 'tracknnnnnn.had.ppm', and 'tracknnnnnn.all.ppm' of the same resolution as the maps. By default, the electromagnetic, muonic, and hadronic maps are taken as the red, green, and blue channels of the RGB image, respectively.

As with increasing shower energy and decreasing threshold the number of tracks increases drastically you should simulate not more than 1 shower at a time to keep the output on the units 55, 56, and 57 at a tolerable size.

Via the keyword (page 123)

```
PLOTSH    T
```

the PLOTSH option is enabled or disabled.

This option is not recommended for ordinary shower production because of the large output files to be produced.

To obtain this program version the PLOTSH option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.28 PLOTSH2 Shower Plot Production Option

In the PLOTSH2 option air showers are visualized in a simple way, without having to deal with a (very large) track file. With the PLOTSH2 option map files are generated separately for the electromagnetic particles, muons, and hadrons. For each particle species maps are generated in all three projections (x-y, x-z, and y-z). The map files are written onto the directory DSN via the unit 55 and are named 'DATnnnnnn.<spec>\_<proj>.map', where <spec> stands for 'em', 'mu', or 'hd', and <proj> stands for 'xy', 'xz', or 'yz'. nnnnnn is the run number specified in the keyword RUNNR (see page 74). These map files are, basically, two-dimensional histograms containing the number of tracks in each xy-/xz-/yz-bin. The resolution of the maps is set via the three integer parameters IXRES, IYRES, and IZRES in the COMMON /CRPLOTSH2/. The unformatted map files consist of a two-word header containing the horizontal and vertical resolution of the respective map file, followed by the rows of map values<sup>57</sup>. These files are then easily processed further, for example by *map2png* (see below).

To control the plotting, the keyword PLAXES followed by 6 parameters (see page 123) might be used in the input file:

```
PLAXES    -5.E5  5.E5  -5.E5  5.E5  0.  3.E6
```

The keyword

```
PLCUTS    0.3  0.3  0.003  0.003  1.E5  T
```

---

<sup>57</sup>Note that writing binary files by FORTRAN, the header and each row are preceded and followed by an additional word.

(see page 124) serves to define energy cuts in the same order as those for the keyword ECUTS (hadrons, muons, electrons, photons) (see page 100). An additional parameter defines an upper bound on the time (in ns) passed since the first interaction, while the final logical determines whether only track segments inside the three-dimensional box given by the axis ranges should be plotted.

Via the keyword

```
PLOTSH    T
```

the PLOTSH2 option is enabled or disabled (see page 123).

This option is not recommended for ordinary shower production.

To obtain this program version the PLOTSH2 option has to be selected when extracting the FORTRAN code from the source file.

The C-program *map2png* (available in `src/utils/`) converts the particle track maps generated by the PLOTSH2 option into PNG image files. The only prerequisite needed by the program is *libpng*, which is included in virtually all recent LINUX distributions. If this library is not available for your system, please use PLOTSH option instead. *map2png* is automatically compiled<sup>58</sup> when the PLOTSH2 option is selected.

The program *map2png* reads in the map files specified on the command line (via the path to *DATnnnnnn*) and combines them into a single PNG image of the same resolution as the maps. By default, the electromagnetic, muonic, and hadronic maps are taken as the red, green, and blue channels of the RGB image, respectively. However, it is possible to specify other colors to be used for the different particle types. The projection to be used as well as whether to use a linear or logarithmic color scale can be specified on the command line. To see a list of options, simply run *map2png* without arguments.

### 3.5.29 PRESHOWER Option

The PRESHOWER option selects code and C-routines [64] to describe the pair production of ultra-high energetic ( $E_0 > 10^{19}$  eV) primary photons and bremsstrahlung interactions of the secondary em-particles within the Earth's magnetic field before reaching the top of atmosphere. Within the atmosphere the resulting swarm of em-particles is treated as one shower. To model correctly the field components of the Earth's magnetic dipole field in the outer space additional input is required to specify the location and the time of the experiment. The keyword (page 84)

```
GCOORD    -69.585  -35.463  2003.  1  0
```

gives the longitudinal and lateral position of the experiment on the Earth's globe, the year (the Earth's magnetic field is time dependent), a print indicator, and a stop indicator (for the case no pre-showering occurred). Details of this input are given in page 84. With

---

<sup>58</sup>`gcc [-Lpath_to_libpng] -lpng -lm map2png.c -o map2png`

The bracketed option is only needed if the path to 'libpng.so' is not included in the LD\_LIBRARY\_PATH variable.

the stop indicator it is possible to skip those events where no pre-showering occurred. Nevertheless the event header and event end sub-blocks are written to MPATAP.

The appropriate C-routines of the preshower2.0 version belonging to the PRESHOWER option are available in the `src/preshower2.0-rev/` subdirectory, compiled and linked by the *Makefile*.

To obtain this program version the PRESHOWER option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.30 REMOTECONTROL

The REMOTECONTROL option enables the user to connect the simulation to an external control server. This makes it possible to exchange information during a simulation run or receive data to control the simulation.

For a full manual/documentation and continuous updates visit the Github web page <https://github.com/tudo-astroparticlephysics/Cor++>. or contact D. Baack<sup>59</sup>.

To obtain this program version, the REMOTECONTROL option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.31 RIGIDITY Option

The RIGIDITY option checks the energy of low-energy primary particles whether they would reach the Earth surface because of their magnetic rigidity which depends on the local Earth magnetic field and the angle of incidence. Particles not reaching ground will be rejected. In the present version of RIGIDITY the parameters are hard-coded for the GRAPES-3 experiment in Ooty, India. The table *gr3.tex* giving the dependence on the zenith and azimuth angles is read in via unit 99.

To obtain this program version the RIGIDITY option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.32 ROOTOUT Option

The ROOTOUT option selects code which directly transmits the particle output (normally directed to MPATAP) to C++-routines<sup>60</sup> to write an output '*DATnnnnnnn.root*' file in *root* format, e.g. for *root* off-line analysis of the particle output data without storing the huge particle output data file MPATAP. The Cherenkov output file MCETAP is not affected by the ROOTOUT selection.

The default COAST ROOTOUT data format is thought for fast and easy off-line analysis of CORSIKA data and is NOT a general replacement for CORSIKA binary data files. Due to the chosen data structure one shower including all particles and all Cherenkov photons is stored in the computer's memory entirely before it is written to disk. This needs a lot of memory for high energy and/or high quality showers.

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<sup>59</sup><dominik.baack@udo.edu>.

<sup>60</sup>Through the enclosed COAST package.

The ROOTOUT option<sup>61</sup> cannot be combined with any other option making use of the COAST interface, i.e. COASTUSERLIB, COREAS, or INCLINED options, nor with histogram producing options ANAHIST, AUGERHIST, or MUONHIST and also the COMPACT option..

To obtain this program version the ROOTOUT option has to be selected in expert mode of *coconut* (`./coconut -e` because of its compatibility with older version on *root* only) when extracting the FORTRAN code from the source file and *root* v5.x or earlier (see `root.cern.ch`) must be installed on your system.

To obtain *root* output file, it is recommended to save normal DAT output files and to convert them using the tools available in `src/utis/coast/CorsikaRead` which are more flexible and can be compiled with current version of *root*. To develop your own CORSIKA file reader on the basis of these files, a doxygen documentation is available in `coast/Documentation` if the information in the files are not enough. After at least one use of `./coconut`, please use the *make* command in this subdirectory. Then you can find all details using `firefox coast/Documentation/COASTDocumentation/html/index.html`

### 3.5.33 SLANT Option

With the SLANT option the longitudinal distributions (page 101 and Sect. 10.5 page 151) are given in slant depth bins along the shower axis instead of vertical depth bins used in the standard case. This slant depth scale is more appropriate to investigations of very inclined showers.

In the SLANT option [65, 66] the ionization energy loss, deflection within the Earth's magnetic field, and the generation of Cherenkov photons is enabled for charged hadronic primaries on their path between entering the atmosphere and the first interaction. (The deflection is disabled in the standard version without using keyword TSTART, page 78.) The arrival time refers to the margin of the atmosphere, which is indicated by a negative value of element 7 of the event header block (page 139).

To obtain this program version the SLANT option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.34 STACKIN Option

With the STACKIN option the parameters of secondary particles will be read into the CORSIKA stack. Thus interactions of very exotic primaries may be treated off-line in a suitable interaction program avoiding a direct coupling of such programs with CORSIKA. The air shower generated by these resulting secondary particles is simulated in CORSIKA, and all options but NUPRIM, PRESHOWER and INTTEST may be combined with STACKIN. To characterize the altitude of the first interaction the keyword FIXHEI<sup>62</sup>

---

<sup>61</sup>This option requires a C++ compiler and compiles only on LINUX and Mac-OSX machines.

<sup>62</sup>With the keyword FIXHEI the parameter N1STTR defining the type of the first target should be set to 0 (random selection) as it acts on that secondary particle which is treated first.

(page 79) must be used, the shower axis is defined by the angles THETAP and PHIP (pages 77 and 77). This externally treated first interaction starts the clock by default (the keyword TSTART (page 78) is disabled). This means that the starting point of the slant longitudinal profile is also defined by FIXHEI in that case (and only in that case). The *all-inputs-st* file in the `run/` subdirectory is an example input file for a run with STACKIN option.

To run a series of showers with fixed first hadronic interaction the keyword OUTFILE (page 80) may be used first in a run without STACKIN option to write out the parameters of the secondary particles to a file (see page 28) which is read-in later with the STACKIN option thus keeping the first interaction independent of the used random seed.

The STACKIN option can also be used to replace the FIXCHI keyword when COAST is used (COASTUSERLIB or INCLINED option for instance). If the input file contains only one primary particle, it will start its propagation at the point defined by FIXHEI.

The file containing the parameters of the particles has to be specified by the keyword INFILE (page 80) and is read in via the logical unit LSTCK (by default 23, page 24). The first line is read with free format (with a leading blank character) and contains the number of secondaries and the primary energy. The following lines are read with the format (2I5,4(1X,E15.7)) containing current particle number, particle type (see Table 4 page 129), total energy (GeV), longitudinal momentum (GeV/c), and transverse momenta (GeV/c). The momenta are taken relative to the direction of the shower axis (direction of the exotic particle). The *all-inputs-stackin* file in the `run/` subdirectory is an example input file.

To obtain this program version the STACKIN option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.35 TAULEP Option

In the standard CORSIKA program the  $\tau$ -leptons cannot be treated. The TAULEP option [49] (as well as the CHARM option) enables the treatment of the  $\tau$ -leptons, their decays are treated by the PYTHIA package [34]. Therefore (except in combination with the DPMJET high-energy interaction code) the PYTHIA 6.411 package in the `pythia/` subdirectory is automatically compiled and linked. In combination with DPMJET the implicit PYTHIA 6.427 package is used and needs no extra linking. The TAULEP option may be combined with all high-energy interaction models.

All possible interactions (bremsstrahlung, pair production, nuclear interactions) of the  $\tau$ -leptons are respected in a manner analogously to the  $\mu$ -leptons. For the treatment of primary  $\nu_\tau$  and  $\bar{\nu}_\tau$  neutrinos the TAULEP (or CHARM) option has to be combined with the NUPRIM option (see page 61). For tracing back PYTHIA problems the PYTHIA keyword is available (page 121).

To obtain this program version the TAULEP option has to be selected when extracting the FORTRAN code from the source file. Please verify that the PYTHIA 6.411 package in the `pythia/` subdirectory is correctly compiled and linked.

### 3.5.36 Option for THINning

For primary energies  $E_0 > 10^{16}$  eV the computing times become excessively long (they scale roughly with the primary energy). To reduce the times to tolerable durations the so-called 'thin sampling' mechanism (also named 'variance reduction' [20]) is introduced [60]. When thinning is active all particles below the adjustable fraction of the primary energy (thinning level  $\varepsilon_{th} = E/E_0$ ) which emerge from an interaction are exposed to the thinning algorithm. Only one of these particles is followed and an appropriate weight is given to it, while the other particles below the thinning level are dropped. Details on this formalism may be found in Refs. [20, 60, 67].

A further improvement [68] to reduce undesired statistical fluctuations of particle densities far from the shower core uses a limitation of the weights. Particles emerging from an interaction which would exceed a specified weight limit are excluded from the thinning algorithm. Using different weight limits for em-particles and hadronic (including muonic) particles enables a drastic reduction of computing time, if the user's interest is focused onto a precise lateral distribution of muons on the expenses of larger fluctuations of the em-part. To have all the thinned particles at the maximum weight and reduce further the number of particles or reduce the maximum weight for a given simulation time, it is possible to activate the "THINning with all particles at MAXimum weight" (THINMAX) option. This is possible by allowing a small energy conservation violation while the "normal" hillas algorithm conserve energy perfectly.

A third algorithm to save space on disks reduces the number of particles close to the shower core where anyway the detectors will saturate. Particles arriving at the detector level within a selected core distance  $rmax$  are selected at random with a probability  $prob \propto (r/rmax)^4$  and, when surviving<sup>63</sup>, their weight factor is multiplied with the inverse of this probability, irrespective of exceeding the weight limit. This radial thinning is not effective for the table output '*DATnnnnnnn.tab*' (see keyword PAROUT page 108 and Sect. 10.6 page 151), nor for the '*datnnnnnnn.lhbook*' file generated in the ANAHIST, AUGERHIST and/or MUONHIST versions. The radial thinning is disabled when the CORECUT option is used.

Via the keyword (page 96)

THIN    EFRCTHN    WMAX    RMAX

you may specify the energy fraction EFRCTHN of the primary energy, below which the thinning process becomes active. Above this energy no thinning will take place. WMAX gives the maximum weight factor<sup>64</sup>, which should not be exceeded. The core distance up to which the radial thinning at detector level takes place is specified by RMAX.

Via the keyword (page 97)

THINH    THINRAT    WEITRAT

---

<sup>63</sup>To ensure the surviving of enough particles close to the shower axis RMAX should be chosen not too large.

<sup>64</sup>In case of a primary energy spectrum (ULIMIT > LLIMIT) the value of WMAX is used for the low energy end of the energy interval and slides to higher values according with the sliding primary energy.



or the keyword (page 97)

THINEM    THINRAT    WEITRAT

differing thin levels and weight limits can be specified for hadronic or electromagnetic particles overwriting the ratios

$$\text{THINRAT} = \varepsilon_{th_{em}}/\varepsilon_{th_{hadr}} \quad \text{and} \quad \text{WEITRAT} = \text{WMAX}_{em}/\text{WMAX}_{hadr}$$

which are set to 1 by default.

$\varepsilon_{th}$	none	$10^{-6}$	$10^{-5}$	$10^{-4}$	$10^{-3}$
Time (min)	98	51	7.2	1.2	0.16
particles	413078	58313	11466	2211	419

Table 2: Computing times and number of particles for various thinning levels, without application of weight limits and radial thinning.

The effect of various thinning levels  $\varepsilon_{th} = E/E_0$  on the computing time and the number of particles on MPATAP may be seen from Table 2 which is established with default CORSIKA parameters (QGSJET option, EGS4 activated) for vertical proton induced showers of  $10^{15}$  eV primary energy (DEC 3000/600 AXP with 175 MHz), with  $\varepsilon_{th_{em}}/\varepsilon_{th_{hadr}} = 1$ , with infinite weight limit, and without radial thinning. As to each particle an appropriate weight must be attributed, the output format described in Sect. 10.2 (page 136 ff.) has to be changed to incorporate this additional parameter of each particle. Consequently MPATAP and MCETAP **output data generated with the THIN option differ** from those generated in simulations without this option (see also Sect. 10.2.3 page 144).

An optimum choice of the various thinning parameters depends on the information which should be drawn from the simulations. To minimize the additional fluctuations (caused by the thin sampling algorithm) for muonic particle densities at large distances  $> 300$  m from the shower core - which is one of the problems of the Auger experiment - a suitable setting [68] of WEITRAT would be

$$\text{WEITRAT} \approx 100$$

while the choice of WMAX is optimized for the primary energy  $E_0$  (given in GeV) and the selected thinning level EFRCTHN for em-particles to

$$\text{WMAX} = \text{EFRCTHN} \cdot E_0 .$$

The *all-inputs-thin* file in the *run/* subdirectory is an example input file for a run with THIN option.

To obtain this program version the THIN option has to be selected when extracting the FORTRAN code from the source file. The THINMAX option should be selected together with the THIN option to be active.

### 3.5.37 TRAJECT Option

The TRAJECT option selects the zenith and azimuth angles in a manner that a gamma (or neutrino) source is followed along its trajectory in the sky. The angles selected by this option override the angle selection according to the VIEWCONE (see Sect. 3.5.40 page 71), VOLUMEDET, and VOLUMECORR (see Sect. 3.5.41 page 72) options or selected by the keywords THETAP (page 77) and PHIP (page 77). This option is activated/disabled by the keyword TRAF LG (page 114). The keywords SRCPOS (page 115), TRATM (page 115), TLAT (page 115), and TLONG (page 116) specify the position of the source in equatorial coordinates, the time information on the observation, and the latitude and longitude of the observing telescope on the globe. The total observation time given by the parameter TDURATION of the keyword TRATM (page 115) specifies the span of the trajectory. This span contains the number of events defined by the keyword NSHOW (page 75).

As the CORSIKA coordinate system refers to magnetic North, the magnetic declination at the telescope site must be specified by the keyword GEODEC (page 116) for a correct transformation from the equatorial coordinates to the CORSIKA coordinate system. For extended sources (like a nebula) moving in the sky a broader trajectory may be followed by specifying a spread radius with the keyword TRARAD (page 116).

A detailed description of the TRAJECT option is given in Ref. [69] which gives also an illustrative example.

To obtain this program version the TRAJECT option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.38 UPWARD Option

The UPWARD option (without CURVED option) selects code which treats the upward traveling particles. For primary particles the zenith angle is restricted to  $0^\circ < \theta < 70^\circ$  and  $110^\circ < \theta < 180^\circ$ . No additional keyword has to be specified.

The UPWARD option might be combined with the CURVED option (page 54). This enables to start showers with arbitrary zenith angles  $0^\circ < \theta < 180^\circ$  and secondary particles with arbitrary zenith angles are followed. [66].

For showers with skimming incidence (zenith angle =  $90^\circ$ ) the minimum altitude of the shower axis above sea level is specified by HIMPACT(*i*) (keyword IMPACT, page 79) and defines the geometry completely. The zenith angle at the entrance into the atmosphere is calculated with it. The angles THETPR(*i*) (see keyword THETAP page 77) have no meaning and are overridden. The keyword IMPACT may be combined with the keywords FIXHEI (page 79) or FIXCHI (page 78) to start the shower before reaching the minimum altitude of the shower axis.

For showers with zenith angles  $> 90^\circ$  (e.g. initiated by neutrinos, page 61) the starting point of the shower resp. the first interaction can be defined by the keywords FIXHEI (page 79) or FIXCHI (page 78, but only if COASTUSERLIB is NOT selected), but otherwise will start at the height of the lowest observation level, and the other observation levels can be defined freely, like without the UPWARD option. Note that only particles coming from

above these observation planes are recorded, so if the shower is too far from the observation level and the zenith angle too large, no particle might be recorded in the particle output file (ground level), but the longitudinal profile is recorded. In that case, the use of COAST based options (page 51) like CoREAS for the radio signal or INCLINED (page 104) to define another observation plane, which might be hit by the shower, is recommended.

To obtain the UPWARD program version the UPWARD option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.39 UPWARDOLD Option

This is the former option for upward going showers with zenith angles  $> 90^\circ$  in which the observation level (page 103) must be chosen either at the top of the atmosphere, or at minimum above the starting point of the shower. (The shower axis must cross the observation level.)

With this option, the particles coming from below the (unique) observation level are recorded in the output file and the particles will stop the propagation here. So the core which define the origin of the coordinate frame is defined as  $(X=0, Y=0, Z=OBSLEV)$  where the shower axis hit this observation level above the starting point of the shower (page 125. This starting point resp. the first interaction must be defined by the keywords FIXHEI (page 79) or FIXCHI (page 78, but only if COASTUSERLIB is NOT selected). This option can be combined with the INCLINED option (page 104) to define a more realistic observation plane, the propagation of the upward going shower being stop by the observation level which can be define directly above the region of interest (to avoid unnecessary calculation). Note that because of the different coordinate system, the position of the plane has to be defined in a different way than with the UPWARD option.

To obtain the UPWARDOLD program version the UPWARDOLD option has to be selected when extracting the FORTRAN code from the source file.

The UPWARDOLD option is not compatible with the CONEX option.

### 3.5.40 VIEWING CONE Option

The VIEWCONE option enables the generation of showers within the viewing cone of e.g. a Cherenkov telescope. Around the fixed incidence angle defined by THETPR(1) and PHIPR(1) (page 77) a (hollow) cone is defined with its tip pointing towards the detector. The inner and outer limiting angles of this cone are defined by the keyword VIEWCONE (page 78). The zenith angular dependence of the chosen detector geometry (see Sect. 3.5.41 below) is maintained for flat horizontal resp. spherical detectors, while the VOLUMECORR option is not supported. Showers originating from those portions of the cone which exceed the allowed range of CORSIKA are not simulated, rather they are skipped and a new angle is selected at random out of the range of the cone.

To obtain this program version the VIEWCONE option has to be selected when extracting the FORTRAN code from the source file.

### 3.5.41 VOLUME DETector and Vertical String Geometry Options

With this options it is possible to select at random the zenith angle in a manner which respects the geometrical acceptance of the detector.

The **default** primary intensity distribution  $I$  goes with the zenith angle  $\theta$  like

$$I \propto \sin\theta \cdot \cos\theta$$

The *sin* term respects the solid angle element of the sky, while the *cos* term takes the geometrical efficiency of a flat horizontal detector into account<sup>65</sup>. This allows to use each shower several times with the shower axis intersecting the detector array with equal distribution in  $x$  and  $y$  at random. The area to be covered by randomly scattering the shower axis is independent from the zenith angle and extends horizontally.

Using the VOLUMEDET option, the primary intensity distribution  $I$  goes with the zenith angle like

$$I \propto \sin\theta$$

respecting only the solid angle elements of the sky. This is appropriate for detectors of approximated spherical geometry, e.g. atmospheric Cherenkov telescopes. To use a shower several times you might scatter it on an area, which has fixed extensions in a plane perpendicular to the shower axis.

Using the VOLUMECORR option, the primary intensity distribution is a more complicated function of the zenith angle, which respects the geometry of a long vertical string detector (e.g. AMANDA [70] or ICECUBE [58] experiment and other neutrino telescopes) with a ratio of  $l/d = \text{length/diameter}$  of the sensitive volume. The functional form of the zenith angle distribution becomes

$$I \propto (d/2)^2 \cdot \pi \cdot \sin\theta \cdot (\cos\theta + 4/\pi \cdot l/d \cdot \sin\theta) \quad .$$

The  $l/d$  ratio (defining the DETector ConFiGuration) has to be read in using the keyword DETCFG (page 107).

The VOLUMECORR option cannot be combined with the VIEWCONE option.

To obtain these program versions the VOLUMEDET resp. VOLUMECORR option has to be selected when extracting the FORTRAN code from the source file.

## 3.6 Combination of Options

In principle most options may be combined. Do not combine any FLUKA version with the present version DPMJET III because of severe clashes caused by duplicate names of COMMONS and subroutines. DPMJET, EPOS, NEXUS, QGSJET, and SIBYLL are tested only with UNIX work stations. You can not select more than one high-energy and one low-energy hadronic interaction model at a time. In principle a combination of NUPRIM with all hadronic interaction models is possible (page 162). The CONEX option can only be used with the high-energy interaction models EPOS, QGSJET, QGSII, or

---

<sup>65</sup>Selecting in the CURVED version the zenith angle  $\theta$  at random one should keep in mind that for the default version the probability vanishes at  $\theta = 90^\circ$ . If the zenith angle range is specified as  $\theta_1 < \theta < \theta_2$  with  $\theta_1 < 90^\circ$  and  $90^\circ < \theta_2$  the zenith angle is selected at random from  $\text{MIN}(\theta_1, 180^\circ - \theta_2) < \theta < \text{MAX}(\theta_1, 180^\circ - \theta_2)$ .

SIBYLL in combination with any low-energy interaction model. The CHARM option makes sense only with the interaction models DPMJET, QGSJET01d, or SIBYLL (with EPOS in preparation), as only these interaction models are producing charmed particles explicitly.

The EHISTORY option cannot be combined with the PARALLEL options, as identical positions in several arrays are used in these two options for completely different purposes. The combination of PRESHOWER with NUPRIM makes no sense as your primary may be either a high-energy gamma ray or a neutrino. Neither PRESHOWER nor NUPRIM may be combined with STACKIN. The COASTUSERLIB, COREAS, INCLINED, and ROOTOUT options cannot be combined with the COMPACT option, nor with histogram producing options ANAHIST, AUGERHIST, INTTEST or MUONHIST.

The combination of INTTEST with PLOTSH or PLOTSH2 is not reasonable as in the INTTEST option the development of showers is suppressed. Also a combination of INTTEST with ANAHIST, AUGERHIST, COREAS, CURVED, EHISTORY, INCLINED, MUONHIST, PRESHOWER, STACKIN, or THIN gives no sense, as you are simulating just the first interaction without development of a complete shower. A combination with UPWARD makes no sense as in INTTEST anyway all upward going particles are respected. The combination of NUPRIM with INTTEST is not supported. The combination of FLUKA with INTTEST is not foreseen.

IACT is only possible with CERENKOV; IACT and ATMEXT have been tested only with UNIX computers. VOLUMEDET and VOLUMECORR exclude each other, as you may use only one detector geometry at a time. Similarly INTCLONG and NOCLONG are excluding each other.

As in the CEFFIC option the atmospheric absorption is treated only in a manner suited for planar atmospheres, you should not combine the CURVED option with CEFFIC. The combination of VOLUMECORR with VIEWCONE is not supported. The AUGERHIST option cannot be combined with the IACT, CURVED, INTTEST, PLOTSH, or PLOTSH2 options.

DYNSTACK is not tested with any parallel execution model, AUGERHIST and COAST. It is not possible to use the ICECUBE1 extensions in combination with the DYNSTACK option.

The NRREXT option cannot be combined with options generating histograms, the combination with PARALLEL and PARALLELIB options is not tested.

Using the *./coconut* script file (see Sect. 2.3.2 page 18) conflicting options are indicated and will be avoided automatically.

## 4 Steering of the Simulation

The simulation of air showers is steered by commands (keywords) that have to be given on unit MONIIN (MONItor INput) in the card image format. A command consists of a keyword usually up to 6 characters long (left shifted, upper or lower case characters) and one or more arguments in the form:

```
KEYWRD   arg1   arg2   arg3   ...   arg$n$   comments
```

The keyword and the arguments must be separated at minimum by one blank. The last argument may be followed by comments up to column 512. Additional comments may be given on separate lines with the first 6 characters (the keyword) kept blank, with the first character a `c` or a `C` followed by a blank, or with the first character a `*`. In the IACT option (page 45) the lines starting with 'IACT' are treated additionally as comment lines. Internally all characters including the keywords are converted to upper case characters, except the characters following the keywords EPOPAR, DATDIR, DIRECT, HISTDS, HOST, IACT, INFILE, NEXPAR, TELFIL, and USER. If you want to specify one of these character arguments by a blank, you should include the blank within apostrophes or quotation marks. The sequence of steering keywords is arbitrary. The valid keywords, the internal argument names, their nature (A = character, F = floating, I = integer, or L = logical), their default settings, their descriptions, and their limitations are listed in the following.

As CORSIKA is primarily designed to simulate EAS in the energy range  $10^{11}$  eV to some  $10^{20}$  eV the code contains parameterizations and approximations which are valid only for a limited range of some arguments. Leaving the recommended range might cause incorrect results or even end the execution of the program with an error message.

### 4.1 Run Number

```
RUNNR     NRRUN
```

Format = (A5, I), Default = 1

NRRUN : Run number of this simulation. This number is used to form part of the name of the various output files.

Limit is:  $0 \leq \text{NRRUN} \leq 999999$

(rsp.  $0 \leq \text{NRRUN} \leq 999999999$  in case of the NRREXT option, see Sect. 3.5.24 page 61).

### 4.2 First Event Number

```
EVTNR     SHOWNO
```

Format = (A5, I), Default = 1

SHOWNO : Event number of first shower. The second shower will get number

SHOWNO+1 and so on.

Limit is:  $1 \leq \text{SHOWNO} \leq 999999$

### 4.3 Random Number Generator Initialization

SEED      ISEED(i,k), i=1... 3

Format = (A4, 3I), Defaults = k, 0, 0

ISEED(1, k) : Contains the seed of the random number sequence k.

ISEED(2..3, k) : Contain the number of calls  $N_{in}$  to the generator that are performed for initialization such that  $N_{in} = \text{ISEED}(2, k) + 10^9 \cdot \text{ISEED}(3, k)$ .

At present at most  $k = 7$  sequences are used: Sequence 1 for the hadron shower, 2 for the EGS4 part, 3 for the simulation of Cherenkov photons (only for CERENKOV option), 4 for the random offset of Cherenkov telescope systems with respect of their nominal positions (only for IACT option) resp. for the offset scattering of the core relative to the Auger detector (AUGERHIT option), 5 for the HERWIG routines in the NUPRIM option, 6 for the PARALLEL option, and 7 for the CONEX option. Their activation follows the sequence of occurrence of the keyword lines. The sequences 9 ... 14 are activated by the keyword MSEED (see Sect. 4.59 page 98).

**At minimum 2 seeds should be activated.** If not sufficient seeds are activated, the default values are taken.

The use of  $\text{ISEED}(2, k) > 0$  and especially of  $\text{ISEED}(3, k) > 0$  should be avoided as pre-setting the random number generator by billions of calls needs considerable computing time. To get different random sequences it is sufficient to modify  $\text{ISEED}(1, k)$ .

When the `eventio` and other separate functions are enabled in the IACT option, an external random generator may be used.

Limit (to get independent sequences of random numbers) is:  $1 \leq \text{ISEED}(1, k) \leq 900\,000\,000$

### 4.4 Number of Showers

NSHOW      NSHOW

Format = (A5, I), Default = 10

NSHOW : Number of showers to be generated in a run.

Limit is:  $\text{NSHOW} \geq 1$

### 4.5 Primary Particle Definition

PRMPAR      PRMPAR(0)

Format = (A6, I), Default = 14

PRMPAR(0) : Particle type of the primary particle. See Table 4 (page 129) for the particle codes.

Limits are:  $1 \leq \text{PRMPAR}(0) \leq 5656$ <sup>66</sup>. Vector mesons, resonances, and charmed particles are excluded because of their short life time. Primary neutrinos can only be used in the NUPRIM option (page 61). Unstable nuclei with mass number 5 or 8 may be used only with SIBYLL or QGSJET, for other models the cross-sections are not defined.

This keyword is not available in the STACKIN option.

## 4.6 Energy Range

ERANGE    LLIMIT    ULIMIT

Format = (A6, 2F), Defaults = 1.E4, 1.E4

LLIMIT : Lower limit and

ULIMIT : Upper limit of the primary particle energy range (in GeV). The primary energy is selected at random out of this interval. If LLIMIT = ULIMIT, the primary energy is fixed at this value.

The energies are total energies and include the particle rest mass.

Limits are: For primary nuclei LLIMIT > HILOW (by default 80 GeV/nucleon for nuclei (30 GeV/nucleon for EPOS), see page 92); below this limit a simple superposition model is used.

For primary photons and electrons ULIMIT  $\leq$  1.E11 GeV (but keep in mind that no LPM-effect is included in NKG!). For primary hadrons and nuclei no upper limit is recommended, but **the user should take care not to over-stretch the selected hadronic interaction model**. See also Ref. [29]. It is recommended for HDPM: ULIMIT  $\leq$   $1 \cdot 10^8$  GeV and for VENUS: ULIMIT  $\leq$   $2 \cdot 10^7$  GeV. For the DPMJET option above ULIMIT >  $10^{10}$  GeV this model becomes uncertain.

This keyword is not available in the STACKIN option.

## 4.7 Slope of Energy Spectrum

ESLOPE    PSLOPE

Format = (A6, F), Default = 0.

PSLOPE : Exponent  $\gamma$  of differential primary energy spectrum. The primary energy is taken at random from an exponential energy spectrum of the form  $dN/dE_0 \propto E_0^\gamma$ . PSLOPE has no meaning in case of fixed primary energy. The energies are total energies and include the particle rest mass.

---

<sup>66</sup>In the combination EPOS+CONEX, the upper limit can be extended to 25099 if the option is selected in *./coconut*



This keyword is not available in the STACKIN option.

## 4.8 Zenith Angle Definition

THETAP      THETPR(1)      THETPR(2)

Format = (A6, 2F), Defaults = 0., 0.

THETPR(1) : Low edge of zenith angle range of primary particle (in °).

THETPR(2) : High edge of zenith angle range of primary particle (in °).

The zenith angle is selected at random out of this interval in a manner which respects equal particle fluxes from all solid angle elements of the sky and a registration by a horizontal flat detector arrangement<sup>67</sup>. THETPR is the angle of incidence at a horizontal detector. THETPR( $i$ ) = 0. is vertical. If THETPR(1) = THETPR(2), the zenith angle is fixed at this value.

Limits<sup>68</sup> are:  $0. \leq \text{THETPR}(i) \leq 70$ .

## 4.9 Azimuth Angle Definition

PHIP      PHIPR(1)      PHIPR(2)

Format = (A4, 2F), Defaults = 0., 0.

PHIPR(1) : Low edge of azimuth angle range of primary particle (in °).

PHIPR(2) : High edge of azimuth angle range of primary particle (in °).

The azimuth angle is selected at random out of this interval.

If PHIPR(1) = PHIPR(2), the azimuth angle is fixed at this value. For  $\phi = 0^\circ$  the shower axis points to magnetic North, for  $\phi = 90^\circ$  it points to West, see Fig. 1 (page 126).

Limits are:  $-360. \leq \text{PHIPR}(i) \leq 360$ <sup>69</sup>.

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<sup>67</sup>In the case you use a volume detector (sphere) or a vertical long string detector instead of a flat horizontal detector, you should respect this by selecting the preprocessor option VOLUMEDET (for sphere) or VOLUMECORR (for vertical long string) to get the angular dependence of the shower intensity as observed with such detectors (see Sect. 3.5.41 page 72).

<sup>68</sup>The zenith angle limitation is recommended for the standard CORSIKA version because of some approximations made in subroutine *NKG*. At  $\theta > 70^\circ$  also the curvature of the Earth's surface must be taken into account. For large zenith angles you should use the CURVED option (page 54).

For the CURVED option the limit is  $0. \leq \text{THETPR}(i) < 90.$ ; for the CURVED option combined with the CERENKOV option the limit is  $0. \leq \text{THETPR}(i) \leq 88$ .

For the UPWARD option with upward going primary the limits are  $110. \leq \text{THETPR}(i) \leq 180$ .

For the CURVED option combined with the UPWARD option THETPR( $i$ ) has to be chosen in a manner excluding the range  $90. < \text{THETPR}(i) < 90. + \delta$ . The angle  $\delta < 90^\circ$  is spanned between the upward going shower axis and the horizontal detector plane above the shower starting point. Further details are given in [66, 71]. The keyword IMPACT (page 79) for skimming horizontal showers overrides THETPR( $i$ ).

<sup>69</sup>In the output file the corresponding  $\phi$  of each shower is given in the range  $[-\pi, \pi]$  radian.

## 4.10 Viewing Cone Specifications

VIEWCONE    VUECON(1)    VUECON(2)

Format = (A8, 2F), Defaults = 0., 0.

VUECON(1) : Inner limiting angle of viewing cone (in °).

VUECON(2) : Outer limiting angle of viewing cone (in °).

The VIEWCONE option (see Sect. 3.5.40 page 71) selects the direction of primaries in a circular cone around the fixed primary direction THETPR(1) and PHIPR(1) (page 77) with the inner opening VUECON(1) and the outer opening VUECON(2). The zenith angular dependence of the selected detector geometry is maintained for flat horizontal resp. spherical detectors (see Sect. 3.5.41 page 72).

Limits:  $0. \leq \text{VUECON}(1) \leq \text{VUECON}(2) < 90$ . The generation of showers with angles beyond the range of the program validity is skipped.

This keyword is only available in the VIEWCONE option.

## 4.11 Starting Gammage

FIXCHI    THICK0

Format = (A6, F), Default = 0.

THICK0 : The vertical starting altitude (in g/cm<sup>2</sup> mass overburden) of the primary particle is set for all showers. This choice is not effective if the height of the first interaction is set by FIXHEI > 0. (see Sect. 4.13 below). With this keyword the development of sub-showers starting at the chosen altitude within the atmosphere may be followed. It will define the starting (0) point of the longitudinal profile with the SLANT option <sup>70</sup>. including when the FIXHEI option is selected. The starting altitude must be above the lowest observation level.

Limit is:  $0. \leq \text{THICK0}$ .

This keyword is not available in the COASTUSERLIB or STACKIN option.

## 4.12 Starting Point of Arrival Timing

TSTART    TMARGIN

Format = (A6, L), Default = F

TMARGIN : Flag indicating the starting point of the arrival time scale. If TMARGIN = .false., the first interaction starts the clock. If TMARGIN = .true., the entrance into the atmosphere (resp. THICK0, see above) is taken for starting the internal clock. Additionally, the ionization energy loss, the deflection within the Earth's magnetic field, and the generation of Cherenkov photons is enabled for charged hadronic or muonic primaries on

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<sup>70</sup>Note that THICK0 is always defined as vertical depth even with the SLANT option.

their path between entering the atmosphere and the first interaction, which otherwise is disabled in the standard version (for the CERENKOV, CURVED, and SLANT options see page 42, 54, resp. 66). For TMARGIN = .true. the height of the first interaction is written negative to element 7 of the event header block.

This keyword is not available in the CONEX, CURVED, SLANT, or STACKIN options, but is set and used by default as .true.; TMARGIN is set .true. by default in the CERENKOV and IACT options, but it may be overwritten using the keyword.

### 4.13 First Interaction Definition

FIXHEI      FIXHEI      N1STTR

Format = (A6, F, I), Defaults = 0., 0

FIXHEI : Fixes the height (in cm) of the first interaction of hadronic primaries (resp. the starting altitude for em-particles) for all showers in a run. If FIXHEI = 0., the height of the first interaction is varied at random according to the appropriate mean free path. In case of unstable hadronic primaries and fixed height the first interaction will not be a decay. The fixed height must be above the lowest observation level. If FIXHEI > 0. is set, the starting grammage FIXCHI of the primary is not effective (see Sect. 4.11 above) for the propagation but will define the starting point of the longitudinal profile unless the STACKIN and SLANT options are selected too. In that case, and only in that case, FIXHEI also define the starting point of the shower profile in slant depth.

In the STACKIN option FIXHEI is needed to specify the altitude of the first, externally treated interaction.

In the CURVED option the keyword FIXHEI cannot be used for em-primary particles.

N1STTR : Fixes the target of the first interaction: 1 = Nitrogen, 2 = Oxygen, 3 = Argon, *else* = random selection according to the atmospheric abundances. This option is only applicable for high-energy hadronic primaries, i.e. primaries with an energy per nucleon of  $E_{lab} \geq \text{HILOW}$  (see page 92). Also in the NUPRIM option it may be used.

With the STACKIN option one should select N1STTR = 0 as in this case it acts on that secondary particle which is treated first.

In case of unstable hadronic primaries and predetermined target the first interaction will not be a decay.

Limits are:  $0. \leq \text{FIXHEI} < \text{border of atmosphere}$  (at 112.8E5 cm for atmospheric models  $1 \leq \text{MODATM} \leq 9$  or  $\text{MODATM} \geq 17$ ).

### 4.14 Skimming Incidence

IMPACT      HIMPACT(1)      HIMPACT(2)

Format = (A6, 2F), Default = 0., 0.

HIMPACT(1) : Lower value (in cm) for minimum altitude of horizontal shower axis.

HIMPACT(2) : Upper value (in cm) for minimum altitude of horizontal shower axis. The actual minimum altitude is selected at random out of this interval with uniform distribution. Zenith angles given by the keyword THETAP (page 77) are overridden by a calculation from the actual minimum altitude. See UPWARD option page 70.

Limits are:  $OBSLEV(1) \leq HIMPACT(i) \leq \min(FIXHEI, FIXCHI, \text{border of atmosphere [at } 112.8E5 \text{ cm for atmospheric models } 1 \leq MODATM \leq 9 \text{ or } MODATM \geq 17])$ .

This keyword is only available in the combination of the CURVED option with the UPWARD option.

## 4.15 Stack Input File Name

INFILE      FILINP

Format = (A6, A132), Defaults = ' '

FILINP : File name to define the name and directory of the input file containing the parameters of secondary particles (see Sect. 3.5.34 page 66). Lower case characters of FILINP are not converted to capitals. Please keep in mind that in FORTRAN an automatic expansion of UNIX names like '\$HOME' is not possible, rather you should give the fully expanded name of the directory and file.

Limit is: FILINP must not begin with a ~ (tilde) character.

This keyword is only available in the STACKIN option.

## 4.16 Stack Output File Name

OUTFILE      FILOUT

Format = (A7, A132), Defaults = ' '

FILOUT : File name to define the name and directory of the output file which will contain the parameters of the secondary particles produced in the first interaction. These particles are written only if the file name is defined. If more than one shower is simulated in one run, only the secondary particles of the first interaction coming from the first shower are written, those from the second and further showers are skipped.

These particles may be read in in a later CORSIKA-run established with the STACKIN option (see Sect. 3.5.34 page 66). Lower case characters of FILINP are not converted to capitals. Please keep in mind that in FORTRAN an automatic expansion of UNIX names like '\$HOME' is not possible, rather you should give the fully expanded name of the directory and file.

Limit is: FILOUT must not begin with a ~ (tilde) character.

This keyword is not available in the CONEX or STACKIN options.

## 4.17 Atmospheric Model Selection

ATMOD      MODATM

Format = (A5, I), Default = 1

MODATM : Gives the number of the atmospheric parameterization.

MODATM = 0: Atmosphere as read in by keywords ATMA, ATMB, ATMC, and ATM-LAY (uppermost layer unchanged).

MODATM = 1: U.S. standard atmosphere as parameterized by Linsley.

MODATM = 2: AT115 Central European atmosphere for Jan. 15, 1993.

MODATM = 3: AT223 Central European atmosphere for Feb. 23, 1993.

MODATM = 4: AT511 Central European atmosphere for May 11, 1993.

MODATM = 5: AT616 Central European atmosphere for June 16, 1993.

MODATM = 6: AT822 Central European atmosphere for Aug. 22, 1993.

MODATM = 7: AT1014 Central European atmosphere for Oct. 14, 1993.

MODATM = 8: AT1224 Central European atmosphere for Dec. 24, 1993.

MODATM = 9: Atmosphere as read in by keywords ATMA, ATMB, ATMC. (Layers as in MODATM = 1 . . . 8.)

MODATM = 10: Atmosphere as read in by keywords ATMA, ATMB, ATMC, and ATM-LAY (uppermost layer also read in).

MODATM = 11: South pole atmosphere for March 31, 1997 (MSIS-90-E).

MODATM = 12: South pole atmosphere for July 01, 1997 (MSIS-90-E).

MODATM = 13: South pole atmosphere for Oct. 01, 1997 (MSIS-90-E).

MODATM = 14: South pole atmosphere for Dec. 31, 1997 (MSIS-90-E).

MODATM = 15: South pole atmosphere for January after Lipari.

MODATM = 16: South pole atmosphere for August after Lipari.

MODATM = 17: U.S. standard atmosphere as parameterized by Keilhauer.

MODATM = 18: Malargüe GDAS model for January after Will/Keilhauer.

MODATM = 19: Malargüe GDAS model for February after Will/Keilhauer.

MODATM = 20: Malargüe GDAS model for March after Will/Keilhauer.

MODATM = 21: Malargüe GDAS model for April after Will/Keilhauer.

MODATM = 22: Malargüe GDAS model for May after Will/Keilhauer.

MODATM = 23: Malargüe GDAS model for June after Will/Keilhauer.

MODATM = 24: Malargüe GDAS model for July after Will/Keilhauer.

MODATM = 25: Malargüe GDAS model for August after Will/Keilhauer.

MODATM = 26: Malargüe GDAS model for September after Will/Keilhauer.

MODATM = 27: Malargüe GDAS model for October after Will/Keilhauer.

MODATM = 28: Malargüe GDAS model for November after Will/Keilhauer.

MODATM = 29: Malargüe GDAS model for December after Will/Keilhauer.

MODATM = 30: South Pole for IceTop for January after S. de Ridder.

MODATM = 31: South Pole for IceTop for February after S. de Ridder.

MODATM = 32: South Pole for IceTop for March after S. de Ridder.

MODATM = 33: South Pole for IceTop for April after S. de Ridder.

MODATM = 34: South Pole for IceTop for May after S. de Ridder.

MODATM = 35: South Pole for IceTop for June after S. de Ridder.  
 MODATM = 36: South Pole for IceTop for July after S. de Ridder.  
 MODATM = 37: South Pole for IceTop for August after S. de Ridder.  
 MODATM = 38: South Pole for IceTop for September after S. de Ridder.  
 MODATM = 39: South Pole for IceTop for October after S. de Ridder.  
 MODATM = 40: South Pole for IceTop for November after S. de Ridder.  
 MODATM = 41: South Pole for IceTop for December after S. de Ridder.  
 The various atmospheric models are described in Appendix F (page 173 ff.).  
 Limits are:  $0 \leq \text{MODATM} \leq 41$  .

#### 4.18 Atmospheric Parameters A(i)

ATMA    AATM1    AATM2    AATM3    AATM4    (AATM5)

Format = (A4, 4F), Defaults = 0., 0., 0., 0. (for ATM0D 0)

Format = (A4, 5F), Defaults = 0., 0., 0., 0., 0. (for ATM0D 10)

AATMi : A-parameters for 4 layers of atmospheric model # 0 (or for 5 layers of atmospheric model # 10). For the 5<sup>th</sup> layer a linear decrease is assumed, which in case of ATM0D = 0 uses the same parameters as the U.S. standard atmosphere. To be used with ATM0D = 0 or 10 .

Limit is:  $0. < \text{AATM5}$

#### 4.19 Atmospheric Parameters B(i)

ATMB    BATM1    BATM2    BATM3    BATM4

Format = (A4, 4F), Defaults = 0., 0., 0., 0.

BATMi : B-parameters for 4 layers of atmospheric model # 0. For the 5<sup>th</sup> layer a linear decrease is assumed with the same parameters as for the U.S. standard atmosphere. To be used with ATM0D = 0 or 10.

Limits are:  $\text{BATM}_i \neq 0$ .

#### 4.20 Atmospheric Parameters C(i)

ATMC    CATM1    CATM2    CATM3    CATM4    (CATM5)

Format = (A4, 4F), Defaults = 0., 0., 0., 0. (for ATM0D 0)

Format = (A4, 5F), Defaults = 0., 0., 0., 0., 0. (for ATM0D 10)

CATMi : C-parameters for 4 layers of atmospheric model # 0 (or for 5 layers of atmospheric model # 10). For the 5<sup>th</sup> layer a linear decrease is assumed, which in case of ATM0D = 0 uses the same parameters as the U.S. standard atmosphere. To be used

with  $ATMOD = 0$  or  $10$ .  
Limits are:  $CATM_i > 0$ .

## 4.21 Atmospheric Layer Boundaries

ATMLAY    HLAY2    HLAY3    HLAY4    HLAY5

Format = (A6, 4F), Defaults = 4.D5, 10.D5, 40.D5, 100.D5

HLAY $i$  : Layer lower boundaries (in cm) for the layers of atmospheric model # 0 and # 10. A value of 0. is adopted for the HLAY1. If not specified, the default values of  $MODATM = 1$  are used for  $MODATM = 0$  and  $10$ . For other models ( $MODATM \neq 0$  and  $\neq 10$ ), the default values correspond with the selected model  $MODATM$ . Should only be used with  $ATMOD = 0$  or  $10$ .

Limits are:  $0. < HLAY_i$ .

## 4.22 External Tabulated Atmosphere

ATMOSPHERE    IATMOX    FREFRX

Format = (A10, I, L), Defaults = 0, F

IATMOX : Use MODTRAN [45] atmospheric model  $IATMOX = i$  (in terms of density and refractive index) instead of CORSIKA built-in model. This requires a file named *atm-profi.dat*. MODTRAN model atmospheres supplied with the 'bernlrohr' package include tropical ( $i = 1$ ), mid-latitude summer (2), mid-latitude winter (3), sub-arctic summer (4), sub-arctic winter (5), and U.S. standard atmosphere 1976 (6). Additionally for the MAGIC Cherenkov telescope on La Palma the summer (7) and winter (8) atmospheres[72] are supplied. The atmosphere (9) gives the winter atmosphere at the South pole. User supplied models are possible ( $i > 9$ ).

FREFRX : If .true., the atmospheric refraction for Cherenkov photons is taken into account (for plane-parallel atmosphere); if .false., refraction is ignored. The value of this second argument is ignored if the CERENKOV option is not selected.

This keyword is only available in the ATMEXT option and needs linking with the (compiled) *atmo.c* routines of the 'bernlrohr' package.

## 4.23 Atmospheric profile file

ATMFILE    FILEATM

Format = (A7, A1024), Defaults = ' '

FILEATM: File name to define the name and directory where to read ATMLAY, ATMA,

ATMB and ATMC parameters, as well as a tabulated altitude profile of the atmospheric refractive index.

The file FILEATM needs to be produced with the `gdastool` script included with CORSIKA (in `src/utills/`), or to have the same format. See Appendix G page 185 for how to use `gdastool`. This allows to use atmospheric profiles for density and refractive index as compiled from the GDAS (Global Data Assimilation System) database. This database contains atmospheric parameters as used in e.g. weather forecasting models. It features atmospheric parameters at 24 altitudes in the atmosphere, at a 1 x 1 degree grid on Earth, at a 3-hour time resolution.

Limit is: FILEATM must not begin with a `~` (tilde) character and should contain the full path of the file if it is not located where CORSIKA is run (DATDIR not used).

This keyword cannot be used together with a manual specification of ATMLAY, ATMA, ATMB and/or ATMC.

## 4.24 Earth's Magnetic Field

MAGNET BX BZ

Format = (A6, 2F), Defaults = 20.40, 43.23

BX : Is the horizontal component of the Earth's magnetic field (in  $\mu\text{T}$ ) to the x-direction of the coordinate system (North) (see Fig. 1 page 126)

BZ : Is the vertical component of the Earth's magnetic field (in  $\mu\text{T}$ ) downwards.

The default values represent the magnetic field for the Karlsruhe location. The values of other locations may be obtained from the program *Geomag* which is available on-line in the world wide web [73]. The value H of *Geomag* corresponds with our BX, the value Z with our BZ. For the orientation of the CORSIKA coordinate system see also Fig. 1 page 126.

Limits are: BX, BZ  $\neq$  0. .

## 4.25 Experiment Coordinates for Pre-Showering and CONEX

GCOORD GLONG GLATI GRFYEAR IPREPR IPRSTP

Format = (A6, 3F, 2I), Defaults = -69.585, -35.463, 2013., 1, 0

GLONG : Gives the geographical longitude (in  $^\circ$ , West length is negative) of the experiment.

GLATI : Gives the geographical latitude (in  $^\circ$ , South latitude is negative) of the experiment.

GRFYEAR : Gives the year of the experiment (the magnetic field is varying with time). These coordinates are used to calculate the magnetic dipole field of the Earth's globe above the atmosphere of the experiment's position in the case of a pre-shower induced by ultra-high energetic primary photons. The default values give the position coordinates of the southern Pierre Auger Observatory at Malargüe (Argentina) for the year 2003.



IPREPR : Print indicator<sup>71</sup>: IPREPR  $\leq 0$  disables pre-shower printing; IPREPR = 1 prints details of pre-shower in case of MAXPRT (page 107) or DEBUG (page 119); IPREPR  $\geq 2$  always prints details of pre-shower.

IPRSTP : If IPRSTP  $\neq 0$  events without pre-showering are skipped.

Limits are:  $-180. \leq \text{GLONG} \leq 180.$ ;  $-90. \leq \text{GLATI} \leq 90.$ ;  $1965. \leq \text{GRFYEAR} \leq 2015.$   
This keyword is only available in the CONEX or PRESHOWER options.

## 4.26 DPMJET Selection Flag

DPMJET    FDPMJT    LEVLDB

Format = (A6, L, I), Defaults = T, 0

FDPMJT : If .true., the DPMJET routines are used to treat the high-energy hadronic interactions. If .false., the HDPM routines are used to treat the high-energy hadronic interactions.

LEVLDB : Gives amount of debug output for the DPMJET code in case of DEBUG. With increasing value up to 8 the debug output becomes more and more detailed. This output cannot be redirected and always appears on unit 6.

Limits are:  $0 \leq \text{LEVLDB} \leq 8$ .

This keyword is only available in the DPMJET option.

## 4.27 DPJSIG Selection Flag

DPJSIG    FDPJSG

Format = (A6, L), Default = T

FDPJSG : If .true., the DPMJET high-energy hadronic cross-sections are used. If .false., the default cross-sections as described in Ref. [3] are used.

This keyword is only available in the DPMJET option.

## 4.28 EPOS Selection Flag

EPOS    FNEXUS    ISHON

Format = (A5, L, I), Defaults = T, 0

FNEXUS : If .true., the EPOS routines are used to treat the high-energy hadronic interactions. If .false., the HDPM routines are used to treat the high-energy hadronic interactions.

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<sup>71</sup>The print statements within the PRESHOWER C-routines write only to the standard output. A redirecting of the standard output to a logical unit differing from 6 (see keyword OUTPUT page 109) usually will not affect the output of the PRESHOWER C-routines.

ISH0N : Determines amount of debug output for the EPOS routines. With increasing number  $ISH0N > 0$  the output becomes more and more detailed. This output appears on the unit MDEBUG.

For more information look into the EPOS documentation. Additional debugging is effective by setting print parameters using EPOPAR *print* . . . . This debug output is written to the *ifch* file (see Table 1 page 26).

Limits are:  $0 \leq ISH0N \leq 9$  .

This keyword is only available in the EPOS option.

## 4.29 EPOS Parameters

EPOPAR      *parcha*

Format = (A6, A74), Defaults = ' '

*parcha* : Command line to be read by subroutine *aread* of program block *epos-bas-lhc.f* .

The possible command lines are described in the EPOS documentation. Use lower case characters. Lower case characters of *parcha* are not converted to capitals. Do not use the commands *application* . . . , *set nevent* . . . , *run*, or *stop* within your input parameters, these will cause unpredictable results or crashes. Only *fname* names might be changed by standard users<sup>72</sup>.

This keyword is only available in the EPOS option.

## 4.30 EPOSIG Selection Flag

EPOSIG      FNEXSG

Format = (A6, L), Default = T

FNEXSG : If .true., the EPOS high-energy hadronic cross-sections are used. If .false., the

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<sup>72</sup>A typical EPOPAR input looks like:

```
EPOPAR input ~ corsika-78010/epos/epos.param
EPOPAR fname pathnx ~ corsika-78010/epos/
EPOPAR fname inics ~ corsika-78010/epos/epos.inics
EPOPAR fname iniev ~ corsika-78010/epos/epos.iniev
EPOPAR fname inirj ~ corsika-78010/epos/epos.inirj
EPOPAR fname initl ~ corsika-78010/epos/epos.initl
EPOPAR fname check ~ corsika-78010/epos/epos.check
EPOPAR fname hpf ~ corsika-78010/epos/urqmd34/tables.dat
EPOPAR fname histo ~ corsika-78010/epos/epos.histo ! for interaction test only
EPOPAR fname data ~ corsika-78010/epos/epos.data ! for debugging only
EPOPAR fname copy ~ corsika-78010/epos/epos.copy ! for debugging only
EPOPAR fname log ~ corsika-78010/epos/epos.log ! for debugging only
EPOPAR printcheck screen ! for debugging only.
```

default cross-sections as described in Ref. [3] are used.  
 This keyword is only available in the EPOS option.

### 4.31 NEXUS Selection Flag

NEXUS      FNEXUS      ISH0N

Format = (A5, L, I), Defaults = T, 0

FNEXUS : If *.true.*, the NEXUS routines are used to treat the high-energy hadronic interactions. If *.false.*, the HDPM routines are used to treat the high-energy hadronic interactions.

ISH0N : Determines amount of debug output for the NEXUS routines. With increasing number  $ISH0N > 0$  the output becomes more and more detailed. This output appears on the unit MDEBUG.

For more information look into the NEXUS documentation. Additional debugging is effective by setting print parameters using NEXPAR *print . . .*. This debug output is written to the *ifch* file (see Table 1 page 26).

Limits are:  $0 \leq ISH0N \leq 9$ .

This keyword is only available in the NEXUS option.

### 4.32 NEXUS Parameters

NEXPAR      *parcha*

Format = (A6, A74), Defaults = ' '

*parcha* : Command line to be read by subroutine *aread* of program block *nexus-bas.f*. The possible command lines are described in the NEXUS documentation. Use lower case characters. Lower case characters of *parcha* are not converted to capitals. Do not use the commands *application . . .*, *set nevent . . .*, *run*, or *stop* within your input parameters, these will cause unpredictable results or crashes. Only *nexus.inixx* names might be changed by standard users<sup>73</sup>.

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<sup>73</sup>A typical NEXPAR input looks like:

```
NEXPAR fname inics ~ corsika-78010/nexus/nexus.inics
NEXPAR fname iniev ~ corsika-78010/nexus/nexus.iniev
NEXPAR fname inirj ~ corsika-78010/nexus/nexus.inirj
NEXPAR fname initl ~ corsika-78010/nexus/nexus.initl
NEXPAR fname check ~ corsika-78010/nexus/nexus.check
NEXPAR fname histo ~ corsika-78010/nexus/nexus.histo ! for interaction test only
NEXPAR fname data ~ corsika-78010/nexus/nexus.data ! for debugging only
NEXPAR fname copy ~ corsika-78010/nexus/nexus.copy ! for debugging only
NEXPAR fname log ~ corsika-78010/nexus/nexus.log ! for debugging only
NEXPAR printcheck screen ! for debugging only.
```

This keyword is only available in the NEXUS option.

### 4.33 NEXSIG Selection Flag

NEXSIG      FNEXSG

Format = (A6, L), Default = T

FNEXSG : If `.true.`, the NEXUS high-energy hadronic cross-sections are used. If `.false.`, the default cross-sections as described in Ref. [3] are used.

This keyword is only available in the NEXUS option.

### 4.34 QGSJET Selection Flag

QGSJET      FQGS      LEVLDQ

Format = (A6, L, I), Defaults = T, 0

FQGS : If `.true.`, the `qgsjet-II-04` (rsp. `QGSJET01d`) routines are used to treat the high-energy hadronic interactions. If `.false.`, the HDPM routines are used to treat the high-energy hadronic interactions.

LEVLDQ : Gives amount of debug output for the `qgsjet-II-04` (rsp. `QGSJET01d`) code in case of DEBUG. With increasing value up to 4 the debug output becomes more and more detailed. This output cannot be redirected and always appears on unit 6.

Limits are:  $0 \leq \text{LEVLDQ} \leq 4$ .

This keyword is only available in the QGSJET option.

### 4.35 QGSSIG Selection Flag

QGSSIG      FQGSSG

Format = (A6, L), Default = T

FQGSSG : If `.true.`, the `qgsjet-II-04` (rsp. `QGSJET01d`) high-energy hadronic cross-sections are used. If `.false.`, the default cross-sections as described in Ref. [3] are used.

This keyword is only available in the QGSJET option.

### 4.36 SIBYLL Selection Flag

SIBYLL      FSIBYL      ISDEBUG

Format = (A6, L), Default = T, 0

FSIBYL : If `.true.`, the SIBYLL 2.3d routines are used to treat the high-energy hadronic interactions. If `.false.`, the HDPM routines are used to treat the high-energy hadronic interactions.

ISDEBUG : Debug level; with increasing level the SIBYLL debug output becomes more and more detailed. This output cannot be redirected and always appears on unit 6. This debugging becomes only active, if CORSIKA debugging is activated by the DEBUG flag (Sect. 4.114 page 119).

This keyword is only available in the SIBYLL option.

### 4.37 SIBSIG Selection Flag

SIBSIG      FSIBSG

Format = (A6, L), Default = T

FSIBSG : If `.true.`, the SIBYLL high-energy hadronic cross-sections are used. If `.false.`, the default cross-sections as described in Ref. [3] are used.

This keyword is only available in the SIBYLL option.

### 4.38 SIBCHM Selection Flag

SIBCHM      FSIBCH

Format = (A6, L), Default = T

FSIBCH : If `.true.`, charm production is active in SIBYLL (default tune in the model). It may be propagated or not in CORSIKA using the CHARM (page 50) option. If `.false.`, no charm particles are produced by SIBYLL.

This keyword is only available in the SIBYLL option.

### 4.39 VENUS Selection Flag

VENUS      FVENUS      ISH00

Format = (A5, L, I), Defaults = T, 0

FVENUS : If `.true.`, the VENUS routines are used to treat the high-energy hadronic interactions. If `.false.`, the HDPM routines are used to treat the high-energy hadronic interactions.

ISH00 : Determines the amount of debug output for VENUS routines. With increasing number  $ISH00 \geq 90$  the output becomes more and more detailed. This output appears on the unit MDEBUG. For more information look into the listing of subroutine *venini*.

Limits are:  $0 \leq ISH00 \leq 98$  .

This keyword is only available in the VENUS option.

#### 4.40 VENUS Parameters

VENPAR      PARCHA(i)      PARVAL(i)

Format = (A6, A6, F), Defaults = '            ', 0.

PARCHA(*i*) : Name of VENUS parameter to be changed.

PARVAL(*i*) : New value of VENUS parameter to be changed.

A maximum of  $i = 100$  VENUS parameters may be set by the user in arbitrary sequence.

The available names and their meaning may be taken from the listing of subroutine *venini*.

The VENUS parameters should not be changed by standard users.

This keyword is only available in the VENUS option.

#### 4.41 VENSIG Selection Flag

VENSIG      FVENSG

Format = (A6, L), Default = T

FVENSG : If *.true.*, the VENUS high-energy hadronic cross-sections are used. If *.false.*, the default cross-sections as described in Ref. [3] are used.

This keyword is only available in the VENUS option.

#### 4.42 HDPM Interaction Parameters & Fragmentation

HADFLG      NFLAIN      NFLDIF      NFLPIO      NFLPIF      NFLCHE      NFRAGM

Format = (A6, 6I), Defaults = 0, 0, 0, 0, 0, 2

Steering flags of the high-energy hadronic interaction model HDPM and of the projectile nucleus fragmentation of all hadronic interaction models.

NFLAIN : The number of interactions of a projectile in a target nucleus may fluctuate (NFLAIN = 0) or is calculated as an average value (NFLAIN  $\neq$  0).

NFLDIF : No diffractive interactions are allowed in case of more than 1 interaction in the target (NFLDIF = 0) or diffractive interactions are possible (NFLDIF  $\neq$  0).

NFLPIO : The rapidity distribution of  $\pi^0$  is taken different from that of charged pions as indicated by collider data (NFLPIO = 0) or is taken as for charged pions (NFLPIO  $\neq$  0).

NFLPIF : The number of  $\pi^0$  fluctuates in the same way as the number of charged pions (NFLPIF = 0) or fluctuates independently as parameterized from collider data (NFLPIF  $\neq$  0).

NFLCHE : Charge exchange reactions for the proj. and target particles are allowed

(NFLCHE = 0) or inhibited (NFLCHE  $\neq$  0).

NFRAGM : A primary nucleus fragments at the first interaction completely into free nucleons (NFRAGM = 0) or successively by assuming that the non-interacting nucleons proceed as one new nucleus (NFRAGM = 1). This new nucleus may evaporate nucleons or alpha-particles with a transverse momentum distribution according to experimental data [74] (NFRAGM = 2, default) or with a transverse momentum distribution according to Goldhaber's theory [75] using 0.090 GeV/nucleon as the average transverse momentum (NFRAGM = 3). NFRAGM = 4 gives identical fragments as NFRAGM = 2 or 3, but without transverse momenta.

The NFRAGM flag is used also to steer the fragmentation in the various interaction models as described for the HDPM routines. EPOS, NEXUS, and VENUS use the same evaporation model as HDPM with the same meaning of NFRAGM, while SIBYLL and QGSJET deliver themselves realistic nuclear fragments with according transverse momenta; they are selected by NFRAGM  $\geq$  2. Therefore the nuclear evaporation as used for HDPM, EPOS, NEXUS, and VENUS is coupled with DPMJET and the meaning of NFRAGM follows HDPM.

Limits are:  $0 \leq \text{all flags} < 100$

#### 4.43 Neutrino Interaction Type Selection

NUSLCT    NUSLCT

Format = (A6, I), Default = 2

NUSLCT : Selects the type of the primary neutrino interaction: 0 = neutral current interaction, 1 = charged current interaction, *else* = type of interaction is selected at random according to the interaction cross sections for the two processes.

This keyword is only available in the NUPRIM option.

#### 4.44 Charm Interaction Cross Section

SIGMAQ    SIGMAQ(i), i=1... 4

Format = (A6, F4), Defaults = 0., 0., 0., 0.

SIGMAQ(i) : The interaction cross-sections for charmed mesons ( $i = 1$ ), charmed baryons ( $i = 2$ ), bottom mesons ( $i = 3$ ), and bottom baryons ( $i = 4$ ) are specified. For SIGMAQ(i) = 0., the parameterizations of Ref. [48] are used; for SIGMAQ(i)  $\neq$  0., the energy independent cross-section of SIGMAQ(i) is used for the corresponding projectile.

Limits are:  $0. \leq \text{SIGMAQ}(i)$ .

This keyword is only available in the CHARM option.

## 4.45 Charm Interaction Selection

PROPAQ    PROPMOD

Format = (A6, I), Default = 1

PROPMOD : If set to 1, the PYTHIA extension [48] is used to treat the interactions of charmed projectiles.

If set to 0, the QGSJET01d interaction model is used to treat the interactions of charmed projectiles.

Limits are:  $0 \leq \text{PROPMOD} \leq 1$ .

This keyword is only available in the CHARM option in combination with the QGSJET option.

## 4.46 Transition Energy between Models

HILOW    HILOELB

Format = (A5, F), Default = 80.

HILOELB : Allows to define the transition energy<sup>74</sup> ( $E_{lab}$  in GeV) between high and low-energy hadronic interaction model.

Because EPOS can describe data very well down to lower energies, the default value for this model is 30 GeV. Limits depend on the used interaction model, for most high-energy hadronic interaction models the low-energy limit is in the range of  $\approx 80$  GeV, for SIBYLL  $\approx 60$  GeV, for EPOS  $\approx 20$  GeV, while most low energy models enable a limit as high as several 100 GeV.

## 4.47 Electromagnetic Interaction Steering Flags

ELMFLG    FNKG    FECS

Format = (A6, 2L), Defaults = T, T

FNKG : If .true., the NKG option is switched on for calculating the electromagnetic sub-cascades analytically. For the electron kinetic energy threshold the value of ELCUT(3) is taken (keyword ECUTS page 100). If .false., the NKG option is disabled<sup>75</sup>.

FECS : If .true., the EGS4 option is selected to calculate all interactions of  $e^+$ ,  $e^-$ , and photons in the atmosphere explicitly. (The second random number sequence should be

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<sup>74</sup>With the INTTEST option the default value of this border is at 49 GeV resp. 101 GeV, depending whether a high- or low-energy hadronic interaction model should be tested. For testing of DPMJET, EPOS, NEXUS, QGSJET, and VENUS the default value is set to 49 GeV, for SIBYLL to 60 GeV. If none of those models is selected, the default value is set to 101 GeV to test the models FLUKA, GHEISHA, or UrQMD.

<sup>75</sup>In the CURVED option the NKG formulas are no longer valid, therefore the NKG flag is disabled automatically in this option. The NKG flag should be disabled in the COMPACT option, as the resulting NKG parameters cannot be written out onto the particle output file.



initialized for use in the EGS4 part. Otherwise the default initialization is taken.) If .false., the EGS4 option is disabled.

In the CERENKOV and PRESHOWER options this flag is obsolete as EGS4 is selected automatically.

The two options may be selected or disabled independently at the same time.

#### 4.48 Electron Multiple Scattering Length Factor

STEPFC STEPFC

Format = (A6, F), Default = 1.

STEPFC : Factor by which the multiple scattering length for electrons and positrons in EGS4 simulations is elongated relative to the value given in [20]. A detailed discussion on the use of the step length is given in [38]. An enlargement of this factor may be tolerated to reduce computing time, but simultaneously the electron lateral distribution on ground becomes slightly narrower. With STEPFC = 10. the CPU-time is reduced by a factor of  $\approx 1.7$  (relative to the default value). A reduction of STEPFC will increase the computing time considerably, e.g with STEPFC = 0.1 by a factor of  $\approx 5$ .

Limits are:  $0. < \text{STEPFC} \leq 10.0$

This keyword is not available in the CERENKOV option.

#### 4.49 Radius of NKG Lateral Range

RADNKG RADNKG

Format = (A6, F), Default = 200.E2

RADNKG : Gives the outer range radius (in cm) within which the lateral NKG distribution is calculated for 10 radii equidistant in logarithmic scale. The inner radius is always kept at 100 cm.

Limit is: RADNKG > 100.

#### 4.50 Flags for Simplified CONEX Threshold Management

CASCADE FCXCAS FCXLCE FCXGHE

Format = (A7, 3L), Defaults = T, F, F

FCXCAS : If .true. optimized parameters are used to run in hybrid mode (MC+CE+MC) with full 3D information. Muons are fully tracked in MC. If .false. the cascade equations are disabled below the high-energy thresholds and only MC is used.

FCXLCE : If .true. default parameters are used to run in hybrid mode (MC+CE) with full 3D information for muons only. For hadrons and e/m particles, only 1D informations

are relevant (longitudinal profile and energy distribution of particles at ground). Muons are fully tracked in 3D in MC. If .false. optimized sampling weights are used to run MC below the default low-energy thresholds. Weights defined by THIN keyword are used for thinning in CONEX and sampling weights are used as thinning weights in CORSIKA (see FCXCE in Sect. 4.53 page 95).

FCXGHE : If .true. default parameters are used to run in hybrid mode (MC+CE) for all particles. Only 1D informations are relevant (longitudinal profile and energy distribution of particles at ground). If .false. optimized sampling weights are used to run MC for muons below the default low-energy thresholds. Weights defined by THIN keyword are used for thinning in CONEX and sampling weights are used as thinning weights in CORSIKA (see FCXCE in Sect. 4.53 page 95).

Limits are: Only the combinations of the three flags T T T, T T F, T F F, or F F F are allowed.

**Notes:** This keyword should be used if the user wants to select easily different hybrid mode. In particulars the first MC part of the shower is not changed by changing the value of CASCADE (or even by introducing this command), thus a shower with the same longitudinal profile can be calculated with different speed and precision (1D or 3D) if the 7<sup>th</sup> SEED sequence is not changed of course.

The sampling weights are adjusted to get a good compromise between precision and speed. If ECUTS are defined below CONEX limits (1 GeV for hadrons and muons and 1 MeV for e-m particles), low energy MC (T F F but with larger weight and lower threshold) is automatically switched on for the relevant part of the shower and to get reasonable precision in 1D, but it will be necessary slower than default T T F or T T T mode. To be able to reproduce a given shower, ECUTS should not be changed while changing CASCADE configuration.

To get the fastest simulation for any ECUTS see the keyword CORSIKA in Sect. 4.54 page 96.

It is recommended not to use this keyword together with the CX2COR or CXWMX keywords since it redefines some of the parameters.

This keyword must be called after the energy cuts are defined (see keyword ECUTS Sect. 4.63 page 100).

This keyword is only available in the CONEX option.

## 4.51 Thresholds for CONEX Cascade Equations

CONEX    CXTHR(1)    CXTHR(2)    CXTHR(3)

Format = (A5, 3F), Defaults = 1.E-3, 1., 1.E-4

CXTHR(*i*) : Fractions of primary energy above which the shower particles are treated individually by Monte Carlo methods. Below these thresholds cascade equations treat the particles. These thresholds apply to hadrons (*i* = 1), muons (*i* = 2), and electrons/photons (*i* = 3).

Limits are:  $0. \leq \text{CXTHR}(i) \leq 1.$

This keyword is only available in the CONEX option.

## 4.52 Thresholds for Transition from CONEX to CORSIKA

CX2COR    CXMCT(1)    CXMCT(2)    CXMCT(3)    CXMCS

Format = (A6, 4F), Defaults = 3.E2, 1.E20, 1.E1, 4.E2

CXMCT(*i*) : Thresholds (in GeV) below which the particles are transferred from CONEX back to the CORSIKA stack. These thresholds apply to hadrons (*i* = 1), muons (*i* = 2), and electrons/photons (*i* = 3).

CXMCS : Minimal vertical depth (g/cm<sup>2</sup>) above observation level below which the transfer of particles from CONEX to CORSIKA treatment is started.

Limits are:  $0. \leq \text{CXMCT}(i) \leq 1.E20$ ;  $0. \leq \text{CXMCS} \leq 1.E20$

In the CONEX option this keyword has to be called after the keyword CASCADE T F F (see Sect. 4.50 page 93) is used to be effective, but it is better not to use it with CASCADE keyword.

This keyword is only available in the CONEX option.

## 4.53 Weight Sampling for CONEX

CXWMX    CXWMT(1)    CXWMT(2)    CXWMT(3)    FCXWMX    FCXCE

Format = (A5, 3F, L), Defaults = -1., -1., -1., F, F

CXWMT(*i*) : Factors (GeV<sup>-1</sup>) to be multiplied with the primary energy to get maximal sampling weights for particles transferred from CONEX to CORSIKA. These factors apply to hadrons (*i* = 1), muons (*i* = 2), and electrons/photons (*i* = 3). For negative or zero values of CXWMT(*i*) the CORSIKA thinning weight limits (see Sects. 4.55, 4.56, and 4.57) are taken as sampling weights too.

FCXWMX : If .true. forces the usage of CONEX sampling weight limits also in thinning (in CONEX and CORSIKA) and these weight limits are determined with the CXWMT(*i*) factors. If .false. and FCXCE = .false. the CORSIKA weight limits are used as thinning weights for both CONEX and CORSIKA (see Sect. 4.50) and forces the usage of the CONEX weight limits also in CORSIKA.

FCXCE : If .true. forces the usage of CONEX sampling weight limits also in thinning for CORSIKA only and these weight limits are determined with the CXWMT(*i*) factors. The CORSIKA weight limits are used as thinning weight for CONEX only. If FCXCAS = .true. (see Sect. 4.50) forces FCXCE to be .true.

Limits are:  $-1. \leq \text{CXWMT}(i) \leq 1.$

In the CONEX option this keyword has to be called after the keyword CASCADE (see page 93) is used to be effective, but it is better not to use it with CASCADE keyword.

This keyword is only available in the CONEX option.

## 4.54 CONEX without CORSIKA

CORSIKA FCORS

Format = (A7, L), Defaults = T

FCORS : If .true. allows CORSIKA to run low energy MC. If .false. all particles which are supposed to go into the CORSIKA stack are lost and only CONEX is run with CE going to the minimum energy. This is equivalent to CASCADE T T T (see Sect. 4.50) but with no influence of ECUTS.

This keyword should be used with care because only the longitudinal energy deposit profile is valid. Even the  $X_{max}$  fit based on the number of electrons could be biased because ECUTS below CONEX limits are not taken into account. This keyword is useful to get very fast 1D energy deposit profile. The change of FCORS doesn't change the first interactions of the shower so with the same SEED the same  $X_{max}$  for the energy deposit will be obtained.

This keyword is only available in the CONEX option.

## 4.55 Thinning Definition

THIN EFRCTHN WMAX RMAX

Format = (A4, 3F), Defaults = 1.E-4, 1.E30, 0.E0

EFRCTHN : Factor  $\varepsilon_{th}$  which defines the energy fraction of the primary energy below which the thinning algorithm becomes active. If the fraction is selected in a manner that this energy is below the lowest energy threshold of  $ELCUT(i)$ ,  $i = 1..4$  (keyword ECUTS page 100), thinning will not become active, but the particle output data structure will contain the weight (= 1.) for each particle.

WMAX : Weight limit for thinning. If the weight of a particle exceeds WMAX, no further thinning is performed<sup>76</sup>.

RMAX : Maximum radius (in cm) at observation level within which all particles are subject to inner radius thinning. Particles are selected with probability  $(r/rmax)^4$ . The weight of surviving particles is multiplied by the appropriate factor (inverse of probability). This thinning neither affects the shower development nor the table output nor the histogram output of the ANAHIST, AUGERHIST, or MUONHIST options, rather only the particle output file written onto MPATAP (and the Cherenkov output file written onto MCETAB). For  $RMAX \leq 0$ . no radial thinning is applied.

If the keyword CORECUT (Sect. 4.62 page 99) is used with  $RCUT > 0$ . the radial thinning will not be effective in favour of the radial cutting.

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<sup>76</sup>See footnote page 68.

Limits are:  $ULIMIT \cdot EFRCTHN \leq 1 \cdot 10^7$  GeV (for ULIMIT see keyword ERANGE Sect. 4.6 page 76);  $0.1 \leq WMAX \leq 1 \cdot 10^{20}$  .

This keyword is only available in the THIN option.

## 4.56 Hadronic Thinning Definition

THINH THINRAT WEITRAT

Format = (A5, 2F), Defaults = 1., 1.

THINRAT : Defines hadronic thinning limit differing from em-thinning limit EFRCTHN by the ratio of  $\varepsilon_{th_{em}}/\varepsilon_{th_{hadr}}$  which gives the ratio between the energy of the em-particles (specified by keyword THIN, see above) and the energy of the hadronic particles below which the thinning algorithm becomes active for these particle species (see also Sect. 4.55 above).

WEITRAT : Defines hadronic weight limit differing from em-weight limit WMAX by the ratio<sup>77</sup> of weight limit of em-particles to weight limit of hadronic particles in case of thinning (see also Sect. 4.55 above).

A simultaneous use of the keyword THINH together with THINEM is not tolerated and will lead to an error stop.

Limits are:  $ULIMIT \cdot EFRCTHN / THINRAT \leq 1 \cdot 10^7$  GeV (for ULIMIT see keyword ERANGE Sect. 4.6 page 76);  $1 \cdot 10^{-4} \leq WEITRAT \leq 1 \cdot 10^6$  .

This keyword is only available in the THIN option.

## 4.57 Electromagnetic Thinning Definition

THINEM THINRAT WEITRAT

Format = (A6, 2F), Defaults = 1., 1.

THINRAT : Defines em-thinning limit differing from hadronic thinning limit EFRCTHN by the ratio of  $\varepsilon_{th_{em}}/\varepsilon_{th_{hadr}}$  which gives the ratio between the energy of the em-particles and the energy of the hadronic particles (specified by keyword THIN, see above) below which the thinning algorithm becomes active for these particle species (see also Sect. 4.55 above).

WEITRAT : Defines em-weight limit differing from hadronic weight limit WMAX by the ratio<sup>77</sup> of weight limit of em-particles to weight limit of hadronic particles in case of thinning (see also Sect. 4.55 above).

A simultaneous use of the keyword THINEM together with THINH is not tolerated and will lead to an error stop.

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<sup>77</sup>In the Slovenian thinning [68] an  $\epsilon$  is defined which gives the inverse of WEITRAT. There the weight limit  $w_{max}$  is defined for em-particles and from this the weight limit for hadrons and muons is derived by  $\epsilon \cdot w_{max}$ .

Limits are:  $ULIMIT \cdot EFRCTHN \cdot THINRAT \leq 1 \cdot 10^7$  GeV (for ULIMIT see keyword ERANGE Sect. 4.6 page 76);  $1 \cdot 10^{-4} \leq WEITRAT \leq 1 \cdot 10^6$  .

This keyword is only available in the THIN option.

## 4.58 Multiple Thinning Definition

MTHINH    EFRCTHN    WMAX    THINRAT    WEITRAT

Format = (A6, 4F), Defaults = 1.E-4, 1.E30, 1., 1.

This keyword has to appear once for each thinning mode to be defined. At maximum up to 6 thinning modes are possible. Their sequence is arbitrary.

EFRCTHN : Factor  $\varepsilon_{thhad}$  which defines the energy fraction of the primary energy below which the thinning algorithm becomes active for hadrons. If the fraction is selected in a manner that this energy is below the lowest energy threshold of ELCUT( $i$ ),  $i = 1 \dots 4$  (keyword ECUTS page 100), thinning will not become active.

WMAX : Weight limit for thinning for hadrons. If the weight of a particle exceeds WMAX, no further thinning is performed<sup>78</sup>.

THINRAT : Defines em-thinning differing from hadronic thinning limit EFRCTHN by the ratio of  $\varepsilon_{them}/\varepsilon_{thhad}$  which gives the ratio between the energy of the em-particles and the energy of the hadronic particles below which the thinning algorithm becomes active for these particle species (see also Sect. 4.55 page 96).

WEITRAT : Defines em-weight limit differing from hadronic weight limit WMAX by the ratio of weight limit of em-particles to weight limit of hadronic particles in case of thinning (see also Sect. 4.55 page 96 and Sect. 4.57 page 97).

Limits are:  $ULIMIT \cdot EFRCTHN \leq 1 \cdot 10^7$  GeV (for ULIMIT see keyword ERANGE Sect. 4.6 page 76);  $0.1 \leq WMAX \leq 1 \cdot 10^{20}$  ;  $ULIMIT \cdot EFRCTHN / THINRAT \leq 1 \cdot 10^7$  GeV (for ULIMIT see keyword ERANGE Sect. 4.6 page 76);  $1 \cdot 10^{-4} \leq WEITRAT \leq 1 \cdot 10^6$  .

This keyword is only available in the MULTITHIN option.

## 4.59 Random Number Generator Initialization for MULTITHIN Modes

MSEED    ISEED(i, j+10), i=1... 3

Format = (A5, 3I), Defaults =  $j + 10$ , 0, 0

ISEED(1,  $j + 10$ ) : Contains the seed of the random number sequence  $j + 10$  .

ISEED(2..3,  $j + 10$ ) : Contain the number of calls  $N_{in}$  to the generator that are performed for initialization such that  $N_{in} = ISEED(2, j + 10) + 10^9 \cdot ISEED(3, j + 10)$ .

At present at most  $j = 6$  sequences are used to initialize the random generator sequence  $j + 10$  for the  $j^{th}$  MULTITHIN mode. Their activation follows the sequence of occurrence

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<sup>78</sup>See footnote page 68.

of the keyword lines.

The use of  $\text{ISEED}(2, j + 10) > 0$  and especially of  $\text{ISEED}(3, j + 10) > 0$  should be avoided as presetting the random number generator by billions of calls needs considerable computing time. To get different random sequences it is sufficient to modify  $\text{ISEED}(1, j + 10)$ . Limit (to get independent sequences of random numbers) is:  $1 \leq \text{ISEED}(1, j + 10) \leq 900\,000\,000$ .

This keyword is only available in the MULTITHIN option.

## 4.60 Maximum Radius in Multiple Thinning

MTHINR      RCUT

Format = (A6, F), Default = 0.

RCUT : Maximum radius (in cm) at observation level within which all particles are discarded. The value of RCUT is stored in the eventheader in EVTH(152).

For  $\text{RCUT} \leq 0$ . no particles are discarded.

RCUT overrides RMAX of the keyword THIN (Sect. 4.55 page 96).

This keyword is only available in the MULTITHIN option.

## 4.61 Multiple Thinning used by COAST

MWEIC      IMWEIC      MTHOUT

Format = (A5, I, L), Default = 0, T

IMWEIC : Index (from 0 to the number of MTHINH lines) of the MTHINH line used by COAST for the weight of particles. The value of IMWEIC is stored in the eventheader in EVTH(223).

For  $\text{IMWEIC} = 0$  unthinned particles are used.

MTHOUT : Enable/Disable writing the weight sub-block from MULTITHIN to get a normal output file of unthinned particles but COAST uses the thinned shower corresponding to IMWEIC.

This keyword is only available in the MULTITHIN option with COASTUSERLIB.

## 4.62 CORECUT Maximum Radius Cut

CORECUT      RCUT

Format = (A7, F), Default = 0. (rsp. 200.E2 for AUGERHIT option)

RCUT : Maximum radius (in cm) at observation level within which all particles are discarded. The value of RCUT is stored in the eventheader in EVTH(152).

For  $\text{RCUT} \leq 0$ . no particles are discarded.

RCUT overrides RMAX of the keyword THIN (Sect. 4.55 page 96).

### 4.63 Energy Cut-Offs

ECUTS      ELCUT(1)      ELCUT(2)      ELCUT(3)      ELCUT(4)

Format = (A5, 4F), Defaults<sup>79</sup> = 0.3, 0.3, 0.003, 0.003

ELCUT(*i*) : The low energy cut-off (in GeV) of the particle kinetic energy may be chosen differently for hadrons (without  $\pi^0$ 's) (*i* = 1), muons (*i* = 2), electrons (*i* = 3), and photons (including  $\pi^0$ 's) (*i* = 4). For nuclei ELCUT(1) is applied to the energy per nucleon. It is in the responsibility of the user to choose the cut-off values in a reasonable way not to eliminate those parent particles which might decay to secondaries which you are looking for in the investigated problem (e.g. decay of muons to electrons).

Limits are: ELCUT(1)<sup>80</sup>  $\geq 0.05$  ; ELCUT(2)<sup>81</sup>  $\geq 0.01$ ; ELCUT(3), ELCUT(4)  $\geq 0.00005$   
The value of ELCUT(3) is also taken as threshold value for the NKG calculation. In this case an upper limit of ELCUT(3)  $< 0.08$  is recommended.

In the CONEX option this keyword has to be called before the keyword CASCADE (see page 93) is used.

### 4.64 Time Cut-Off

TIMLIM      DSTLIM      LTMLMPR

Format = (A6, F, L), Defaults = 1.D8, F

DSTLIM : Gives the distance (in cm) a particle would travel with velocity of light downstream the detector before cut away by the time limit. An additional security time of 25  $\mu$ sec (corresponding with  $\approx 7$  km) is taken into account.

DSTLIM can be positive or negative, but below 20km you must be careful. With the INCLINED option, the minimum time is computed with respect to the inclined observation plane origin.

This option has to be used carefully not to lose particles which should be counted.

LTMLMPR : Flag which enables (T) or disables (F) printing of particles which exceed the time limit.

This keyword is only available in the CURVED option.

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<sup>79</sup>For the INTTEST option: Defaults = 0., 0., 0., 0. and all limits are ELCUT(*i*)  $\geq 0$ .

<sup>80</sup>ELCUT(1) is used also for neutrinos in the NEUTRINO option. In the FLUKA versions ELCUT(1) should be  $\geq 0.02$  GeV; in the URQMD option ELCUT(1) should be  $\geq 0.3$  GeV.

<sup>81</sup>ELCUT(2) is used also for  $\tau$ -leptons in the CHARM or TAULEP option.



## 4.65 Longitudinal Shower Development

LONGI    LLONGI    THSTEP    FLGFIT    FLONGOUT

Format = (A5, L, F, 2L), Defaults = F, 20.0, F, F

LLONGI : If .true., the longitudinal development of particle numbers for gammas (EGS4), positrons (EGS4), electrons (EGS4), positive and negative muons, hadrons, all charged, nuclei, and Cherenkov photons (CERENKOV) is sampled. Moreover the longitudinal development of the energy content in the various particle species (same order as before, but without Cherenkov photons) is sampled. Additionally the longitudinal development of energy deposit by ionization energy loss and by angular or energy cuts is sampled. See also Sect. 10.1 page 133. To get the sampling in slant depth instead of the (default) vertical depth you should use the SLANT option (page 66). If .false., the longitudinal development is not sampled.

THSTEP : Vertical step width (rsp. slant step width in the SLANT option) for sampling of the longitudinal development (in g/cm<sup>2</sup>). The sampling is done in vertical (rsp. slant) depth. The altitudes are not depending on the zenith angle of the primary particle (except the preprocessor option SLANT has been selected). In the CURVED option the minimum step size has to be selected in a manner that no more than 14998 steps are needed to pass through the complete atmosphere.

FLGFIT : If .true. and LLONGI also .true., the longitudinal development of all charged particles number is fitted. If .false., the fit is suppressed.

FLONGOUT : If .true. and LLONGI also .true., the longitudinal distributions of particle numbers and energy deposit for the various particle groups are written to the file 'DATnnnnnnn.long' (rsp. 'DATnnnnnnnnnn.long' with option NRREXT, see Sect. 3.5.24 page 61) (see Sect. 10.5 page 151).

If .false. and LLONGI .true., the longitudinal distributions only of the particle numbers for the various particle species are written out to the particle output file 'DATnnnnnnn' (rsp. 'DATnnnnnnnnnn') in extra 'LONG' sub-blocks (see Sect. 10.2, Table 6 page 136 and Table 15 page 145).

Limits are: 1. ≤ THSTEP ≤ 14998.

20. ≤ THSTEP ≤ 14998 for the SLANT option and horizontal incidence.

Normally only to the number distribution of all charged particles a function of the Gaisser-Hillas type [76]

$$N(t) = N_{max} \cdot \left( \frac{t - t_0}{t_{max} - t_0} \right)^{\frac{t_{max} - t_0}{a + bt + ct^2}} \cdot \exp\left( \frac{t_{max} - t}{a + bt + ct^2} \right)$$

is fitted to describe the dependence on the atmospheric depth  $t$  and the resulting 6 parameters  $N_{max}$ ,  $t_0$ ,  $t_{max}$ ,  $a$ ,  $b$ , and  $c$  and the  $\chi^2/dof$  are stored in the event end block. The longitudinal development of the electromagnetic particles is only sampled if EGS4 is selected (see ELMFLG). If only NKG is activated the fit is applied to the NKG longitudinal distribution which consists of particle numbers from only  $\leq 10$  levels. If neither EGS4

nor NKG is selected the charged particle distribution contains only muons and charged hadrons. In the AUGERHIST option also a Gaisser-Hillas type function is fitted to the longitudinal energy deposit, if EGS4 is selected.

In the Cherenkov versions the longitudinal distribution of photons is given in differential mode (i.e. the number of photons generated within each step) as default. By the preprocessor option INTCLONG the integral mode is selected (i.e. accumulated number of generated Cherenkov photons for each step) which needs additional computing time. If both kinds of the distribution are of no interest, you may deselect the Cherenkov photon distribution completely by the preprocessor option NOCLONG thus saving computing time.

## 4.66 Muon Multiple Scattering Treatment

MUMULT FMOLI

Format = (A6, L), Default = T

FMOLI : If `.false.`, the muon multiple scattering angle is selected by Gauss approximation. If `.true.`, the muon multiple scattering angle is selected for large steps by Molière's theory and for small steps by adding many single Coulomb scattering events.

## 4.67 Additional Electromagnetic Particle Information

EMADDI FEMADD

Format = (A6, L), Default = F

FEMADD : If `.false.`, no additional information on electromagnetic (EM) particles is written to particle output file. If `.true.`, additional information on mother and grandmother hadrons of EM particles at the origin of the EM subshower is written to the particle output file. Details are similar to those given in Ref. [57] for muons in the EHISTORY option (Sect. 3.5.15 page 55). In case of em-particle primaries the additional particle information is suppressed (by automatically setting FEMADD = `.false.`) as for the bulk of the em-particles no hadronic mother or grandmother particles exists.

This keyword is only available in the EHISTORY option and not in the INTTEST option.

## 4.68 Additional Muon Information

MUADDI FMUADD

Format = (A6, L), Default = F

FMUADD : If `.false.`, no additional muon information is written to particle output file. If `.true.`, additional information on muons at their origin is written to the particle output

file. This additional muon information consists of 7 data words according to Table 10 (page 141) and precedes the corresponding muon particle on the particle output file. The first data word contains the particle identification 75 (for  $\mu^+$ ) or 76 (for  $\mu^-$ ) combined with 3 digits of the hadronic generation counter (instead of the 2 digits + 1 digit for observation level), which in this case may differentiate between muons originating from K-decay (normal generation counter) and  $\pi^\pm$ -decay (generation counter incremented by 50) and can be larger than 100 in some case of electromagnetic interaction. To indicate clearly the electromagnetic origin of the muon (muon pair production or photonuclear interaction in muon history) the hadronic generation counter is increased by 500. The 7<sup>th</sup> data word contains the altitude (in cm) of the muon birth instead of time (see also Table 10 page 141).

In combination with the EHISTORY option (see Sect. 3.5.15 page 55) an extended additional muon information is written to the particle output file. In the combination of the MUPROD option (see Sect. 3.5.22 page 60) with the EHISTORY option the extended muon information is written also for muons which decay or interact before reaching the lowest observation level. Details are given in Ref. [57].

This keyword is not available in the INTTEST option.

## 4.69 Additional Neutrino Information

NUADDI      FNUADD

Format = (A6, L), Default = F

FNUADD : If .false., no additional neutrino information is written to particle output file. If .true., additional information on mother and grandmother hadrons of neutrinos at their origin is written to the particle output file. Details are similar to those given in Ref. [57] for muons in the EHISTORY option (Sect. 3.5.15 page 55). An extended additional neutrino information is written to the particle output file.

This keyword is only available in the EHISTORY option and not in the INTTEST option.

## 4.70 Observation Level Definition

OBSLEV      OBSLEV(i)

Format = (A6, F), Default = 110.E2

OBSLEV(i) : Observation level  $i$  above sea level (in cm). This keyword has to appear once for each level to be defined. At maximum up to 10 observation levels are possible<sup>82</sup>. Their sequence is arbitrary.

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<sup>82</sup>Only one observation level is possible in the CURVED option. Up to 20 levels might be specified for the production of histograms in the AUGERHIST option, the lowest observation level must be at minimum 1 g/cm<sup>2</sup> above sea level.

In the UPWARD option (page 70) for upward going primaries the observation level should be chosen preferentially at the top of atmosphere, but at minimum above the starting point of the shower. The value of OBSLEV has to be selected in a manner that the shower axis crosses the observation level.

Limits are<sup>83</sup>:  $0 \leq \text{OBSLEV}(i) < \text{top of atmosphere}$

In the CURVED option only one observation level can be defined.

## 4.71 Inclined Observation Plane

INCLIN    XPINCL    YPINCL    ZPINCL    THINCL    PHINCL    TDINCL

Format = (A6, 5F), Defaults = 0., 0., *OBSLEV*(1), 0., 0., 0.

XPINCL : X-coordinate (in cm) of reference point in inclined observation plane.

YPINCL : Y-coordinate (in cm) of reference point in inclined observation plane.

ZPINCL : Z-coordinate (in cm) of reference point in inclined observation plane.

THINCL :  $\theta$ -angle (in deg) of normal vector of inclined observation plane.  $\theta = 0.$  points to the zenith (opposite to Fig. 1 page 126) and defines a horizontal plane.

PHINCL :  $\phi$ -angle (in deg) of normal vector of inclined observation plane.

TDINCL : thickness (in g/cm<sup>2</sup>) of the origin of an observation plane perpendicular to the shower axis (automatic definition of all other numbers which are not used if TDINCL<sub>i</sub>0).

The behavior of this parameter depends on the SLANT option. If SLANT is defined, TDINCL define the slant depth along the shower axis, where the perpendicular plane should be placed. The origin to count the slant depth depends on the use of FIXHEI. If this parameter is defined, then the slant depth is counted from this point. Otherwise it will count from the starting point of the shower. Without SLANT, TDINCL will define the fixed vertical depth (= fixed height) at which the plane should be placed. In both cases, CORSIKA automatically defines THINCL=THETAP and PHINCL=PHIP (shower by shower), to have a plane perpendicular to the shower axis, and compute XPINCL, YPINCL and ZPINLC such that the axis is at the origin of the frame of the inclined plane at the right vertical or slant depth. The actual values of all the parameters including the thickness are registered is EVTH (see Tab. 9 page 140).

All coordinates are defined in the standard CORSIKA output coordinate system (see Fig. 1 page 126) with its origin at the point where the shower axis hits the observation level. Detailed description of the coordinate definition is given in Fig. 3 page 149.

**ATTENTION:** CORSIKA discards all particles BELOW the height of the lowest standard observation level (see Sect. 4.70 page 103 for definition of OBSLEV), so as soon as the inclined observation plane drops below the lowest standard (horizontal) observation level, there will be no particles recorded any more. In case of UPWARDOLD option, the particles ABOVE OBSLEV are discarded. This can also be used to limit the computation time placing the horizontal observation level at the right place compared to the inclined

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<sup>83</sup>For atmospheric models  $1 \leq \text{MODATM} \leq 9$  and  $\text{MODATM} \geq 17$  limits are:  $-1.E5 < \text{OBSLEV}(i) < 112.8E5$

plane (in particular with UPWARDOLD for upward going showers). TIMLIM (Sect. 4.64 page 100) has a special behavior and can be used carefully to limit the computation time when the CURVED option is selected.

Particles on the inclined plane will be stored in a coordinate system within the inclined plane, with the origin defined at (XPINCL, YPINCL, ZPINCL). See Sect. 10.3 page 148 for details.

This keyword uses the COAST package and is available *only* when the INCLINED option (see page 28) in `./coconut` has been selected.

Technically, the INCLINED option uses a particular COASTUSERLIB library (Sec. 3.5.8 page 51), namely `InclinedPlane`, which is shipped with COAST in the `CorsikaOptions` directory by default.

This keyword is only available in the INCLINED option.

## 4.72 Observation Level Curvature

CURVOUT      FCURVOUT

Format = (A7, L), Default = T

FCURVOUT : If `.true.` the observation level is a sphere following the Earth curvature at an altitude  $H = OBSLEV(1)$  and  $(X, Y)$  in the output file are replaced by  $(X', Y')$  which can be used to calculate the angles necessary to obtain  $(X, Y, Z)$  in a Cartesian frame: defining  $\theta$  and  $\phi$  the spherical coordinates of the particle as

$$\theta = \frac{\sqrt{X'^2 + Y'^2}}{R_{Earth} + H}$$

$$\phi = \text{atan2}(Y', X')$$

with  $R_{Earth} = 637131500$  cm. At the observation level ( $H = OBSLEV(1)$ ) the Cartesian coordinates can be obtained using  $D = (R_{Earth} + H) \cdot \sin \theta$  as

$$X = D \cdot \cos \phi$$

$$Y = D \cdot \sin \phi$$

$$Z = (R_{Earth} + H) \cdot \cos \theta - R_{Earth}$$

If  $\theta$  is small (close to shower core) then  $(X', Y') = (X, Y)$ . If `.false.` the observation level is considered as flat and particle positions  $(X, Y)$  are defined in a Cartesian frame with the origin  $(0, 0)$  at the core position. Particles are tracked only until  $Z = OBSLEV(1)$  in this Cartesian frame (apparent height). As a consequence, far from the core, some particles can be discarded before reaching the ground making a bias in the longitudinal profile. To avoid such a problem, `FCURVOUT = .false.` is forbidden by default if the zenith angle is between  $85^\circ$  and  $95^\circ$ .

As usually the Cherenkov detector geometries are small, the `FCURVOUT` is set `.false.` in the CERENKOV option.

The keyword **FLATOUT** = .not.FCURVOUT can still be used for backward compatibility.

This keyword is only available in the CURVED option and not in the ANAHIST or AUGERHIST options.

### 4.73 Array Rotation

ARRANG ARRANG

Format = (A6, F), Default = 0.

ARRANG : Defines a rotation angle (in °) between the detector array x-direction and magnetic north direction; positive if detector array x-direction points to the West.

**Attention:** The default value of ARRANG is set differently to  $-92.08^\circ$  if the AUGERHIT option (see Sect. 3.5.5 page 49) is selected.

Limits are:  $-180. \leq \text{ARRANG} \leq 180.$

This keyword is not available in the COASTUSERLIB option.

### 4.74 Auger Detector Scattering

AUGSCT MAUGPOS DRADIUS DETDIS FTANKSHADOW FANYMODE

Format = (A6, I, 2F, 2L), Default = 1, 35., 1500., .true., .true.

MAUGPOS : Number of uses of each event for the Auger detector. For MAUGPOS = 0 the scatter positions are expected to be given via the keyword AUGHIT (see Sect. 4.75). For MAUGPOS > 0 the scatter positions are taken at random employing the SOBOL quasi-random generator [43].

DRADIUS : Half width (in meter) of the stripes in which the Auger detector rows are aligned.

DETDIS : Shortest distance (in meter) between two tanks on the hexagonal tank array.

FTANKSHADOW : Enable (=true.) or disables (=false.) the check whether a particle falls in a tank shadow. One should select DRADIUS large enough for long shadows of very inclined particles. The tank dimensions are fixed at 1.8 m radius and 1.2 m height as used for the Auger detectors.

FANYMODE : Enable (=true.) or disable (=false.) the check whether a particle survives any thin mode of the MULTITHIN option (see Sect. 3.5.20 page 59).

Limits are:  $0 \leq \text{MAUGPOS} \leq 20$ ;  $\text{DRADIUS} \geq 0.$ ;  $\text{DETDIS} \geq 0.$

This keyword is only available in the AUGERHIT option.

## 4.75 Auger Detector Scattering Positions

AUGHIT XSHCORE(i) YSHCORE(i)

Format = (A6, 2F), Default = 0., 0.

XSHCORE(i) : Detector position in x-direction (in meter) of the  $i^{th}$  scattering relative to the shower core.

YSHCORE(i) : Detector position in y-direction (in meter) of the  $i^{th}$  scattering relative to the shower core.

This keyword has to be activated by setting the parameter MAUGPOS = 0 of the keyword AUGSCT (see Sect. 4.74). The positions have to be specified in the detector coordinate frame (with the positive X-axis pointing to East), the length units are meter. For each scatter position its own input line is needed and up to 20 lines are tolerated in arbitrary sequence.

This keyword is only available in the AUGERHIT option.

## 4.76 String Detector Configuration

DETCFG DETCFG

Format = (A6, F), Default = 0.

DETCFG : Gives the geometry configuration of a long vertical string detector as the ratio *height/diameter*.

Limit is: DETCFG > 0.

This keyword is only available in the VOLUMECORR option.

## 4.77 Event Printout

MAXPRT MAXPRT

Format = (A6, I), Default = 1

MAXPRT : Is the maximum number of events that produce a detailed printout during the simulation run.

Limit is: MAXPRT  $\geq$  1

## 4.78 Particle Printout

ECTMAP ECTMAP

Format = (A6, F), Default = 1.E11

ECTMAP : Defines a cut in the particle  $\gamma$  factor (or energy in GeV for em-particles and

neutrinos) above which they are printed out on the logical unit MONIOU when passing an observation level.

## 4.79 Output Directory

DIRECT      DSN

Format = (A6, A239), Defaults = 'anynameupto239characters/'

DSN : May be used to define a name of an output directory. Lower case characters of DSN are not converted to capitals. Do not use capitals with the ANAHIST, AUGERHIST, MUONHIST, or INTTEST options as the HBOOK routines use only lower case characters. To suppress the output you should use the keyword PAROUT and not the /dev/null path which is not allowed anymore.

If you want to write into the directory from where you are starting your CORSIKA run, you should give: DIRECT ./ or DIRECT ' ' or DIRECT " " (a blank enclosed in apostrophes or quotation marks). Please keep in mind that in FORTRAN an automatic expansion of UNIX names like '\$HOME' is not possible, rather you should give the fully expanded name of the directory ending with a / (slash) character.

Limit is: DSN must not begin with a ~ (tilde) character.

## 4.80 Table Output

PAROUT      FPAROUT      FTABOUT

Format = (A6, 2L) Defaults = T, F

FPAROUT : If .false., the particle output onto MPATAP is suppressed. This might be of advantage with the CERENKOV option to suppress the particle output file but keeping the Cherenkov output file (see Sect. 4.90).

FTABOUT : If .true., the tabular output of the charged particle development is written out to the file 'DATnnnnnnn.tab' (rsp. 'DATnnnnnnnnnn.tab' with option NRREXT, see Sect. 3.5.24 page 61) onto the output directory DSN (see Sect. 4.79 above).

## 4.81 Compact Output

COMOUT      COMOUT

Format = (A6, L), Default = T

COMOUT : If .true., the particle output is written in COMPACT form (see page 52). If .false., the standard CORSIKA particle output is written.

This keyword is only available in the COMPACT option.



## 4.82 Printer Output Unit

OUTPUT    MONNEW

Format = (A6, I), Default = 6

MONNEW : Logical unit of simulation control output on line printer. Make sure that your selection of MONNEW is not conflicting with existing definitions (see Table 1 page 26).

In the PARALLEL + PARALLELIB option the default value of MONNEW is changed to 89 for writing individual output files with names as the *cutfile* seeds, see Ref. [63].

## 4.83 Data Table Directory

DATDIR    DATDIR

Format = (A6, A132), Default = './'

DATDIR : Can be used to specify a common directory, where CORSIKA will try to find all required input data tables (except those belonging to the FLUKA versions). Lower case characters of DATDIR are not converted to capitals.

Limit is: DATDIR must not begin with a ~ (tilde) character.

## 4.84 Parameters for Parallel Treatment

PARALLEL    ECTCUT    ECTMAX    MPIID    FECTOUT

Format = (A8, 2F, I, L), Defaults = 1.E4, 1.E7, 1, F

ECTCUT : Threshold energy (in GeV) for subshowers. All particles with energies above ECTCUT will have a seed from the 6<sup>th</sup> sequence of random numbers (see keyword SEED page 75).

ECTMAX : Maximum energy (in GeV) for a complete subshower.

MPIID : Identification number of the parallel task which will be changed internally by MPI.

FECTOUT : If .true., particles with energy above ECTCUT are written additionally in an external *DATnnnnnnn-iiiiiii-kkkkkkkkkk.cut* file. This is needed when parallel runs are performed via shell scripting. These cut-files are written in binary mode. For the case of parallel runs via MPI treatment this is optional and could be used for debugging or for additional re-simulations of secondary subshowers to escape re-running the whole simulation from start/primary particle. Those cut-files are written in ASCII.

This keyword has to be specified before the keyword CUTFILE (see Sect. 4.85 below) is used.

This keyword is only available in the PARALLEL option <sup>84</sup>.

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<sup>84</sup>The MPI-Runner [63] reads this keyword and transmits the arguments to the subroutine *corsika*.

## 4.85 Reading of Cutted Particle File for Parallel Treatment

CUTFILE      CFILINP      I1CUTPAR      I2CUTPAR

Format = (A7, A255, 2I), Defaults = ' ', 0, 0

CFILINP : Input file name for cut particles to be read into second stack.

I1CUTPAR : Index of first particle to be used for the actual run.

I2CUTPAR : Index of last particle to be used for the actual run.

This keyword <sup>84</sup> has to be used when compilation is made with the PARALLEL option (see Sect. 3.5.26) to run parallel CORSIKA simulations distributed with the help of shell scripts. It must be used after the keyword PARALLEL (see Sect. 4.84 above) has been specified.

This keyword is only available in the PARALLEL option, but not interpreted by CORSIKA in the PARALLELIB option with MPI.

## 4.86 Parameters of Cutted Particles for Parallel Treatment

CUTPAR      CUTPAR(0...18)

Format = (A6, 19Z)

CUTPAR(0...18) : Parameters (in hex. format) of particle to be read into the second stack.

This keyword is only available in the PARALLEL option, but not in the PARALLELIB option.

## 4.87 Cherenkov Detector Array Definition

CERARY      NCERX      NCERY      DCERX      DCERY      ACERX      ACERY

Format = (A6, 2I, 4F), Defaults = 27, 27, 1500., 1500., 100., 100.

NCERX : Number of Cherenkov detectors in X-direction.

NCERY : Number of Cherenkov detectors in Y-direction.

DCERX : Grid spacing (in cm) of Cherenkov detectors in X-direction. The DCERX value has no relevance in case of NCERX = 1.

DCERY : Grid spacing (in cm) of Cherenkov detectors in Y-direction. The DCERY value has no relevance in case of NCERY = 1.

ACERX : Length (in cm) of each Cherenkov detector in X-direction.

ACERY : Length (in cm) of each Cherenkov detector in Y-direction.

The altitude of this array is at the lowest observation level. For the definition of the X and Y-directions see Fig. 1 (page 126) and keyword ARRANG (page 106).

Limits are: NCERX, NCERY  $\geq 1$  ; DCERX, DCERY, ACERX, ACERY  $\geq 1$ .

This keyword is only available in the CERENKOV option, but not in the IACT option for Cherenkov telescopes.

## 4.88 Cherenkov Wavelength Band

CWAVLG    WAVLGL    WAVLGU

Format = (A6, 2F), Defaults = 300., 450.

WAVLGL : Lower limit (in nm) of the wavelength band for Cherenkov radiation production.

WAVLGU : Upper limit (in nm) of the wavelength band for Cherenkov radiation production.

Limits are:  $100. \leq \text{WAVLGL} < \text{WAVLGU} \leq 700.$  (rsp. 2000 for the IACT option)

This keyword is only available in the CERENKOV, AUGCERLONG, or AUGERHIST options.

## 4.89 Cherenkov Bunch Size Definition

CERSIZ    CERSIZ

Format = (A6, F), Default = 0.

CERSIZ : Defines the maximal bunch size of Cherenkov photons that are treated together. If set to 0., by the subroutine *getbus* the program calculates a bunch size which is found to be appropriate for the HEGRA-array.

Limit is:  $\text{CERSIZ} \geq 0.$

This keyword is only available in the CERENKOV, AUGCERLONG, or AUGERHIST options.

## 4.90 Cherenkov Output Steering

CERFIL    MCERFI

Format = (A6, I), Default = 1 (rsp. 2 for CERWLEN option)

MCERFI : Steers the Cherenkov output to MPATAP or MCETAP.

MCERFI = 0: Output goes to MPATAP.

MCERFI 0 < MCERFI < 3: Output goes to MCETAP.

MCERFI  $\geq$  3: Output goes to MPATAP with name extension.

In the CERENKOV option (Cherenkov telescopes) with MCERFI > 0, the output file name DSN (specified by keyword DIRECT) should be set to `/dev/null`<sup>85</sup> to suppress the normal Cherenkov output file, as in the IACT option the Cherenkov telescope output will be written to the `eventio` output file<sup>86</sup>.

MCERFI = 1: is set automatically in the COMPACT option to prevent a writing of

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<sup>85</sup>The existence of `/dev/null` is assumed, see footnote of Sect. 4.79.

<sup>86</sup>Details on the `eventio` format may be found in the documentation supplied with the 'bernlhr' package.

Cherenkov photons to the COMPACT output.

MCERFI = 2: In the THIN option the last (8<sup>th</sup>) item of a bunch (see Table 11 page 141) gives the wavelength instead of the weight factor.

MCERFI ≥ 3: Acts like MCERFI = 2, but now the 7<sup>th</sup> item of each bunch gives the distance from the photon emission point to the detector array center instead of the photon emission height (See Table 11 page 141). Moreover the name convention of the Cherenkov data output files are changed to DATnnnnnnn.cher-teliii (see Sect. 3.4.3 page 45) (rsp. DATnnnnnnnnnn.cher-teliii with option NRREXT, see Sect. 3.5.24 page 61).

Limit is:  $0 \leq \text{MCERFI}$ .

This keyword is only available in the CERENKOV option.

## 4.91 Cherenkov Quantum Efficiency

CERQEF CERQEF CERATA CERMIR

Format = (A6, 3L), Defaults = F, F, F

CERQEF : If .true., quantum efficiency of detector photomultiplier is taken into account. It needs reading in the *quanteff.dat* file.

CERATA : If .true., the atmospheric absorption of Cherenkov photons is taken into account. It needs reading in the *atmabs.dat* file.

CERMIR : If .true., the mirror reflectivity of Cherenkov telescopes is taken into account. It needs reading in the *mirreff.dat* file.

Respecting these effects at an early stage of the Cherenkov photon simulation drastically reduces computing time and storage requirements for Cherenkov photon output. For the influence onto the longitudinal distribution of Cherenkov photons see Sect. 3.4.5 page 46 and keyword LONGI page 102.

This keyword is only available in the CEFFIC option together with the CERENKOV option.

## 4.92 Cherenkov Wavelength Range

CWAVRANG WLMIN WLMAX

Format = (A8, 2F), Defaults = 180., 700.

This keyword is only necessary for user supplied data files with a Cherenkov wavelength range different from the standard files delivered with CORSIKA (see Sect. 3.4.5 page 46).

WLMIN : Lower limit (in nm) of the wavelength band available in the *atmabs.dat*, *quanteff.dat*, and *mirreff.dat* data files.

WLMAX : Upper limit (in nm) of the wavelength band available in the *atmabs.dat*, *quanteff.dat*, and *mirreff.dat* data files.

Limits are:  $0. \leq \text{WLMIN} < \text{WLMAX} \leq 1000$ .

This keyword is only available in the CERENKOV, AUGCERLONG, or AUGERHIST

options.

### 4.93 Multiple Use of Cherenkov Events

CSCAT ICERML XSCATT YSCATT

Format = (A5, I, 2F), Defaults = 0, 0., 0.

ICERML : Number of uses of each event.

XSCATT : Maximum scattering of core location in  $\pm X$  direction (in cm). See Sect. 3.4.1 page 42 ff.

YSCATT : Maximum scattering of core location in  $\pm Y$  direction (in cm). See Sect. 3.4.1 page 42 ff.

Limits are:  $0 \leq \text{ICERML} \leq 20$  ;  $\text{XSCATT}, \text{YSCATT} \geq 0$ .

In case of Cherenkov telescopes ICERML telescope arrays are simulated randomly (see keyword SEED page 75) in the specified area which is a circle of radius XSCATT, if YSCATT = 0., or within a rectangle of area  $2 \text{ XSCATT} \cdot 2 \text{ YSCATT}$ .

This keyword is only available in the CERENKOV option.

### 4.94 Core Position of Scattered Cherenkov Event

COREPOS CERXOS(i) CERYOS(i)

Format = (A7, 2F), Defaults = 0., 0..

CERXPOS(i) :  $i^{\text{th}}$  position of scattered core in  $\pm X$  direction (in cm).

CERYPOS(i) :  $i^{\text{th}}$  position of scattered core in  $\pm Y$  direction (in cm).

If the keyword COREPOS appears before the CSCAT keyword, the scattering by the CSCAT input is disabled. IF this keyword appears after the CSCAT keyword the core positions defined by COREPOS are appended as additional scatter positions after the randomly selected core positions.

Limit:  $0 \leq i \leq 20$  core positions are possible.

This keyword is only available in the CERENKOV option.

### 4.95 Cherenkov Telescope Dimensions

Without IACT option

TELESCOPE X Y Z R ID

or with IACT option

TELESCOPE X Y Z R

Format = (A9, 4F, I)

X, Y, Z : Coordinates of Cherenkov telescope (in cm) relative to the center of the observation level. So at (0,0,0) means at the core of the shower at the observation level altitude.

R : Radius of sphere (in cm) within which the telescope is fully contained.

ID : Identification number of Cherenkov telescope. This number must be given without the IACT option. It is used to specify the extension of the output file name CERnnnnnnn-TELiii (rsp. CERnnnnnnnnnn-TELiii) where *iii* is the telescope ID (see Sect. 3.4.3 page 45). For ID = 0 the Cherenkov output file name extension '-TELiii' is omitted.

This keyword adds a new telescope at position X, Y, Z with radius R. At least one telescope has to be specified. It is recommended to define at least  $Z = R$  not to cut half of the sphere by the observation level. For the definition of the X and Y-directions see Fig. 1 (page 126) and keyword ARRANG (page 106).

Limits are:  $0 < R$ ;  $0 \leq \text{ID of telescope} \leq 999$ .

This keyword is only available in the CERENKOV rsp. IACT option.

## 4.96 Cherenkov Telescope Data File Name

TELFIL    TELFNM

Format = (A6, A512), Default = ' '

TELFNM : The telescope-specific data are to be written to a file named TELFNM in **eventio** format<sup>87</sup>. Lower case characters of TELFNM are not converted to capitals. If this file exists and is write-enabled, new data are appended. After ending the run the file will be set read-only to avoid accidental overwriting. The file name `/dev/null`<sup>88</sup> suppresses the output file.

If you want to write into the directory from where you are starting your CORSIKA run, you should give: `TELFIL ./` or `TELFIL ' '` or `TELFIL " "` (a blank enclosed in apostrophes or quotation marks). Please keep in mind that in FORTRAN an automatic expansion of UNIX names like '\$HOME' is not possible, rather you should give the fully expanded name of the directory ending with a / (slash) character.

Limit is: TELFNM must not begin with a ~ (tilde) character.

This keyword is only available in the CERENKOV option together with the IACT option for Cherenkov telescopes.

## 4.97 Trajectory Selection Flag

TRAFGL    TLOGIC

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<sup>87</sup>Details on the **eventio** format may be found in the documentation supplied with the 'bernlshr' package.

<sup>88</sup>The existence of `/dev/null` is assumed, see footnote of Sect. 4.79.

Format = (A6, L), Default = T

TLOGIC : If .true., the zenith and azimuth angles of each shower event are determined by the TRAJECT option to follow a source object movement in the sky. The determination of the angles by other options (VIEWCONE, VOLUMECORR, or VOLUMEDET) or by the keywords THETAP and PHIP is disregarded.

This keyword is only available in the TRAJECT option.

## 4.98 Source Position Parameters

SRCPOS RA DEC

Format = (A6, 2F), Defaults = 5.57, 22.0

RA : Defines the right ascension (in hours) of the simulated source in equatorial coordinates.

DEC : Defines the declination (in °) of the simulated source in equatorial coordinates.

The default values give the Crab Nebula.

Limits are:  $0. < RA < 24.$ ;  $-90. < DEC < +90.$  .

This keyword is only available in the TRAJECT option.

## 4.99 Trajectory Time Parameters

TRATM TYEAR TMONTH TDAY THOUR TMINUTE TSECOND TDURATION

Format = (A5, 7I), Defaults = 2000, 1, 1, 21, 0, 0, 3600

TYEAR, TMONTH, TDAY, THOUR, TMINUTE, TSECOND : Define the start time of the observation (in year, month, day, hour, min, sec).

TDURATION : Duration of observation (in sec).

Limits are: Only valid values for date and time are admitted;  $0 < TDURATON$ .

This keyword is only available in the TRAJECT option.

## 4.100 Lateral Telescope Site Parameters

TLAT TLATDGR TLATMIN TLATSEC TLATDIR

Format = (A4, 3F, A1), Defaults = 28., 45., 42.462, 'N'

TLATDGR : Latitude of the telescope site (in °).

TLATMIN : Latitude of the telescope site (in min).

TLATSEC : Latitude of the telescope site (in sec).

TLATDIR : Direction North = 'N', South = 'S' of the latitude of the telescope site.

The default values give the site of the MAGIC telescope.

Limits are:  $0. \leq TLATDGR \leq 90.$ ;  $0. \leq TLATMIN \leq 60.$ ;  $0. \leq TLATSEC \leq 60.$ ;

TLATDIR = 'N' or 'S'.

This keyword is only available in the TRAJECT option.

#### 4.101 Longitudinal Telescope Site Parameters

TLONG TLONGDGR TLONGMIN TLONGSEC TLONGDIR

Format = (A5, 3F, A1), Defaults = 17., 53., 26.525, 'W'

TLONGDGR : Longitude of the telescope site (in °).

TLONGMIN : Longitude of the telescope site (in min).

TLONGSEC : Longitude of the telescope site (in sec).

TLONGDIR : Direction East = 'E', West = 'W' of the longitude of the telescope site.

The default values give the site of the MAGIC telescope.

Limits are:  $0. \leq \text{TLONGDGR} \leq 180.$ ;  $0. \leq \text{TLONGMIN} \leq 60.$ ;  $0. \leq \text{TLONGSEC} \leq 60.$ ; TLONGDIR = 'E' or 'W'.

This keyword is only available in the TRAJECT option.

#### 4.102 Geomagnetic Declination of Telescope

GEODEC GEODECL

Format = (A6, F), Default = -6.35

GEODECL : Defines the geomagnetic declination<sup>89</sup>, i.e. the directional deviation of the magnetic North from the geographic North (in °).

The default value corresponds with the site of the MAGIC telescope.

Limit is:  $-45. \leq \text{GEODECL} \leq +45.$  .

This keyword is only available in the TRAJECT option.

#### 4.103 Trajectory Broadening Parameter

TRARAD TRAD

Format = (A6, F), Default = 0.

TRAD : Defines the radius of a spread around the calculated trajectory (in arcmin).

Limit is:  $0. \leq \text{TRAD} \leq 3600.$  .

This keyword is only available in the TRAJECT option.

---

<sup>89</sup>The sign of the declination is defined positive for eastward declination, negative for westward declination, see also Ref. [73]



#### 4.104 Interesting Energy Threshold

EINTER ENERGY\_INTER

Format = (A6, F), Default = 1000.

ENERGY\_INTER : Threshold energy (in GeV) above which particles are interesting for production of high-energy neutrinos for the IceCube [59] experiment.

Limit is:  $0 \leq \text{ENERGY\_INTER}$ .

This keyword is only available in the ICECUBE1 option.

#### 4.105 PIPE Output Redirection Flag

PIPE PIPE\_OUTPUT

Format = (A4, L), Default = F

PIPE\_OUTPUT : If .true., the output is redirected to a pipe-buffer. This buffer should serve as on-line input to a detector simulation program.

This keyword is only available in the ICECUBE2 option.

#### 4.106 Compress Output Flag

COMPRESS GZIP\_OUTPUT

Format = (A8, L), Default = T

GZIP\_OUTPUT : If .true., the output is compressed using the gzip procedure to reduce storage space. The gain in space reduction will not be very dramatic.

This keyword is only available in the ICECUBE2 option.

#### 4.107 Size of Stack

DYNSTACK N1 N2 ... NX

Format = (A8, A500), Default = "

Forwards a variable number of integer numbers from the steering card to the DYNSTACK setup routine. In normal use cases this should contain the size of the stack.

This keyword is only available in the DYNSTACK option.

#### 4.108 Parameter to Stack

DYNSTACK\_P A500

Format = (A10, A500), Default = ”

Forwards a variable number of characters, separated at whitespaces, to the DYNSTACK setup routine. This call should be used to set any stack-specific settings used in the current simulation.

This keyword is only available in the DYNSTACK option.

#### 4.109 Remote IP

REMOTECONTROL\_IP A20

Format = (A16, A20), Default = ”

Forwards a IP from the external server to connect the network thread to. The format equals the default IP format, e.g. 127.0.0.1:1337, where the part subsequent to the colon denotes the port number.

This keyword is only available in the REMOTECONTROL option.

#### 4.110 Parameter to Remotecontrol

REMOTECONTROL\_P A500

Format = (A10, A500), Default = ”

Forwards a variable number of characters, separated at whitespaces, to the REMOTECONTROL setup routine. This call should be used to set any connection-specific settings used in the current simulation.

This keyword is only available in the REMOTECONTROL option.

#### 4.111 Write Data Base File

DATBAS FDBASE

Format = (A6, L), Default = F

FDBASE : If .true., all essential run parameters are written to the file '*DATnnnnnnn.dbase*' (rsp. '*DATnnnnnnn.info*' or '*DATnnnnnnnnnn.info*' in the AUGERINFO option) onto the output directory DSN (keyword DIRECT page 108). This file may be used to build a data base for examining the content of an air shower library (page 151).

This keyword is only effective in the UNIX options.

#### 4.112 User Name

USER USER

Format = (A4, A60), Defaults = ' '

USER : A user name is read in to be written to the 'DATnnnnnnn.dbase' file. If the NRREXT option is used the file name gets 'DATnnnnnnnnnn.dbase' (see Sect. 3.5.24 page 61). Lower case characters of USER are not converted to capitals.

#### 4.113 Host Name

HOST HOST

Format = (A4, A60), Defaults = ' '

HOST : A host name is read in to be written to the 'DATnnnnnnn.dbase' file. If the NRREXT option is used the file name gets 'DATnnnnnnnnnn.dbase' (see Sect. 3.5.24 page 61). Lower case characters of HOST are not converted to capitals.

#### 4.114 Debugging

DEBUG DEBUG MDEBUG DEBDEL NDEBDL

Format = (A5, L, I, L, I), Defaults = F, MONIOU, F, 100000

DEBUG : If .false., debugging is disabled. If .true., additional output for debugging purposes is given on logical unit MDEBUG.

MDEBUG : Logical unit where to write debugging information. Make sure that your selection of MDEBUG is not conflicting with existing definitions (see Table 1 page 26).

DEBDEL : If .true., the debugging printouts are activated after NDEBDL particles above the ECTMAP energy have been printed. If .false., delayed debugging is disabled. This feature helps to trace run time errors that have occurred in long simulation runs.

NDEBDL : See DEBDEL

Limit is:  $0 \leq \text{NDEBDL}$ .

#### 4.115 Debugging EGS

EGSDEB JCLOCK

Format = (A6, I), Default = 2147483647

JCLOCK : Counter for delayed start of EGS4 debugging. After activation of debug by DEBUG or by NDEBDL (see Sect. 4.114 above) each pass of subroutine *electr* or *photon* is counted. If the counter exceeds JCLOCK, the debug statements within the EGS4 portion are activated. This output appears on the unit MDEBUG.

Limit is:  $0 \leq \text{JCLOCK}$ .

#### 4.116 FLUKA Printing

FLUDBG      FFLUDB

Format = (A6, L), Default = F

FFLUDB : If .true. the two files '*DATnnnnnnn.flout*' for additional information on the parameters used by FLUKA and '*DATnnnnnnn.flerr*' on possible FLUKA error messages are written onto the output directory DSN (keyword DIRECT page 108). In case of the NRREXT option these files get the names '*DATnnnnnnnnnn.flout*' resp. '*DATnnnnnnnnnn.flerr*'. If .false. in the LINUX option the two files are written to 'SCRATCH' i.e. they are written to a temporary file which at the end of the program will be deleted. Without the LINUX option the files are opened directly to the directory */dev/null*<sup>90</sup>.

This keyword is only available in the FLUKA options.

#### 4.117 GHEISHA Debugging

GHEIDB      GHEISDB

Format = (A6, L), Default = F

GHEISDB : If .true., in the DEBUG case also the GHEISHA routines produce debug output. This output appears on the unit MDEBUG.

This keyword is only available in the GHEISHA option.

#### 4.118 URQMD Debugging

URQMD      FURQMD      IUDEBUG

Format = (A5, L, I), Default = T, 0

FURQMD : If .true., the UrQMD routines are used for the low-energy hadronic interactions. If .false., the program will stop.

IUDEBUG : If  $> 0$ , in the DEBUG case also the routines of UrQMD produce some output. With increasing value of IUDEBUG this printout becomes more and more detailed. This output appears on the unit MDEBUG.

Limit is:  $0 \leq \text{IUDEBUG} \leq 3$ .

This keyword is only available in the URQMD option.

---

<sup>90</sup>The existence of */dev/null* is assumed, see footnote of Sect. 4.79 page 108.

## 4.119 PYTHIA Printing

PYTHIA IFLGPYW IFLGPYE

Format = (A6, 2I), Default = 0, 0

IFLGPYW : If set  $> 0$  the printing of PYTHIA warnings is enabled. IFLGPYW gives the number of warnings after which PYTHIA becomes silent (see MSTU(26) of Pythia manual [34]).

IFLGPYE : If set  $> 0$  the printing of PYTHIA errors is enabled. IFLGPYE gives the number of errors after which PYTHIA becomes silent (see MSTU(22) of Pythia manual [34]).

Limits are:  $0 \leq \text{IFLGPYW}$ ;  $0 \leq \text{IFLGPYE}$ .

This keyword is only available in the DPMJET, CHARM, or TAULEP options.

## 4.120 Cherenkov Debugging

CDEBUG LCERDB

Format = (L), Default = F

LCERDB : If .false., Cherenkov debug output is disabled. If .true., the Cherenkov routines produce debug output. This output appears on the unit MDEBUG.

This keyword is only available in the CERENKOV option.

## 4.121 Interaction Test Target Definition

INTTST ITTAR MCM

Format = (A6, 2I), Defaults = 0, 0

ITTAR : Defines the target for the interaction test option<sup>91</sup>: 1 = proton; 2 = neutron; 9 = Beryllium; 12 = Carbon; 14 = Nitrogen; 16 = Oxygen; 40 = Argon; 99 = air.

MCM : Defines the reference system for which the interaction products are plotted. 1 = rest system of 1 target nucleon and 1 projectile nucleon<sup>92</sup>; 2 = laboratory system; 3 = rest system of all secondary particles (but not the spectators).

This keyword is only available in the INTTEST option.

## 4.122 Interaction Test Decay

INTDEC LPIO LETA LHYP LKOS

---

<sup>91</sup>For the DPMJET high-energy model and  $\text{ITTAR} < 14$  the Glauber data file *glaubint.glb* has to be used instead of the *glaubtar.glb* file.

<sup>92</sup>In photo-nuclear interactions the cm-system of the photon with one nucleon.

Format = (A6, 4L), Defaults = T, T, T, T

LPI0 : If .true. the  $\pi^0$  particles decay before gathering them in the interaction test.

LETA : If .true. the  $\eta$  particles decay before gathering them in the interaction test.

LHYP : If .true. all hyperons decay before gathering them in the interaction test.

LK0S : If .true. the  $K_s^0$  particles decay before gathering them in the interaction test.

This keyword is only available in the INTTEST option.

#### 4.123 Interaction Test Spectator Definition

INTSPC LSPEC

Format = (A6, L), Default = F

LSPEC : If .true. spectators are plotted, if .false. spectators are not plotted in the interaction test.

This keyword is only available in the INTTEST option.

#### 4.124 Interaction Test Diffraction Flag

DIFOFF NDIF

Format = (A6, I), Default = 0

NDIF : Allows to select diffractive or non-diffractive interactions only<sup>93</sup>. 0 = diffractive and non-diffractive interactions mixed; 1 = non-diffractive interactions only; 2 = diffractive interactions only. With the QGSJET01d interaction model NDIF = 2 cannot be selected. With the EPOS, NEXUS and QGSJET-II models only NDIF = 0 is possible.

This keyword is only available in the INTTEST option.

#### 4.125 Interaction Test Trigger Condition

TRIGGER NTRIG

Format = (A7, I), Default = 0

NTRIG : Allows to select various trigger conditions for comparison with experimental data:

0 = accepts all events;

1 = accepts only events according to the UA5-experiment [77] trigger;

2 = accepts only events according to the CDF-experiment [78] trigger.

3 = accepts only events according to the P238-experiment [79] trigger.

---

<sup>93</sup>For photo-nuclear interactions the meaning is: 0 = multi-hadron production and vector meson production mixed; 1 = no vector meson production, only multi-hadron production; 2 = only vector meson production, no multi-hadron production.

NTRIG  $\neq$  0 may be combined only with NDIF = 0 .  
This keyword is only available in the INTTEST option.

#### 4.126 Interaction Test Histogram Output

HISTDS HISTDS

Format = (A6, A120), Defaults = 'histo.corsika.inttest'

HISTDS : May be used to specify a name of the histogram output directory and data file. Lower case characters of HISTDS are not converted to capitals. Do not use capitals as the HBOOK routines use only lower case characters. The data file name is extended by a string containing information about projectile, target, energy, and the type of interaction which has been selected. At the end of the data file name .hbook is appended such that the total data file name would look like

histo.corsika.inttest.p0014t14e100e3.diffractive.hbook for a proton projectile on nitrogen target with a lab energy of 100E3 GeV including diffractive events.

If you want to write into the directory from where you are starting your CORSIKA run, you should give: HISTDS ./ or HISTDS ' ' or HISTDS " " (a blank enclosed in apostrophes or quotation marks). Please keep in mind that in FORTRAN an automatic expansion of UNIX names like '\$HOME' is not possible, rather you should give the fully expanded name of the directory ending with a / (slash) character.

Limit is: HISTDS must not begin with a ~ (tilde) character.  
This keyword is only available in the INTTEST option.

#### 4.127 Plot Output

PLOTSH PLOTSH

Format = (A6, L), Default = F

PLOTSH : If .true., the track start- and endpoints of the electromagnetic, muonic, and hadronic component of the shower are given out separately and may be used to plot the shower development.

This keyword is only available in the PLOTSH and PLOTSH2 option.

#### 4.128 Plot Axes Definition

PLAXES X1 X2 Y1 Y2 Z1 Z2

Format=(A6,6F), Defaults = -500000., 500000., -500000., 500000., 0., 3000000.

X1, X2 : They denote the X-axis range (in cm) to be plotted in the map.

Y1, Y2 : They denote the Y-axis range (in cm) to be plotted in the map.

Z1, Z2 : They denote the Z-axis range (in cm) to be plotted in the map.

The point of first interaction determines the zero point of the X and Y axes (see Fig. 1 page 126). Depending on the choice of these parameters, the whole shower may be visualized, or one can 'zoom in' on interesting regions of the shower.

Limits are:  $X1 < X2$ ,  $Y1 < Y2$ ,  $Z1 < Z2$  .

This keyword is only available in the PLOTSH2 option.

#### 4.129 Plot Energy Cut Definition

PLCUTS    ELCUTS(1...4)    TCUT    FBOXCUT

Format=(A6,5F,L), Defaults 0.3, 0.3, 0.003, 0.003, 100000., F

ELCUTS(1...4) : ELCUTS denote the energy cuts in the same order as those for the keyword ECUTS (hadrons, muons, electrons, photons) (see page 100).

TCUT : This is an upper bound on the time (in ns) passed since the first interaction. If, at the end point of a track, the time is above TCUT, the track is **not** plotted. This cut allows a visualization of the shower development.

FBOXCUT : This flag determines whether only track segments inside the three-dimensional box given by the axis ranges should be plotted. If .true., all track segments whose start- and endpoints both fall outside this box are not plotted.

This keyword is only available in the PLOTSH2 option.

#### 4.130 End of Steering

EXIT

Format = (A4)

This keyword ends the keyword input.



## 5 Input Example

The keyword input file for a CORSIKA run including QGSJET and CERENKOV options may look like the following list.

```
RUNNR 1                number of run
EVTNR 100400           no of first shower event
SEED 100401 0 0        seed for hadronic part
SEED 100402 0 0        seed for EGS4 part
SEED 100403 0 0        seed for Cherenkov part
NSHOW 10               no of showers to simulate
PRMPAR 5626            primary particle code (iron)
ERANGE 2.00E4 4.00E4   energy range of primary (GeV)
ESLOPE -2.7            slope of energy spectrum
THETAP 0. 10.          range zenith angle (deg)
PHIP -180. 180.        range azimuth angle (deg)
QGSJET T 0             QGSJET for high energy & debug level
QGS SIG T             QGSJET cross-sections enabled
HADFLG 0 0 0 0 0 2    HDPM interact.flags & fragmentation flag
ELMFLG T T            elmag. interaction flags NKG, EGS4
STEPFC 1.             multiple scattering step length factor
RADNKG 200.E2          outer radius (cm) of NKG elect. distrib.
MAGNET 20.4 43.23      magnetic field central Europe (/uT)
ECUTS .3 .3 .015 .015 energy cuts: hadr. muon elec. phot. (GeV)
LONGI T 20. T T        longitud, stepsize(g/cm^2), fit, out
MUMULT T              muon multiple scattering by Moliere
MUADDI T              additional muon information
OBSLEV 110.E2          observation level (cm)
ARRANG 18.25           angle between north to array-grid (deg)
MAXPRT 10             max. no of printed events
ECTMAP 1.E2           printout gamma factor cut
DIRECT /home/user/corsika/run/ directory of particle output
CERARY 10 8 1200. 1500. 80. 50. Cherenkov detector grid (cm)
CWAVLG 300. 450.      Cherenkov wavelength band (nm)
CERSIZ 5.             bunch size Cherenkov photons
CERFIL 1              Cherenkov output file
CSCAT 5 1000. 1000.   scatter Cherenkov events (cm)
DATBAS T              write data base file
USER your              user name for data base file
HOST your_host        host name for data base file
DEBUG F 6 F 999999999 debug flag, log. unit, delayed debug
EXIT
```

## 6 Coordinate System

The coordinates in CORSIKA are defined with respect to a Cartesian coordinate system with the positive  $X$ -axis pointing to the magnetic North, the positive  $Y$ -axis to the West, and the  $Z$ -axis upwards. The origin is located at sea level and at the vertical of the core position, such that the without particular option the particles in the output file are distributed around the point ( $x = 0, y = 0$ ) for each observation levels at a given altitude.

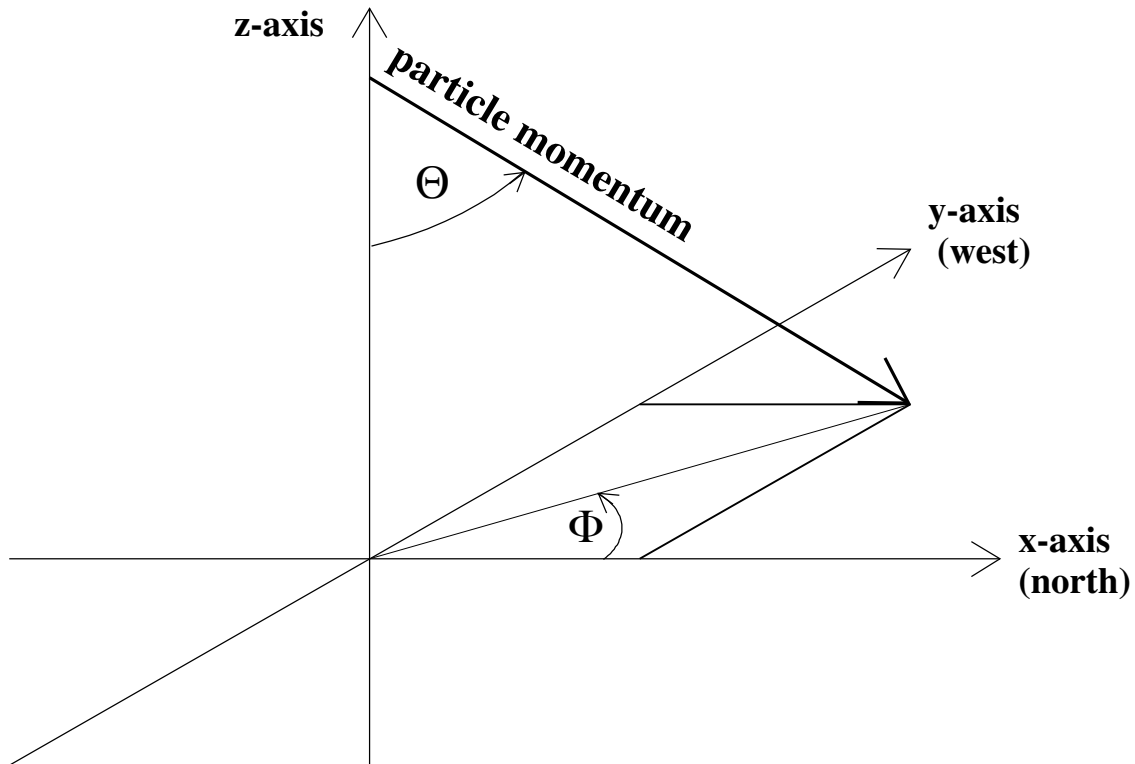


Figure 1: Coordinate system in CORSIKA.

This definition is necessary, because the Earth's magnetic field is taken into account. By default the magnetic field is implemented for the location of Karlsruhe (49° N, 8° E) as described at the keyword MAGNET (page 84). The zenith angle  $\Theta$  of a particle trajectory is measured between the particle momentum vector and the negative  $Z$ -axis, and the azimuthal angle  $\Phi$  between the positive  $X$ -axis and the horizontal component of the particle momentum vector (i.e. with respect to North) proceeding counterclockwise. This is shown in Fig. 1.

In the EGS4 part and in CONEX the definitions of the coordinate system differ from that used in CORSIKA (see the in-line comments in these program parts). For EGS4 the positive  $Z$ -axis points downwards, the positive  $Y$ -axis points to East, the  $X$ -axis remains unchanged. For CONEX the positive  $X$ -axis points to East, the positive  $Y$ -axis points to North, the  $Z$ -axis remains unchanged. But this concerns only internal definition, selecting CONEX do not change the coordinate system of CORSIKA experienced by the user.

In the case of upward going shower ( $\Theta > 90^\circ$ ), the definition of the angles remains the same ( $\Phi = 0$  means toward the magnetic North) but the origin of the axis will depend

whether UPWARD or UPWARDOLD options are used :

- In the case of UPWARDOLD (Fig. 2 a), the observation level is at the END of the shower propagation. It means that the shower direction is from the sea level toward the origin of the frame. When FIXHEI increase, the starting point of the shower get closer to the (0,0) point.
- In the case of UPWARD (Fig. 2 b), the observation level is at the BEGINNING of the shower propagation. It means that the shower direction is from the origin of the frame toward the top of the atmosphere. When FIXHEI increase, the starting point of the shower get further away from the (0,0) point.

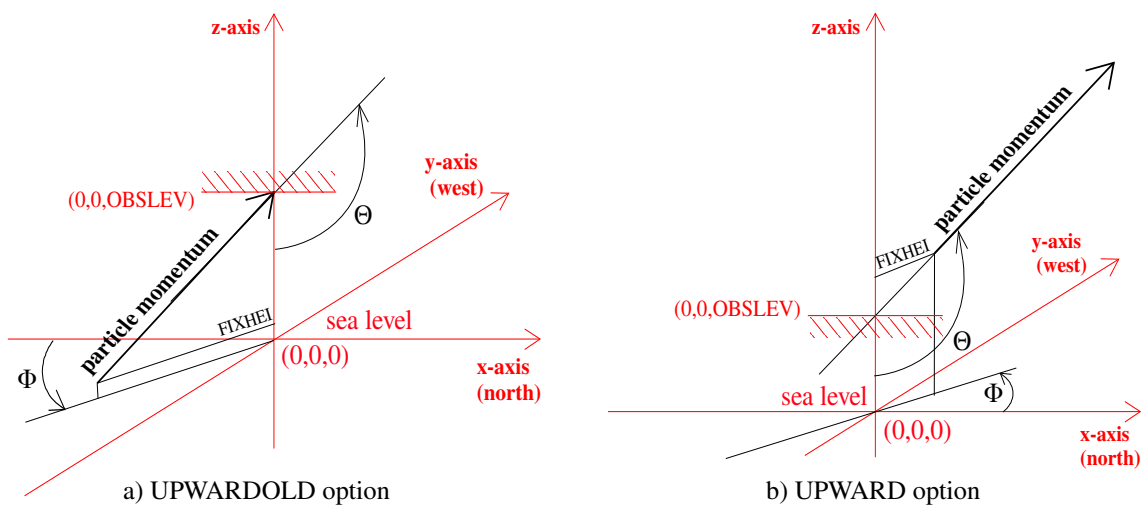


Figure 2: Coordinate system in CORSIKA for upward goind showers.

Since the origin of the frame is used to define the position of the inclined observation plane when the INCLINED option is used (see Sect. 10.3 page 148), it is important to take it into account when choosing between these 2 options.

**Attention:** The definition of the coordinate system in CORSIKA **disagrees** from definitions of other air-shower simulation programs (e.g. AIRES [80]) and from the conventions of the Auger experiment !!

## 7 Units in CORSIKA

Within CORSIKA uniform units for the various dimensions are used as far as possible. But there are deviations at that program parts which are coupled to CORSIKA. Table 3 gives an overview on the used units.

Quantity	CORSIKA	EGS4	FLUKA GHEISHA UrQMD HERWIG PYTHIA	SIBYLL	DPMJET EPOS NEXUS QGSJET VENUS	CONEX
length	cm <sup>1)</sup>	cm				m
energy	GeV	MeV	GeV	GeV <sup>4)</sup>	GeV	GeV
mass	GeV	MeV	GeV	GeV	GeV	GeV
time	sec <sup>2)</sup>	sec				sec
magn. field	$\mu\text{T}$					<sup>5)</sup>
density	$\text{g}/\text{cm}^3$	$\text{g}/\text{cm}^3$				$\text{g}/\text{cm}^3$
mass overburden	$\text{g}/\text{cm}^2$					$\text{g}/\text{cm}^2$
angle	rad <sup>3)</sup>	rad				rad
wavelength	nm					
	<sup>1)</sup> In some subroutines also m is used. <sup>2)</sup> For output files also nsec is used. <sup>3)</sup> For in- and output files also $^\circ$ is used. <sup>4)</sup> In some subroutines also TeV is used. <sup>5)</sup> No Earth magnetic field considered.					

Table 3: Units used in CORSIKA and the coupled programs.

## 8 Particles in CORSIKA

CORSIKA is able to treat the particles that are listed in Table 4. The particle codes have in general been chosen according to the convention in the GEANT detector simulation code [36]. Exceptions are vector mesons ( $\omega$  and  $\rho$ ) and resonances ( $K^*$  and  $\Delta$ ), the  $\eta$  particles which are split in 4 types according to their decays in the HDPM routines, the different types of neutrinos (to be generated in the NEUTRINO option), the nuclei, and the Cherenkov photons. The codes of the charmed particles correspond with those of DPMJET. The particles with codes between 116 and 173 are only available in the CHARM or TAULEP option. Cherenkov photons can not be a primary particle for an air shower simulation.

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<sup>94</sup>Only available in MUPROD option.

Particle identifications			
Identification	Particle	Identification	Particle
1	$\gamma$	50	$\omega$
2	$e^+$	51	$\rho^0$
3	$e^-$	52	$\rho^+$
		53	$\rho^-$
5	$\mu^+$	54	$\Delta^{++}$
6	$\mu^-$	55	$\Delta^+$
7	$\pi^0$	56	$\Delta^0$
8	$\pi^+$	57	$\Delta^-$
9	$\pi^-$	58	$\overline{\Delta}^{--}$
10	$K_L^0$	59	$\overline{\Delta}^-$
11	$K^+$	60	$\overline{\Delta}^0$
12	$K^-$	61	$\overline{\Delta}^+$
13	$n$	62	$K^{*0}$
14	$p$	63	$K^{*+}$
15	$\bar{p}$	64	$\overline{K}^{*-}$
16	$K_S^0$	65	$\overline{K}^{*0}$
17	$\eta$	66	$\nu_e$
18	$\Lambda$	67	$\bar{\nu}_e$
19	$\Sigma^+$	68	$\nu_\mu$
20	$\Sigma^0$	69	$\bar{\nu}_\mu$
21	$\Sigma^-$		
22	$\Xi^0$	71	$\eta \rightarrow \gamma\gamma$
23	$\Xi^-$	72	$\eta \rightarrow 3\pi^0$
24	$\Omega^-$	73	$\eta \rightarrow \pi^+\pi^-\pi^0$
25	$\bar{n}$	74	$\eta \rightarrow \pi^+\pi^-\gamma$
26	$\bar{\Lambda}$	75	$\mu^+$ add. info.
27	$\overline{\Sigma}^-$	76	$\mu^-$ add. info.
28	$\overline{\Sigma}^0$		
29	$\overline{\Sigma}^+$	85	decaying $\mu^+$ at start <sup>94</sup>
30	$\overline{\Xi}^0$	86	decaying $\mu^-$ at start <sup>94</sup>
31	$\overline{\Xi}^+$		
32	$\overline{\Omega}^+$	95	decaying $\mu^+$ at end <sup>94</sup>
48	$\eta'$	96	decaying $\mu^-$ at end <sup>94</sup>
49	$\phi$		

Table 4: Particle identifications as used in CORSIKA (to be continued).

Particle identifications (continued)			
Identification	Particle	Identification	Particle
116	$D^0$	155	$\overline{\Xi}'_c{}^-$
117	$D^+$	156	$\overline{\Xi}'_c{}^0$
118	$\overline{D}^-$	157	$\overline{\Omega}_c{}^0$
119	$\overline{D}^0$		
120	$D_s^+$	161	$\Sigma_c^{*++}$
121	$\overline{D}_s^-$	162	$\Sigma_c^{*+}$
122	$\eta_c$	163	$\Sigma_c^{*0}$
123	$D^{*0}$		
124	$D^{*+}$	171	$\overline{\Sigma}_c^{*-}$
125	$\overline{D}^{*-}$	172	$\overline{\Sigma}_c^{*0}$
126	$\overline{D}^{*0}$	173	$\overline{\Sigma}_c^{*+}$
127	$D_s^{*+}$		
128	$\overline{D}_s^{*-}$	176	$B^0$
		177	$B^+$
130	$J/\psi$	178	$\overline{B}^-$
131	$\tau^+$	179	$\overline{B}^0$
132	$\tau^-$	180	$B_s^0$
133	$\nu_\tau$	181	$\overline{B}_s^0$
134	$\overline{\nu}_\tau$	182	$B_c^+$
		183	$\overline{B}_c^-$
137	$\Lambda_c^+$	184	$\Lambda_b^0$
138	$\Xi_c^+$	185	$\Sigma_b^-$
139	$\Xi_c^0$	186	$\Sigma_b^+$
140	$\Sigma_c^{++}$	187	$\Xi_b^0$
141	$\Sigma_c^+$	188	$\Xi_b^-$
142	$\Sigma_c^0$	189	$\Omega_b^-$
143	$\Xi_c'^+$	190	$\overline{\Lambda}_b^0$
144	$\Xi_c'^0$	191	$\overline{\Sigma}_b^+$
145	$\Omega_c^0$	192	$\overline{\Sigma}_b^-$
		193	$\overline{\Xi}_b^0$
149	$\overline{\Lambda}_c^-$	194	$\overline{\Xi}_b^+$
150	$\overline{\Xi}_c^-$	195	$\overline{\Omega}_b^+$
151	$\overline{\Xi}_c^0$		
152	$\overline{\Sigma}_c^{*-}$		
153	$\overline{\Sigma}_c^-$		
154	$\overline{\Sigma}_c^0$		
$A \times 100 + Z$	nucleus of Z protons and A – Z neutrons ( $2 \leq A \leq 56$ )		
8888jjj	weights of preceding particle (MULTITHIN option)		
9900	Cherenkov photons on particle output file		

Table 4: (continued) Particle identifications as used in CORSIKA.

## 9 Running the CORSIKA Program

Depending on the program options CORSIKA needs at minimum  $\approx 96$  Mbyte memory for QGSJET01 with GHEISHA without other options. Other high-energy or low-energy models need additionally memory. Especially EPOS, QGSJET-III, SIBYLL, and FLUKA need a huge amount of memory. Empirical (approximated) values of required memory for 64-bit LINUX computers are given in Table 5.

Option	QGSJET01 + GHEISHA	DPMJET	EPOS	QGSJET-II	SIBYLL	FLUKA	UrQMD	CONEX
approx. Memory (Mbyte)	96	+ ?	+ 551	+ 267	+ 4957	+ 431	+ 5	+ 657

Table 5: Required memory for various CORSIKA options (64-bit LINUX).

A simplified flow diagram of CORSIKA is given in Appendix B (page 165).

The sequence of the initializing procedures is given in Appendix C (page 166).

CORSIKA runs fastest (full simulation adopted without THINning) when using no EGS4, no DPMJET, no EPOS, no NEXUS, no VENUS, no FLUKA, no UrQMD, no NEUTRINO, no Cherenkov light generation, no CURVED, and no options producing plots, histograms or ROOT files. For this program version the computing time on an Intel 6600 Core2Duo (with 2.4 GHz) is  $\approx 7.3$  sec per shower for primary protons of energy  $10^6$  GeV, vertical incidence, NKG enabled, with one observation level at 110 m a.s.l. and with the hadron and muon energy cut at 0.3 GeV. Under the same conditions an iron induced shower consumes  $\approx 10.6$  sec. The computing time scales roughly with the primary energy. The full EGS4 option with longitudinal profile is roughly 200 times slower than the fastest version mentioned above (with ELCUT(3,4) at 3 MeV). There is not much experience what the time consumption is for the CERENKOV option, but it will be much higher. QGSJET01 run fastest in combinatin with GHEISHA. EPOS or NEXUS (NKG, no EGS) need  $\approx 50$  times the computing time of QGSJET01. FLUKA needs  $\approx 7$  times computing time of GHEISHA (+ QGSJET01), UrQMD needs  $\approx 40$  times that of GHEISHA.

In case of THINning (Sect. 3.5.36 page 68) the computing time strongly depends on the energy fraction below which thinning becomes active, see Table 2 page 69.

The interested CORSIKA user may find CPU-times for various interaction models and options in Ref. [30].

The particle output format is described in Sect. 10.2 (page 136 ff.). For each particle that reaches an observation level 7 words (for THINning 8 words) with 4 bytes each are stored on the particle output file. Proton showers at  $10^6$  GeV deliver at sea level roughly 0.8 MB particle output when calculated with the NKG option ( $E_h, E_\mu > 0.3$  GeV). With EGS4 about 30 MB output are written due to the huge amount of photons and electrons that

are explicitly stored ( $E_e, E_\gamma > 0.003$  GeV)<sup>95</sup>. The particle output in the CERENKOV option is additionally increased as the Cherenkov photons are stored, either separately on the Cherenkov output file, or together with the particles on the particle output file.

While running, the interactions produce plenty of secondary particles which are stored in an intermediate stack which is foreseen for 200000 particles. After completing an interaction all particles (in case of THINning only a fraction of them) are moved to the internal stack. The size of this internal stack is large enough to keep 680 (in case of thinning 624) particles. If the size of the internal stack is full, 680 (624) particles are temporarily added to the external stack. If the internal stack is empty 340 (312) particles are re-read from the external stack. Only half of the internal stack is read back to avoid a permanent shifting of data from and to the internal stack if it is just full and the number of secondaries oscillates with a small amplitude around 680 (624).

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<sup>95</sup>These energy cuts correspond with the sensitivities of the KASCADE array detectors.



## 10 Outputs

There are two major output files produced by a simulation run. The control printout (.lst file) allows to survey the simulation run. The particle output file is written to the data file '*DATnnnnnn*' (rsp. '*DATnnnnnnnnnn*' if the NRREXT option is selected, see Sect. 3.5.24 page 61) with *nnnnnn* (rsp. *nnnnnnnnnn*) being the run number specified in the keyword RUNNR (page 74). This file becomes very large when simulating showers in great detail (EGS4, low thresholds, ...). In the CERENKOV version an additional file '*CERnnnnnn*' (rsp. '*CERnnnnnnnnnn*') might be written. Optionally a tabular output ( $\gamma$ ,  $e^\pm$ , and  $\mu^\pm$  particles at ground and longitudinal development of charged particles) is written out to the file '*DATnnnnnn.tab*' (rsp. '*DATnnnnnnnnnn.tab*'). Further on the longitudinal distributions of particle numbers and energy deposits may be written to file '*DATnnnnnn.long*' (rsp. '*DATnnnnnnnnnn.long*'). The output file '*DATnnnnnn.dbase*' (rsp. '*DATnnnnnn.info*' or '*DATnnnnnnnnnn.dbase*' or '*DATnnnnnnnnnn.info*') is destined to be used in a data base for examining the content of an air shower library.

If the INCLINED option (Sect. 3.5.17 page 57) is used in combination with the INCLIN keyword (page 104) an output file '*DATnnnnnn.inclined*' (rsp. '*DATnnnnnnnnnn.inclined*') is written, which has the same binary structure as a standard CORSIKA particle output file '*DATnnnnnn*' (rsp. '*DATnnnnnnnnnn*') see Sect. 10.3 (page 148).

The INCLINED option internally makes use of the COAST package.

### 10.1 Control Printout (.lst File Output)

The simulation run produces a printout (.lst file, unit MONIOU, by default standard output)<sup>96</sup> that allows to control the simulation and informs about the general run, the program version with interaction model, the selected options, steering keywords, physical constants, the used atmospheric model, and the primary particle (about 200 lines). For each shower it prints roughly 400 lines containing the random number generator status, time at beginning of a shower, the primary particle at the place of the first interaction, the number of secondaries reaching the observation levels with energies above ECTMAP (page 107), the stack statistics, internal and external stack usage, energy-multiplicity and energy-elasticity relations, interaction statistics for nucleons, pions, kaons, and strange baryons per kinetic energy interval, an interaction length statistics for the above particles and a decay statistics for muons, summaries of secondaries for each observation level, NKG electromagnetic shower information, and the longitudinal shower development.

The NKG output (keyword ELMFLG, page 92) comprises a table on the longitudinal development of the electromagnetic shower component giving every 100 g/cm<sup>2</sup> the number of electrons and the longitudinal pseudo-age parameter<sup>97</sup>. For the lateral electron distri-

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<sup>96</sup>Renaming the standard (log) output to '*DATnnnnnn.lst*' (rsp. '*DATnnnnnnnnnn.lst*') and redirecting it to the directory specified by the keyword DIRECT (page 108) is convenient (page 134) as by the shell commands 'dir' or 'ls -l' all files belonging to one run are displayed consecutively which facilitates book-keeping.

<sup>97</sup>See footnote page 41.

bution the densities (in *electrons/cm<sup>2</sup>*) are calculated on a grid of 80 points (8 directions separated by 45° with 10 distances between 1 m and RADNKG (page 93) for each direction) and the lateral pseudo-ages for those various distances are determined. The lateral distribution is calculated only for the lowest two observation levels.

For the longitudinal development (keyword LONGI page 101) three tables are generated:

- The first one gives the particle numbers of  $\gamma, e^+, e^-, \mu^+, \mu^-$ , hadrons, all charged, nuclei, neutrinos, and Cherenkov photons as function of atmospheric depth.
- The second one reflects the energy content within the various particle species  $\gamma, e^+, e^-, \mu^+, \mu^-$ , hadrons, charged particles, nuclei, neutrinos, and the energy sum as function of atmospheric depth. For all particle species the particle rest mass is included within the energy.
- The third table gives the energy dissipated within the atmosphere specified for various processes: Energy contained within  $\gamma$ 's falling below (energy or angular) cut, ionization energy loss of  $e^\pm$ , energy contained within  $e^\pm$  falling below (energy or angular) cuts, ionization energy loss of  $\mu^\pm$ , energy contained within  $\mu^\pm$  falling below (energy or angular) cuts, ionization energy loss of hadrons, energy contained within hadrons falling below (energy or angular) cuts, energy taken away by  $\nu$ 's, and the energy sum for each depth bin. For  $\pi^\pm$  and  $K^\pm$  falling below the energy or angular cuts 1/4 of the energy is attributed to the hadronic energy, while 3/4 is attributed to neutrinos. For the neutral  $K_{L,S}^0$  this ratios are 1/2.

The last bin of the cut energies reflects the energy content of particles arriving at detector level. In this table the rest mass of nucleons and electrons is **not** counted, while the rest mass of the corresponding anti-particles (which might undergo annihilation) is counted **twice**. Energies of unstable hadrons and muons are added up including their rest mass. This counting is necessary to respect the target nucleons or electrons involved into the shower development, thus enabling a correct energy balancing.

For writing out the longitudinal tables to the '*DATnnnnnnn.long*' (rsp. '*DATnnnnnnnnnnn.long*') file or as 'LONG' blocks to the particle output file '*DATnnnnnnn*' (rsp. '*DATnnnnnnnnnnn*') see the FLONGOUT flag (page 101).

The control printout contains as well all kinds of warnings and error messages. System errors may be redirected in UNIX systems<sup>98</sup> to the standard output<sup>99</sup> (.lst file) by the & character following immediately the > character as given in the example:

```
./corsika <inputs >& /home/user/corsika/run/DATnnnnnnn.lst
```

assuming that all output should go to the directory /home/user/corsika/run/ as given in the example on page 125. Users are advised to check this printout carefully for any indications of problems during the run time and keeping it together with the particle output for later consultation.

<sup>98</sup>The usage of the korne-shell is assumed.

<sup>99</sup>For the naming of standard output see the footnotes page 27 and 133.

When errors occur the DEBUG option may help in localizing the bug. This option entails a very detailed printout of the simulation process that easily fills large disks when enabled for many showers.

All these informations are printed per event. For low energy primaries and high statistics this printout per event may not be necessary and can be suppressed by selecting a maximum number of showers to be printed (keyword MAXPRT page 107). At the end of each run a short run summary is printed with similar tables as for single showers but averaged over all showers of the actual run.

## 10.2 Normal Particle Output

The particle and Cherenkov photon output files contain the information about the simulation run and about all particles reaching observation levels. This is what has to be analyzed for detailed energy spectra and distributions. These files are written to the directory DSN (as defined by keyword DIRECT, page 108) as 'sequential' 'unformatted' FORTRAN files. They may be read by the utility programs *corsikaread.f* resp. *corsikaread\_thin.f* available in the subdirectory `src/utills/` or by some COAST based programs such as the ones available in the subdirectory `src/utills/coast/CorsikaRead`. The particle output file and the Cherenkov photon output file are structured as shown in Table 6, with the sub-blocks as given in Tables 7 to 17. All quantities are written as single precision real numbers by the subroutine *tobuf.c* (independently of 32-bit or 64-bit processors).

Block structure
RUN HEADER nrun
EVENT HEADER 1
DATABLOCK
DATABLOCK
...
...
(LONG 1:1)
...
(LONG 1:n)
EVENT END 1
EVENT HEADER 2
DATABLOCK
DATABLOCK
...
...
(LONG 2:1)
...
(LONG 2:n)
EVENT END 2
...
...
EVENT HEADER nevt
DATABLOCK
DATABLOCK
...
...
(LONG nevt:1)
...
(LONG nevt:n)
EVENT END nevt
RUN END nrun

Table 6: Block structure of the particle and Cherenkov photon output files. (The LONG blocks eventually occur only in the particle output file.)

Run header sub-block: (once per run)	
No. of word	Contents of word (as real numbers R*4)
1	'RUNH'
2	run number
3	date of begin run ( yymmdd )
4	version of program
5	number of observation levels (maximum 10)
5 + $i$	height of observation level $i$ in cm
16	slope of energy spectrum
17	lower limit of energy range
18	upper limit of energy range
19	flag for EGS4 treatment of em. component
20	flag for NKG treatment of em. component
21	kin. energy cutoff for hadrons in GeV
22	kin. energy cutoff for muons in GeV
23	kin. energy cutoff for electrons in GeV
24	energy cutoff for photons in GeV
	physical constants and interaction flags:
24 + $i$	$C(i), i = 1, 50$
75	XPINCL X-displacement of inclined observation plane
76	YPINCL Y-displacement of inclined observation plane
77	ZPINCL Z-displacement of inclined observation plane
78	THINCL $\theta$ angle of normal vector of inclined observation plane
79	PHINCL $\phi$ angle of normal vector of inclined observation plane
80	TDINCL Depth on the axis angle for perpendicular observation plane
80 + $i$	0, $i = 1, 11$ (no longer used)
92	ARRANG rotation angle (deg) between between array x-direction and North
93	NSHOW number of showers to be generated
94	flag for SLANT option
94 + $i$	CKA( $i$ ), $i = 1, 40$
134 + $i$	CETA( $i$ ), $i = 1, 5$
139 + $i$	CSTRBA( $i$ ), $i = 1, 11$
150 + $i$	0, $i = 1, 97$ (no longer used)
248	XSCATT scatter range in x direction for Cherenkov
249	YSCATT scatter range in y direction for Cherenkov
249 + $i$	HLAY( $i$ ), $i = 1, 5$
254 + $i$	AATM( $i$ ), $i = 1, 5$
259 + $i$	BATM( $i$ ), $i = 1, 5$
264 + $i$	CATM( $i$ ), $i = 1, 5$
270	NFLAIN
271	NFLDIF
272	NFLPIO+100×NFLPIF
273	NFLCHE+100×NFRAGM

Table 7: Structure of the run header sub-block.

Event header sub-block: (once per event)	
No. of word	Contents of word (as real numbers R*4)
1	'EVTH'
2	event number
3	particle id (particle code or $A \times 100 + Z$ for nuclei)
4	total energy in GeV
5	starting altitude in g/cm <sup>2</sup>
6	number of first target if fixed
7	z coordinate (height) of first interaction in cm (negative, if tracking starts at margin of atmosphere, see TSTART)
8	px momentum in x direction in GeV/c
9	py momentum in y direction in GeV/c
10	pz momentum in -z direction in GeV/c (pz is positive for downward going particles)
11	zenith angle $\theta$ in radian
12	azimuth angle $\phi$ in radian
13	number of different random number sequences (max. 10)
$11 + 3 \times i$	integer seed of sequence $i$
$21 + 3 \times i$	number of offset random calls (mod $10^6$ ) of sequence $i$
$31 + 3 \times i$	number of offset random calls ( $/ 10^6$ ) of sequence $i$
44	run number
45	date of begin run (yymmdd)
46	version of program
47	number of observation levels
$47 + i$	height of level $i$ in cm
58	slope of energy spectrum
59	lower limit of energy range in GeV
60	upper limit of energy range in GeV
61	cutoff for hadrons kinetic energy in GeV
62	cutoff for muons kinetic energy in GeV
63	cutoff for electrons kinetic energy in GeV
64	cutoff for photons energy in GeV
65	NFLAIN
66	NFLDIF
67	NFLPIO
68	NFLPIF
69	NFLCHE
70	NFRAGM
71	x component of Earth's magnetic field in $\mu\text{T}$
72	z component of Earth's magnetic field in $\mu\text{T}$
73	flag for activating EGS4
74	flag for activating NKG

Table 8: Structure of event header sub-block (to be continued).

Event header sub-block: (continued)	
No. of word	Contents of word (as real numbers R*4)
75	low-energy hadr. model flag (1.=GHEISHA, 2.=UrQMD, 3.=FLUKA)
76	high-energy hadr. model flag (0.=HDPM,1.=VENUS, 2.=SIBYLL, 3.=QGSJET, 4.=DPMJET, 5.=NEXUS, 6.=EPOS)
77	CERENKOV flag <sup>100</sup>
78	NEUTRINO flag
79	CURVED flag (0=standard, 2=CURVED)
80	computer flag (3=UNIX, 4=Macintosh)
81	lower edge of $\theta$ interval (in $^\circ$ )
82	upper edge of $\theta$ interval (in $^\circ$ )
83	lower edge of $\phi$ interval (in $^\circ$ )
84	upper edge of $\phi$ interval (in $^\circ$ )
85	Cherenkov bunch size in the case of Cherenkov calculations
86	number of Cherenkov detectors in x-direction
87	number of Cherenkov detectors in y-direction
88	grid spacing of Cherenkov detectors in x-direction in cm
89	grid spacing of Cherenkov detectors in y-direction in cm
90	length of each Cherenkov detector in x-direction in cm
91	length of each Cherenkov detector in y-direction in cm
92	Cherenkov output directed to particle output file (= 0.) or Cherenkov output file (= 1.)
93	angle (in rad) between array x-direction and magnetic north
94	flag for additional muon information on particle output file
95	step length factor for multiple scattering step length in EGS4
96	Cherenkov wavelength lower end (in nm)
97	Cherenkov wavelength upper end (in nm)
98	number $i$ of uses of each Cherenkov rsp. Auger event
98 + $i$	x coordinate of $i^{th}$ core location for scattered events (in cm)
118 + $i$	y coordinate of $i^{th}$ core location for scattered events (in cm)
139	SIBYLL interaction flag (0.= no SIBYLL, 1.=vers.1.6; 2.=vers.2.1; 3.=vers.2.3c; 4.=vers.2.3d; 5.=vers.2.3e; 6.=vers.2.3e-star)
140	SIBYLL cross-section flag (0.= no SIBYLL, 1.=vers.1.6; 2.=vers.2.1; 3.=vers.2.3c; 4.=vers.2.3d/e)
141	QGSJET interaction flag (0.=no QGSJET, 1.=QGSJETOLD, 2.=QGSJET01d, 3.=QGSJET-II)

Table 9: (continued) Structure of event header sub-block (to be continued).

<sup>100</sup>EVTH(77) has the following contents if converted to an integer with suitable rounding applied:

bit 1 CERENKOV option compiled in

2 IACT option compiled in

3 CEFFIC option compiled in

4 ATMEXT option compiled in

5 ATMEXT option used with refraction enabled

6 VOLUMEDET option compiled in

7 CURVED option compiled in (see also EVTH(79))

9 SLANT option compiled in

11-21 table number for external atmosphere table (but limited to 1023 if the number is larger).

Event header sub-block: (continued)	
No. of word	Contents of word (as real numbers R*4)
142	QGSJET cross-section flag (0.=no QGSJET, 1.=QGSJETOLD, 2.=QGSJET01d, 3.=QGSJET-II)
143	DPMJET interaction flag (0.=no DPMJET, 1.=DPMJET)
144	DPMJET cross-section flag (0.=no DPMJET, 1.=DPMJET)
145	VENUS/NEXUS/EPOS cross-section flag (0=neither, 1.=VENUSSIG, 2./3.=NEXUSSIG, 4.=EPOSSIG, 5.=LHCRSIG)
146	muon multiple scattering flag (1.=Molière, 0.=Gauss)
147	NKG radial distribution range in cm
148	EFRCTHN energy fraction of thinning level hadronic
149	EFRCTHN·THINRAT energy fraction of thinning level em-particles
150	actual weight limit WMAX for thinning hadronic
151	actual weight limit WMAX·WEITRAT for thinning em-particles
152	max. radius (in cm) for radial thinning (THIN) rsp. cutting (CORECUT)
153	inner angle of viewing cone VIEWCONE (in °)
154	outer angle of viewing cone VIEWCONE (in °)
155	transition energy high-energy/low-energy model (in GeV)
156	skimming incidence flag (0.=standard, 1.=skimming)
157	altitude (cm) of horizontal shower axis (skimming incidence)
158	starting height (cm)
159	flag indicating that explicite charm generation is switched on
160	flag for hadron origin of electromagnetic subshower on particle tape
161	minimal vertical depth for transfer of particles from CONEX to CORSIKA
162	high-energy threshold for treatment of hadrons by cascade equations in CONEX
163	high-energy threshold for treatment of muons by cascade equations in CONEX
164	high-energy threshold for treatment of em-particles by cascade equations in CONEX
165	low-energy threshold for treatment of hadrons by cascade equations in CONEX
166	low-energy threshold for treatment of muons by cascade equations in CONEX
167	low-energy threshold for treatment of em-particles by cascade equations in CONEX
168	flag for observation level curvature (CURVOUT) (0.=flat, 1.=curved)
169	actual weight limit whmax for thinning hadronic in CONEX
170	actual weight limit wtmax for thinning em-particles in CONEX
171	actual weight limit whmax for sampling hadronic in CONEX
172	actual weight limit wtmax for sampling muons in CONEX
173	actual weight limit wtmax for sampling em-particles in CONEX
174	half width of stripes parallel to detector rows in AUGERHIT option (cm)
175	detector distance between neighbouring detectors (AUGERHIT) (cm)
176	(reserved for AUGERHIT)
177	number of MULTITHIN modes (max. 6)
177 + j	energy fraction of hadronic thinning for $j^{th}$ MULTITHIN mode
183 + j	actual weight limit of hadronic thinning for $j^{th}$ MULTITHIN mode
189 + j	energy fraction of em thinning for $j^{th}$ MULTITHIN mode
195 + j	actual weight limit of em thinning for $j^{th}$ MULTITHIN mode
199 + 3 × j	integer seed of random sequence for $j^{th}$ MULTITHIN mode
205 + 3 × j	# of offset random calls (mod $10^6$ ) for $j^{th}$ MULTITHIN mode
211 + 3 × j	# of offset random calls (exceeding $10^6$ ) for $j^{th}$ MULTITHIN mode
220	threshold energy above which particles are interesting (ICECUBE)
221	flag indicating that output is compressed by gzip (ICECUBE)
222	flag indicating that output is written to pipe buffer (ICECUBE)
223	index of MULTITHIN weight used in COAST
224	flag indicating MULTITHIN output is active
225	X cartesian coordinate of the origin of the inclined observation plane relative to shower core
226	Y cartesian coordinate of the origin of the inclined observation plane relative to shower core
227	Z apparent height of the origin of the inclined observation plane relative to sea level at the vertical of shower core 140
228	zenith angle $\Theta$ of the normal to the inclined observation plane in radian
229	zenith angle $\Phi$ of the normal to the inclined observation plane in radian
230	depth (slant on the shower axis or vertical) of the origin of the inclined observation plane
231	FLUKA version (MAJOR*1E5+MINOR*100+RESPIN or -1 for FLUKACERN)
232...273	not used

Table 9: (continued) Structure of event header sub-block.



Particle data sub-block : (up to 39 particles, 7 words each)	
No. of word	Contents of word (as real numbers R*4)
$7 \times (n - 1) + 1$	particle description encoded as: part. id $\times$ 1000 + hadr. generation <sup>101</sup> $\times$ 10 + no. of obs. level [for additional muon information or id 95/96: part. id $\times$ 1000 + hadr. generation <sup>102</sup> ]
$7 \times (n - 1) + 2$	px, momentum in x direction in GeV/c
$7 \times (n - 1) + 3$	py, momentum in y direction in GeV/c
$7 \times (n - 1) + 4$	pz, momentum in -z direction in GeV/c
$7 \times (n - 1) + 5$	x position coordinate in cm
$7 \times (n - 1) + 6$	y position coordinate in cm
$7 \times (n - 1) + 7$	t time since first interaction (or since entrance into atmosphere) <sup>103</sup> in nsec [for additional muon information or id 95/96: z coordinate in cm]
	for $n = 1 \dots 39$ if last block is not completely filled, trailing zeros are added

Table 10: Structure of particle data sub-block.

Cherenkov photon data sub-block : (up to 39 bunches, 7 words each)	
No. of words	Contents of word (as real numbers R*4)
$7 \times (n - 1) + 1$	number of Cherenkov photons in bunch <sup>104</sup> [in case of output on the particle output file: 99.E5 + 10 $\times$ NINT(number of Cherenkov photons in bunch) + 1]
$7 \times (n - 1) + 2$	x position coordinate in cm
$7 \times (n - 1) + 3$	y position coordinate in cm
$7 \times (n - 1) + 4$	u direction cosine <sup>105</sup> to x-axis
$7 \times (n - 1) + 5$	v direction cosine <sup>105</sup> to y-axis
$7 \times (n - 1) + 6$	t time since first interaction (or since entrance into atmosphere) <sup>103</sup> in nsec
$7 \times (n - 1) + 7$	height of production of bunch in cm <sup>106</sup>
	for $n = 1 \dots 39$ if last block is not completely filled, trailing zeros are added

Table 11: Structure of Cherenkov photon data sub-block.

<sup>101</sup>Without EHISTORY option, the hadronic generation counter  $h$  saturates at 99 (mainly EM particles from muons). With EHISTORY,  $h$  is last 2 digits of the full generation counter  $G$  (see Section 3.5.15 page 55).

<sup>102</sup>With muon additional information the hadronic generation counter  $h$  is fully given (3 digits). If the muon is produced after photonuclear interaction or was created by muon pair production +500 is added to  $h$ .

<sup>103</sup>See keyword TSTART page 78.

<sup>104</sup>For the THIN option the photon bunch size is multiplied with the thinning weight. See Sect. 4.90 page 111.

Mother particle data : (up to 39 particles, 7 words each)	
No. of word	Contents of word (as real numbers R*4)
$7 \times (n - 1) + 1$	particle description encoded as: –(part. id $\times$ 1000 + hadronic generation <sup>102</sup> of mother)
$7 \times (n - 1) + 2$	px, momentum in x direction in GeV/c
$7 \times (n - 1) + 3$	py, momentum in y direction in GeV/c
$7 \times (n - 1) + 4$	pz, momentum in -z direction in GeV/c
$7 \times (n - 1) + 5$	x position coordinate in cm
$7 \times (n - 1) + 6$	y position coordinate in cm
$7 \times (n - 1) + 7$	z position coordinate in cm at <b>creation point</b>
	in data line after muon additional information (for muons only) and before grandmother particle data line

Table 12: Structure of mother particle data in particle data sub-block.

Grandmother particle data : (up to 39 particles, 7 words each)	
No. of word	Contents of word (as real numbers R*4)
$7 \times (n - 1) + 1$	particle description encoded as: –(part. id $\times$ 1000 + electromagnetic generation <sup>107</sup> of <b>mother</b> )
$7 \times (n - 1) + 2$	px, momentum in x direction in GeV/c
$7 \times (n - 1) + 3$	py, momentum in y direction in GeV/c
$7 \times (n - 1) + 4$	pz, momentum in -z direction in GeV/c
$7 \times (n - 1) + 5$	ext. electromagn. gen. <sup>108</sup> counter of final particle
$7 \times (n - 1) + 6$	penetration depth in g/cm <sup>2</sup>
$7 \times (n - 1) + 7$	–z position coordinate in cm at <b>interaction point</b> (negative to make a difference with mother !)
	in data line after mother particle data line and before particle data line

Table 13: Structure of grandmother particle data in particle data sub-block.

<sup>105</sup>The direction cosines  $u$  and  $v$  are  $u = \sin(\theta) \cdot \cos(\phi)$  and  $v = \sin(\theta) \cdot \sin(\phi)$ , see Fig. 1 page 126.

<sup>106</sup>For MCERFI = 3 (see Sect. 4.90 page 111) the production height is replaced by the distance between production point and detector array center.

<sup>107</sup>The mother electromagnetic generation counter is the sum of the photon generation counter  $\gamma$  with 100 times the log of the electron generation counter  $e$  ( $\gamma + 100 \times \ln e$ ) as defined in Section 3.5.15 page 55.

<sup>108</sup>The extended electromagnetic generation counter  $\tilde{G}$  can be associated to the normal hadronic counter  $h$  given in the particle description field of the particle sub-block to get the full generation counter  $G = \tilde{G} \times 100 + h$  described in Section 3.5.15 page 55.

'Multithin' weights in particle data sub-block if MTHOUT active <sup>109</sup> : (preceded by particle in particle data line, 7 words each)	
No. of word	Contents of word (as real numbers R*4)
$7 \times (n - 1) + 1$	identification word 8888jjj.
$7 \times (n - 1) + 2$	weight of 1 <sup>st</sup> thinning mode
$7 \times (n - 1) + 3$	weight of 2 <sup>nd</sup> thinning mode
$7 \times (n - 1) + 4$	weight of 3 <sup>rd</sup> thinning mode
$7 \times (n - 1) + 5$	weight of 4 <sup>th</sup> thinning mode
$7 \times (n - 1) + 6$	weight of 5 <sup>th</sup> thinning mode
$7 \times (n - 1) + 7$	weight of 6 <sup>th</sup> thinning mode
	for $n = 1 \dots 39$

Table 14: Structure of 'multithin' weights in particle data sub-block.

### 10.2.1 Size Reduction of Normal Particle Output File

There exist different methods to reduce the size of the normal particle output file:

- a) Employing the AUGERHIT option: Only those particles are written, which arrive in the neighbourhood of detectors, while particles far from detectors are suppressed. Suppression factors of up to 500 can be achieved [47]. For details see Sect. 3.5.5 page 49 and Sects. 4.74 resp. 4.75 page 106 resp. 107.
- b) Without THIN option: The keyword CORECUT (Sect. 4.62 page 99) can be employed to discard those particles close to the shower axis which anyway would saturate the detector. By a suitable choice of the parameter RCUT of this keyword the size of the output can be adjusted to the particle data output file size limits.
- c) With THIN option: By suitably setting the thin-sampling parameters with keywords THIN (Sect. 4.55 page 96), THINH (Sect. 4.56 page 97) resp. THINEM (Sect. 4.57 page 97) the size of the particle data output file is influenced. Similarly to case b) a suitable choice of RMAX in the keyword THIN reduces the number of those particles which arrive close to the shower axis.  
Alternatively the writing of particles which arrive close to the shower core may be completely disabled using the CORECUT keyword with suitably selected RCUT. It should be noted that RCUT of keyword CORECUT resp. MTHINR overrides RMAX of keyword THIN.

Both cases b) without thinning resp. c) with thinning can be combined with the method a) employing the AUGERHIT option.

<sup>109</sup>See Sect. 3.5.20 page 59 and Table 10 page 141.

### 10.2.2 Normal Particle Output: Block Structure without Thinning

The information is stored unformatted in a fixed block structure with a block length of 22932 bytes. A block consists of 5733 words each 4 bytes long<sup>110</sup>. Each block consists of 21 sub-blocks of 273 words. These sub-blocks can be a RUN HEADER, EVENT HEADER, DATA BLOCK, LONGitudinal, EVENT END, or a RUN END sub-block (see Table 6). The contents of the sub-blocks are listed in Tables 7 to 17. This block structure is also maintained in the MULTITHIN option.

### 10.2.3 Normal Particle Output: Block Structure with Thinning

To take the weight parameter for each particle the data structure of the version without thinning has to be extended for the THIN option. The structure of the output (see Table 6 page 136) remains unchanged, but the blocksize is enlarged to a length of 26208 bytes. Now a block has 6552 words each 4 bytes long<sup>110</sup>, consisting of 21 sub-blocks of 312 words. The ends of the sub-blocks RUN HEADER (see Table 7 page 137), EVENT HEADER (see Table 9 page 139), LONGitudinal (see Table 15 page 145), EVENT END (see Table 16 page 146), and RUN END (see Table 17) are filled up with zeros, while the data blocks (Tables 10 and 11 page 141) contain 8 words for each particle resp. Cherenkov bunch, the last one being the weight<sup>111</sup>. Again 39 particles are collected within one data sub-block.

### 10.2.4 Normal Particle Output: Compact Option

The compact particle output available in the COMPACT option (Sect. 3.5.9 page 52) is organized similar to the normal particle output (see Sect. 10.2 page 136 ff.) with the following modifications:

- The block size of the records has variable length.
- In the block structure (Table 6 page 136) the 'event end' blocks are omitted completely.
- Only the first event header block has its full length (273 resp. 312 4-byte words) and starts with the characters 'EVTH'. The headers of subsequent events are shortened and contain only the first 12 parameters of Table 9 (page 139) and start with the characters 'EVHW'.
- The data blocks have variable length, trailing zeros are omitted. They are headed by a 4-byte integer which defines the length (in 4 byte words) of the appended data block. The maximal length is to 1+273 4-byte words (resp. 1+312 4-byte words for thinning).

Therefore a special reading routine is necessary for this output. The Cherenkov output is not affected by the COMPACT option, but it must be directed to the Cherenkov output file MCETAP (page 111, MCERFI  $\geq 1$ ).

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<sup>110</sup>On UNIX installations using the GNU g77 compiler, on HP UNIX stations, and some other machines the blocks comprise two additional words giving the record length.

<sup>111</sup>Concerning the weight calculation see Sect. 3.5.36 page 68 ff.

<sup>112</sup>Nuclei ( $A > 1$ ) are not counted with the 'hadron' species. They are assumed to be completely

'Longitudinal' sub-block: (up to 26 depth steps/block)	
No. of word	Contents of word (as real numbers R*4)
1	'LONG'
2	event number
3	particle id (particle code or $A \times 100 + Z$ for nuclei)
4	total energy in GeV
5	(total number of longitudinal steps) $\times$ 100 + number of longitudinal blocks/shower
6	current number $m$ of longitudinal block
7	altitude of first interaction in g/cm <sup>2</sup>
8	zenith angle $\theta$ in radian
9	azimuth angle $\phi$ in radian
10	cutoff for hadron kinetic energy in GeV
11	cutoff for muon kinetic energy in GeV
12	cutoff for electron kinetic energy in GeV
13	cutoff for photon energy in GeV
$10 \times n + 4$	vertical (rsp. slant) depth of step $j$ in g/cm <sup>2</sup>
$10 \times n + 5$	number of $\gamma$ -rays at step $j$
$10 \times n + 6$	number of $e^+$ particles at step $j$
$10 \times n + 7$	number of $e^-$ particles at step $j$
$10 \times n + 8$	number of $\mu^+$ particles at step $j$
$10 \times n + 9$	number of $\mu^-$ particles at step $j$
$10 \times n + 10$	number of hadronic particles at step $j$
$10 \times n + 11$	number of all charged particles at step $j$
$10 \times n + 12$	number of nuclei <sup>112</sup> at step $j$
$10 \times n + 13$	number of Cherenkov photons at step $j$
	for $n = 1, 26$ and for $j$ longitudinal steps  for 1 <sup>st</sup> 'LONG' block: 1 ... $j$ ... 26 for 2 <sup>nd</sup> 'LONG' block: 27 ... $j$ ... 52 ..... for $m^{\text{th}}$ 'LONG' block: $(m - 1) \cdot 26 + 1$ ... $j$ ... $m \cdot 26$  if last block is not completely filled, trailing zeros are added

Table 15: Structure of 'longitudinal' sub-block. (These type of blocks are written only if 'LONGI' is enabled and 'FLONGOUT' is disabled, page 101.)

Event end sub-block : (once per event)	
No. of word	Contents of word (as real numbers R*4)
1	'EVTE'
2	event number
3	statistics for one shower :
4	weighted number of photons arriving at observation level(s)
5	weighted number of electrons arriving at observation level(s)
6	weighted number of hadrons arriving at observation level(s)
7	weighted number of muons arriving at observation level(s)
	number of weighted particles written to particle output file MPATAP. (This number includes also Cherenkov bunches, if Cherenkov output is directed to MPATAP, but excludes additional muon information.)
	NKG output (if selected) :
$7 + i$	$i = 1, 21$ lateral distribution in x direction for 1. level in $\text{cm}^{-2}$
$28 + i$	$i = 1, 21$ lateral distribution in y direction for 1. level in $\text{cm}^{-2}$
$49 + i$	$i = 1, 21$ lateral distribution in xy direction for 1. level in $\text{cm}^{-2}$
$70 + i$	$i = 1, 21$ lateral distribution in yx direction for 1. level in $\text{cm}^{-2}$
$91 + i$	$i = 1, 21$ lateral distribution in x direction for 2. level in $\text{cm}^{-2}$
$112 + i$	$i = 1, 21$ lateral distribution in y direction for 2. level in $\text{cm}^{-2}$
$133 + i$	$i = 1, 21$ lateral distribution in xy direction for 2. level in $\text{cm}^{-2}$
$154 + i$	$i = 1, 21$ lateral distribution in yx direction for 2. level in $\text{cm}^{-2}$
$175 + i$	$i = 1, 10$ electron number in steps of $100 \text{ g/cm}^2$
$185 + i$	$i = 1, 10$ pseudo-age in steps of $100 \text{ g/cm}^2$
$195 + i$	$i = 1, 10$ distances for electron distribution in cm
$205 + i$	$i = 1, 10$ local pseudo-age 1. level
$215 + i$	$i = 1, 10$ height of levels for electron numbers in $\text{g/cm}^2$
$225 + i$	$i = 1, 10$ height of levels for electron numbers in cm
$235 + i$	$i = 1, 10$ distance bins for local pseudo-age in cm
$245 + i$	$i = 1, 10$ local pseudo-age 2. level
$255 + i$	$i = 1, 6$ parameters of longitudinal distribution of charged particles
262	$\chi^2$ per degree of freedom of fit to longitudinal distribution
263	weighted number of photons written to particle output file
264	weighted number of electrons written to particle output file
265	weighted number of hadrons written to particle output file
266	weighted number of muons written to particle output file
267	number of em-particles emerging from pre-shower
268..273	not used

Table 16: Structure of event end sub-block.

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stripped and therefore counted with their charge  $Z$  in the 'all charged' species.

Run end sub-block : (once per run)	
No. of word	Contents of word (as real numbers R*4)
1	'RUNE'
2	run number
3	number of events processed
4	not used
	in case of PARALLEL:
3	run number
4	number of cores used in this parallel run
5...273	not used yet

Table 17: Structure of run end sub-block.

### 10.2.5 Normal Particle Output with EHISTORY option

With the EHISTORY option, after activation, two additional data sub-block containing information about mother particle (first) and grandmother particle (second) are added before the data sub-block of the particle saved at the observation level. Their identification code is negative to avoid any confusion with final particles. In case of a muon, these two sub-blocks are placed between the sub-block of muon additional information and the sub-block for the muon itself. The contents of the sub-blocks are described in Table 12 for the mother and Table 13 page 142 for the grandmother.

## 10.3 .inclined File Output

This file is produced in the INCLINED option (Sect. 3.5.17 page 57) for inclined observation planes, using the keyword INCLIN (page 104). The format and structure is identical to the normal binary output data file as described in Sect. 10.2 (page 136), besides the fact that all coordinates are given within the inclined plane. The origin can be either defined by its coordinate in a flat frame defined in the standard CORSIKA coordinate system described in Sect. 6 (page 125) or automatically set at the intersection of the shower axis with the observation plane perpendicular to it. ( $X=0, Y=0, Z=0$ ) is in that case the shower axis. In both cases,  $\vec{Z}$  is the normal to the inclined plane (along the shower axis in the direction of the shower development in the automatic mode).  $\vec{X}$  is then in the plane defined by  $\vec{Z}$  and the vertical z-axis at the shower core.  $\vec{Y}$  is defined such that it defines an orthonormal coordinate system with  $(\vec{X}, -\vec{Z})$ . Since there is no  $Z$  coordinate in the output file (particles are per definition at  $Z = 0$ ), the negative sign is for the definition of the momentum which has a z-component positive in the direction of the shower development (exactly like in the standard coordinate system).

As shown in Fig. 3-a  $\theta$  and  $\phi$  angles are used to fix the direction of the normal of the plane. By definition if  $\theta = 0$  and  $\phi = 0$ , and  $z_0 = \text{OBSLEV}$ , the inclined shower plane is like the standard observation level (with  $X = x$  and  $Y = y$  with two very fundamental



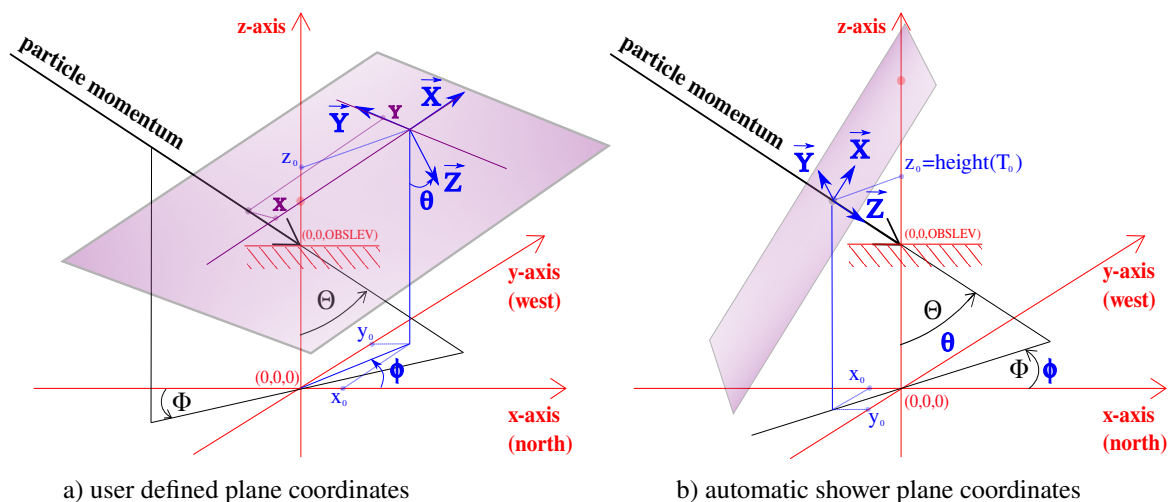


Figure 3: Coordinate system in inclined plane.

different properties: it doesn't stop the shower propagation and particles crossing in both direction are counted. If  $\theta = \Theta$  and  $\phi = \Phi$  the angle defining the shower direction, then the plane is perpendicular to the shower axis. If  $\theta = 90^\circ$ ,  $\phi = \Phi + 90^\circ$ , the shower axis will be contained in the inclined plan and the shower can be observed from the side (longitudinal development but only for particles at the axis !). The momentum  $p_z$  can be used to determine if a particle goes in the direction of the normal ( $p_z > 0$ ) or in the opposite direction ( $p_z < 0$ ).

To ease the theoretical studies of particles in a plane perpendicular to the shower plane, the last argument of the INCLIN keyword, TDINCL, is used to define the depth  $T_0$  of the position of the plane as shown in Fig. 3-b. If the SLANT option is used,  $T_0$  is the slant depth along the shower axis starting from the starting point of the shower. If FIXHEI is used (with or without STACKIN),  $T_0$  is the slant depth from the point on the shower axis defined by FIXHEI. Without the SLANT option,  $T_0$  is a fixed vertical depth which will be used to properly define  $(x_0, y_0, z_0)$  such that the shower axis is at  $(X=0, Y=0)$  in the plane perpendicular to the shower axis at this height. Without SLANT, FIXHEI has no influence on the position of the plane. It is simply a fixed height but define in  $\text{g}/\text{cm}^2$  (between 0 (top of the atmosphere) and  $1036 \text{ g}/\text{cm}^2$  (see level for US Std Atmosphere)).  $T_0$  cannot be defined beyond the observation level since no particle can reach there.

Note that a wrong definition of the plane coordinates can lead to an empty output file. Please refer to Sect. 6 (page 125) for a proper definition of the coordinates in particular using UPWARD or UPWARDOLD. The frame doesn't change but the position of the origin is different.

histo number	dimension	content of histograms
1 - 28	1-dim	longitudinal distributions
99	2-dim	particle codes vs. $\log_{10}(r)$
101 - 331	1-dim	ground particle distributions
341 - 451	2-dim	ground particle distributions
999	1-dim	density normalization
9111	1-dim	slant depth of muons
9112	2-dim	slant depth vs. $\log_{10}(\text{energy})$ of muons
9114	2-dim	slant depth vs. $\log_{10}(p_t)$ of muons
9115	2-dim	$\log_{10}(\text{energy})$ vs. $\log_{10}(p_t)$ of muons
9121	1-dim	mass overlay of muons
9122	2-dim	mass overlay vs. $\log_{10}(\text{energy})$ of muons
9123	2-dim	mass overlay vs. $\log_{10}(p_t)$ of muons
9127	2-dim	mass overlay vs. $\log_{10}(y\text{-coordinate})$ of muons
8000-8200	2-dim	$\log_{10}(\text{energy})$ vs. $\log_{10}(p_t)$ for slant depths between 0 and 2000 g/cm <sup>2</sup> in steps of 10 g/cm <sup>2</sup>
10iijj	1-dim	lateral particle distributions
20ii10	1-dim	Cherenkov photons vs. emission angle
30iijj	1-dim	energy spectra
40iijj	1-dim	deposited energy vs. distance
41iijj	1-dim	releasable energy vs. distance

Table 18: Histogram numbering of the 'datnnnnnn.lhbook' file.

## 10.4 .lhbook File Output

The optional file 'datnnnnnn.lhbook' (rsp. 'datnnnnnnnn.lhbook') written to the directory DSN (see keyword DIRECT, page 108) contains histograms produced by the ANAHIST (page 47), AUGERHIST (page 48), and/or the MUONHIST (page 60) options. Do not use capitals in the DSN directory name as the HBOOK routines tolerate only lower case characters. Table 18 gives the numbering of the hbook histograms. The histograms with numbers <1000 are produced by ANAHIST, those with numbers between 8000 and 9200 by MUONHIST, while those with numbers above 100000 come from the AUGERHIST option. If more than one of those options are selected simultaneously, the series of numbers appear in parallel without interferences. All series of histograms are added up for many showers. In case of ANAHIST and AUGERHIST they are normalized correctly before writing them to output. The AUGERHIST histograms are produced for up to 20 levels (to be specified by the keyword OBSLEV, page 103) which are denoted by ii with  $01 \leq ii \leq 20$ , and for various particle types denoted by jj with  $00 \leq jj \leq 04$  with the meaning: 00 = *gamma*, 01 = *electron*, 02 = *positron*, 03 = *muon*, 04 = *hadron*.

## 10.5 .long File Output

The optional file '*DATnnnnnnn.long*' (rsp. '*DATnnnnnnnnnn.long*') written to the directory DSN (keyword DIRECT page 108) contains a table of the longitudinal distribution of various particle numbers, arranged in the columns:

*depth,  $\gamma$ ,  $e^+$ ,  $e^-$ ,  $\mu^+$ ,  $\mu^-$ , hadrons, charged particles, nuclei, Cherenkov photons*<sup>113,114</sup>.

In a second table the longitudinal distributions of energy deposit (in GeV) by various particle species are given in columns: *depth,  $\gamma$  energy cut,  $e^\pm$  ionization,  $e^\pm$  energy cut,  $\mu^\pm$  ionization,  $\mu^\pm$  energy cut, hadron ionization, hadron energy cut, neutrino*<sup>115</sup>, *sum of all*.

For both tables the binning is in vertical depth (in g/cm<sup>2</sup>) as specified by the keyword LONGI, rsp. slant depth (in g/cm<sup>2</sup>) if the SLANT option (page 66) has been selected. This table output is activated by the keyword LONGI (page 101).

## 10.6 .tab File Output

The optional file '*DATnnnnnnn.tab*' (rsp. '*DATnnnnnnnnnn.tab*') written to the directory DSN (keyword DIRECT page 108) contains information on the particles arriving at the lowest detector level. Activation is done by the keyword PAROUT (page 108). There are 3 tables containing separately the number of  $\gamma$ ,  $e^\pm$ , and  $\mu^\pm$  particles binned into energy (40 bins ranging from 100 keV to 10 TeV in logarithmic steps), time delay relative to a spherical shower front (30 bins ranging from 10 nsec to 10  $\mu$ sec in logarithmic steps), and core distance (20 bins ranging from 50 m to 5 km in logarithmic steps).

Additionally a fourth table is written containing the longitudinal development of  $\gamma$ ,  $e^+$ ,  $e^-$ ,  $\mu^+$ ,  $\mu^-$ , *hadrons*, and *charged particles* (see Sect. 10.1 page 133) in steps as defined by keyword LONGI (page 101). This fourth table output is activated by the parameter LLONGI of keyword LONGI.

## 10.7 .dbase File and .info File Output

To build up a data base as a directory of a shower library which enables a computer aided search for specific shower events, the '*DATnnnnnnn.dbase*' (rsp. '*DATnnnnnnnnnn.dbase*') file (page 118) may be used. The content of the '*DATnnnnnnn.dbase*' file consists of parameter words enclosed within # marks, followed by the information on the corresponding parameter. The list of parameter words and their contents is given in Table 19 (page 152). In the AUGERINFO version this file is named '*DATnnnnnnn.info*' (rsp. '*DATnnnnnnnnnn.info*') and for each parameter a new line is started omitting the # mark separators. It should be noted, that some of the parameters listed in Table 19 are present only with the selected options e.g. for thinning (page 68), for external atmospheres (page 48), for viewing cone (page 71), or for Cherenkov telescopes (page 45).

---

<sup>113</sup>The Cherenkov photon longitudinal distributions are influenced by selecting the preprocessor options INTCLONG and NOCLONG (page 47, see also page 102) and AUGCERLONG (page 48).

<sup>114</sup>In case of the NEUTRINO or NUPRIM options the *neutrinos* are given instead of the *Cherenkov photons*.

<sup>115</sup>The counting of energy going into neutrinos is described in the third item of Sect. 10.1 (page 134).

Content of <i>DATnnnnnn.dbase</i> (rsp. <i>DATnnnnnn.info</i> ) file	
Keyword	Content (Format)
#version#	version number (F6.3)
#versiondate#	date of version YYYYMMDD (I9)
#modelversion#	version number of high-energy hadronic interaction program (F8.3)
#rundate#	date of run YYYYMMDD (I9)
#computer#	computer option (I2) (3=UNIX/LINUX, 4=Macintosh)
#curved#	number indicating CURVED option (I2)(2=curved, 0=else)
#neutrino#	number indicating NEUTRINO option (I2)
#cerenkov#	number indicating CERENKOV option (I2)
#runnumber#	run number (I7)
#primary#	particle code of primary particle (I5)
#e_range_l#	lower end of primary energy range (E14.7)
#e_range_u#	upper end of primary energy range (E14.7)
#slope#	slope of primary energy spectrum (E15.7)
#nkg#	number indicating use of NKG option (I2)
#egs#	number indicating use of EGS4 option (I2)
#model#	high-energy hadr. int. model (0=HDPM, 1=VENUS, 2=SIBYLL, 3=QGSJET, 4=DPMJET, 5=NEXUS/EPOS) (I2)
#gheisha#	low-energy hadr. model(1=GHEISHA, 2=UrQMD, 3=FLUKA) (I2)
#isobar#	low-energy hadr. model(1=GHEISHA, 2=UrQMD, 3=FLUKA) (I2)
#model+crosssect#	high-energy hadr. model and cross-sections (0=HDPM, 11=VENUS, 22=SIBYLL, 33=QGSJET, 44=DPMJET, 55=NEXUS, 66=EPOS) (I2);
#hadflag1#	number indicating NFLAIN (I2)
#hadflag2#	number indicating NFLDIF (I2)
#hadflag3#	number indicating NFLPI0 (I2)
#hadflag4#	number indicating NFLPIF (I2)
#hadflag5#	number indicating NFLCHE (I2)
#hadflag6#	number indicating NFRAGM (I2)
#longi#	number indicating use of longitudinal sampling (I2)
#longistep#	step width for longitudinal sampling (E14.7)
#magnetx#	horizontal component of Earth's magnetic field (E15.7)
#magnetz#	vertical component of Earth's magnetic field (E15.7)
#nobslev#	number of observation levels (I3)
#obslev1#	height of highest observation level (E15.7)
#obslev2#	height of 2nd observation level (E15.7)
#obslev3#	height of 3rd observation level (E15.7)
#obslev4#	height of 4th observation level (E15.7)
#obslev5#	height of 5th observation level (E15.7)
#obslev6#	height of 6th observation level (E15.7)
#obslev7#	height of 7th observation level (E15.7)
#obslev8#	height of 8th observation level (E15.7)
#obslev9#	height of 9th observation level (E15.7)
#obslev10#	height of 10th observation level (E15.7)
#hcut#	energy for hadron cut (E14.7)
#mcut#	energy for muon cut (E14.7)
#ecut#	energy for electron cut (E14.7)

Table 19: Content of *DATnnnnnn.dbase* (rsp. *DATnnnnnn.info*) file (to be continued).

Content of <i>DATnnnnnnn.dbase</i> (rsp. <i>DATnnnnnnn.info</i> ) file (continued)	
Keyword	Content (Format)
#gcut#	energy for gamma cut (E14.7)
#theta.l#	lower end of $\theta$ range (E14.7)
#theta.u#	upper end of $\theta$ range (E14.7)
#phi.l#	lower end of $\phi$ range (E15.7)
#phi.u#	upper end of $\phi$ range (E15.7)
#fixhei#	fixed height of first interaction (E14.7)
#n1sttr#	first target (I3) (0=random air, 1=Nitrogen, 2=Oxygen, 3=Argon)
#fixchi#	starting altitude of primary particle (E14.7)
#stepfc#	multiple scattering step length factor (E14.7)
#arrang#	array rotation angle (E15.7)
#himpact1#	lower limit of horizont. shower axis (skimming incid.) (E14.7)
#himpact2#	upper limit of horizont. shower axis (skimming incid.) (E14.7)
#muaddi#	number indicating use of additional muon information (I2)
#nseq#	number of used sequences for random generator (I2)
#seq1seed1#	seed of sequence 1 (I9)
#seq1seed2#	number of calls of sequence 1 (I9)
#seq1seed3#	billions of calls of sequence 1 (I9)
#seq2seed1#	seed of sequence 2 (I9)
#seq2seed2#	number of calls of sequence 2 (I9)
#seq2seed3#	billions of calls of sequence 2 (I9)
#seq3seed1#	seed of sequence 3 (I9)
#seq3seed2#	number of calls of sequence 3 (I9)
#seq3seed3#	billions of calls of sequence 3 (I9)
#size#	size of particle tape output (I10)
#dsn_events#	data set name of particle tape output (A59)
#dsn_prtout#	data set name of .lst file output (A9)
#tape_name#	name of data tape (A10)
#backup#	name of backup tape (A10)
#howmanyshowers#	number of showers to generate (I10)
#host#	host computer name (A20)
#user#	user name (A20)
#atmosphere#	Modtran atmosphere model number (I3)
#refract#	number indicating use of refractive index (I2)
#viewcon.l#	inner limiting angle of viewing cone (E14.7)
#viewcon.u#	outer limiting angle of viewing cone (E14.7)
#telescope <i>i</i> #	coordinates $x, y, z, r$ of telescope $i$ (4F11.1)
#cscat #	number and range of scattering in $x, y$ (2F10.1)
#thinning#	number indicating use of thinning (I2)
#thinnlev_had#	thinning level hadronic (E14.7)
#thinnlev_em#	thinning level em (E14.7)
#maxweight_had#	weight limit hadronic (E14.7)
#maxweight_em#	weight limit em (E14.7)
#rad_max#	maximum radius for radial thinning [m] (E14.7)
#energy_prim#	primary energy of first shower (E14.7)
#theta_prim#	primary's $\theta$ of first shower (E14.7)
#phi_prim#	primary's $\phi$ of first shower (E14.7)

Table 19: (continued) Content of *DATnnnnnnn.dbase* (rsp. *DATnnnnnnn.info*) file.

## 11 Hints for Programmers

If you need any option, addition, or other extension which is not (yet) covered by CORSIKA, feel free to program it. (Please contact T. Pierog before doing so.) If your addition is of general interest, it might be worth to take it over into the next official CORSIKA release. At the beginning of the *corsika.h* file you find a rather complete list of all global variables used in */COMMONS/*. This enables to check the names which you give your variables for conflicts with names already used within CORSIKA. The names of the CORSIKA commons start with CR... to prevent conflicts with common or subroutine names of linked interaction program packages.

As the largest fraction of the CORSIKA routines is written in CAPITAL letters, you are advised to use lower case characters for your private program extensions. This facilitates to distinguish your programming from the official code.

We advise you to use:

```
./coconut [help] [dev] [options] [configure options]
```

to handle your compilation/installation processes.

It is recommended to run *./coconut* (without arguments) to configure and install CORSIKA.

Various options are available in *./coconut* for expert use only:

- `--*`: Any additional option with trailing `--` will be passed on to *./configure* directly (try e.g. `--help`).
- `-h, --help`: Display a help file and stop.
- `-b, --batch`: No user interaction, just take DEFAULT and CACHED configuration.
- `-c, --clean`: Execute `make clean` (if CORSIKA is already installed once). Remove all object files and stop.
- `-d, --distclean`: Execute `make distclean` (if CORSIKA is already installed once). Remove all files produced by the installation (but not the binaries in *run/*) and stop.
- `-e, --expert`: Activate expert mode with additional configuration steps.
- `--dev`: Activate development mode without predefined system compilation flags and without optimization flag.
- `-i, --install`: Execute `make install` (if CORSIKA is already installed once). Compile, link and copy binaries into *run/*, but *./configure* is not called at all (except if it was never done before. If *Makefiles* are not yet installed, *./configure* will be called once).

- `-m, --make`: Execute `make all` (if CORSIKA is already installed once). Compile and link, but do not copy binaries to `run/`, and `./configure` is not called at all (no option selection ... except if it was never done before. If *Makefiles* are not yet installed, `./configure` will be called once).
- `-n, --no-cache`: Do not use cached configuration.
- `-t, --dist`: To run `make dist` (skips configuration).
- **Configure options**: Any option can be transferred to *configure*, for more information use `-h` or `--help`<sup>116</sup>, but `--prefix`, `--bindir` and `--libdir` are already defined in `./coconut`. You should not change this options. Values defined in `$CC`, `$CCLAGS`, `$CXX`, `$CXXFLAGS`, `$CPP`, `$CPPFLAGS` and `$LDFLAGS` environment variable are always used if defined. To use the values defined in `$F77` and `$FFLAGS` environment variable in `./coconut`, the following environment variable has to be defined:

```
setenv CORSIKA_USER_COMP yes
```

If `$CORSIKA_USER_COMP` is not empty, then `./coconut` does not set any FORTRAN compiler name or flags. If they are not defined by you (see A.1 to know the recommended flags), `./configure` will give some default value (compilation not guaranteed). So don't use this option if you don't know exactly what you are doing (no support for this).

Run it a first time to select the options you want to use (saved in `$libdir/include/config.h`) and then work directly on the `src/corsika.F` file (FORTRAN+preprocessor commands). To compile this modified source without calling *configure* again, use `./coconut [dev] make` for debugging. It's equivalent just to go into the proper subdirectory (depending on `dev`) and to type

```
rm -f compilefile.f
make
```

This will update the *compilefile.f* with your modifications and then compile it. When the compilation is successful, you can link objects and libraries into `run/corsika<VER>-<OS>-<HIGH>-<low>`<sup>117</sup> by typing

```
./coconut --install
```

---

<sup>116</sup>Call *configure* help.

<sup>117</sup>Where `<VER>` is the version number, `<OS>` is the operating system used for compilation, `<HIGH>` is the chosen high energy hadronic interaction model and `<low>` is the chosen low energy hadronic interaction model.

in your `corsika-78010/` directory<sup>118</sup>.

When the development phase is over, you can use the standard procedure

```
./coconut [dev]
```

to have binary files with different options (If `dev` is not used any more, you will have to type `./coconut --distclean` or `make distclean` in `corsika-78010/` first.).

We strongly recommend to use this scheme (instead of the traditional `compilefile.f`) because it is the easiest way for you and for us, if later this modification has to be implemented in the official release of CORSIKA (with proper preprocessor commands) or if you want to use your modifications with different CORSIKA options.

If you prefer to work on `compilefile.f` anyway, you can use the produced *Makefiles* to compile it. To produce the *Makefiles*, use `./coconut dev` (note the `dev` option<sup>119</sup>), select your options and, at the end, save the source code and do not compile. Then, you only have to rename your source file `corsika<VER><OS>_<HIGH>_<low>_compilefile.f` to `compilefile.f` and then you have to type `make` in `src/` to compile or `make install` in your `corsika-78010/` directory to (compile and) link to get a binary file. In this case, you should not use `./coconut` any more, because this will not take into account your modifications (and even overwrite the `compilefile.f`). Of course, if you already used some "home-made" *Makefiles* with a former version of CORSIKA, you can use it with this `compilefile.f`.

## 12 In Case of Problems ...

In spite of our care to avoid faults in the physics model and errors in the programming of CORSIKA and in writing this CORSIKA GUIDE, you may have problems of technical or physical nature with the code or the results you obtain from it. Please report all problems to the address mentioned on page 3. When applying for help in cases of crash, **it is recommended to transmit the control printout file (.lst file)**, as it contains the selected preprocessor options as well as it echos the employed input keywords with their parameters.

Special interest exists in comparisons of CORSIKA simulations with experimental cosmic ray data.

---

<sup>118</sup>The `./coconut` compiles and installs all the libraries needed by CORSIKA to be linked with and then creates a binary file in `run/`. Then, if you change `corsika.F`, doing simply `./coconut --make` will produce a `corsika` executable file in `src/`. You can then copy this file into `run/` to use it with all the data files.

<sup>119</sup>If you don't use the `dev` option when running `./coconut`, you can follow the same scheme but using *Makefiles* located in `corsika-78010/lib/'uname'/'src/` and `corsika-78010/lib/'uname'/'src/` (for instance `corsika-78010/lib/Linux/src/` if you work on a LINUX system) instead of just `corsika-78010/`.



Suggestions to improve CORSIKA in any respect are welcome.

## 13 Acknowledgments

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## A Compilation and Linking

Using the `./coconut` script, it will call `make` automatically and nothing has to be done. But for expert people who want to do their own `Makefile`, a description of what has to be done is given in this appendix.

### A.1 Compilation

After running `./coconut` asking not to compile at the end of the selection process, you get a source file `corsika<VER><OS>_<HIGH>_<low>_compilefile.f` in the `src/` subdirectory. This subdirectory is the working directory for the following commands. This `corsika_compilefile.f` FORTRAN file and the `dpmjet` FORTRAN files, the `epos-xxx.f`, `nerus-xxx.f`, `qgsjet-II-04.f` (rsp. `qgsjet01d.f`), `sibyll2.3d.f`, `venus.f`, and `gheisha_2002d.f` sources and/or the `urqmd` FORTRAN files have to be compiled before linking them together to obtain an executable module. To prevent overwriting of local variables by optimization of some compilers SAVE statements (which replace the '-static' option to be used for the previous release) have been implemented throughout in the FORTRAN files of `corsika_compilefile.f`, `gheisha_2002d.f`, the FORTRAN files of `dpmjet`, `qgsjet-II-04.f` (rsp. `qgsjet01d.f`), `sibyll2.3d.f`, and `venus.f`. It is recommended to use the 'bounds check'<sup>120</sup> option for first trials to inhibit uncontrolled array operations outside the allowed index range.

- For DEC-UNIX machines the procedure looks like  
`f77 -c -check_bounds $1.f 2>$1.err`  
(This means: Suppress the loading phase; generate code to perform runtime checks on subscript.)
- Procedures for LINUX computers with GNU g77 compilers should be used without optimization<sup>121</sup> and are  
`g77 -c -O0 $1.f 2>$1.err`  
to ensure correct simulations.
- Procedures for LINUX hosts (with 64bit AMD CPU's) with GNU g77 compilers should be used without optimization<sup>121</sup> and are  
`g77 -c -O0 $1.f 2>$1.err`  
to ensure correct simulations with the correct data format of the binary output files.
- Hosts with Portland pgf77 compiler available might use  
`pgf77 -c -O2 $1.f 2>$1.err`  
which gives a fast and reliable executable on LINUX hosts.

---

<sup>120</sup>But not using GFORTRAN compiler, it will result in unexpected stop.

<sup>121</sup>Do not use the optimization without carefully checking the results. There is bad experience with GNU g77 (v0.5.24 and egcs-2.91.66) which frequently brings NaN in the particle output file for the x and y coordinates of particles or results in unidentified hang-ups within the QGSJET routines. Also the g77 optimization causes the DADMUL integration routine to end with an error stop (message: DBRSGM: IFAIL= 1 ...) despite the correct programming respecting all FORTRAN standards.

- Compile procedures for IBM RS6000 are  
`xlf -c -C -O -qextname -qsave -qmaxmem=-1 $1.f 2>$1.err`
- For HP-UX processors the compilation procedures look like  
`fort77 -c -K +E1 +ppu +Dportable -O2 +Onolimit $1.f 2>$1.err`  
 and the +E1 option should also be used in the link step.

These compiler procedures should be used for all needed FORTRAN programs. Our experience with machines other than DEC or LINUX are sparse.

### **DPMJET option:**

The compiler procedures of the standard case (see begin of this subsection A.1) should be used. As the files *dpmjet.f*, *phojet.f*, and *pythia-6427\_cors.f* are extremely large you presumably have to give in advance<sup>122</sup>

```
limit datasize unlimited
limit stacksize unlimited
```

to overcome the small default values of many compilers which lead to an error stop during compilation.

**EPOS option:** The compiler procedures of the standard case (see begin of this subsection A.1) should be used.

To compile the *epos\*.f* source files in `../epos/` you should use the *Makefile.ka* available with EPOS (omitting the link step).

For compiling and linking EPOS you presumably have to give<sup>123</sup>

```
limit datasize unlimited
limit stacksize unlimited
```

to overcome the small default values of many compilers which lead to an error stop during linking. Linking is performed e.g. by

```
f77 corsika_compilefile.o gheisha_2002d.o ../epos/*.o -o corsika\vers
```

where the compiled EPOS program parts are expected to be available within the subdirectory `../epos/`.

**NEXUS option:** The compiler procedures of the standard case (see begin of this subsection A.1) should be used.

To compile the *nexus-xxx.f* source files in `../nexus/` you should use the *Makefile.ka* available with NEXUS (omitting the link step).

For compiling and linking NEXUS you presumably have to give<sup>123</sup>

```
limit datasize unlimited
limit stacksize unlimited
```

---

<sup>122</sup>Usage of *cs*h shell is assumed. In other shells (*sh* or *ksh*) you should use `ulimit -d unlimited` and `ulimit -s unlimited`.

<sup>123</sup>See footnote at DPMJET compilation page 160.

to overcome the small default values of many compilers which lead to an error stop during linking. Linking is performed e.g. by

```
f77 corsika_compilefile.o gheisha_2002d.o ../nexus/*.o -o corsika\vers
```

where the compiled NEXUS program parts are expected to be available within the subdirectory `../nexus/`.

**FLUKA CERN option:** For compiling the CORSIKA-FLUKACERN versions of the `corsika_compilefile.f`, the FLUKACERN include files should be available in the subdirectory from which you are calling the compiler to include them into the CORSIKA-FLUKACERN linking routines at the appropriate places. Preferentially you use the gfortran compiler with the option

```
-Iflukadirectory/include
```

to indicate the compiler where to find the include files. For all steps using the FLUKA packages you presumably have to give<sup>123</sup>

```
limit datasize unlimited  
limit stacksize unlimited
```

to overcome the small default values of many compilers which lead to an error stop during linking.

**FLUKA option:** For compiling the CORSIKA-FLUKA versions of the `corsika_compilefile.f`, the FLUKA include files should be available in the subdirectory from which you are calling the compiler to include them into the CORSIKA-FLUKA linking routines at the appropriate places. Preferentially you use the g77 resp. gfortran compiler with the option

```
-Iflukadirectory/flukapro
```

to indicate the compiler where to find the include files. For all steps using the FLUKA packages you presumably have to give<sup>123</sup>

```
limit datasize unlimited  
limit stacksize unlimited
```

to overcome the small default values of many compilers which lead to an error stop during linking.

**URQMD option:** For compiling the CORSIKA-URQMD version of the `corsika_compilefile.f`, the UrQMD include files `boxinc.f`, `colltab.f`, `comres.f`, `coms.f`, `inputs.f`, `newpart.f`, and `options.f` should be available in the directory from which you are calling the compiler to include them into the CORSIKA-URQMD linking routines at the appropriate places. Alternatively you may use the g77 resp. gfortran compiler with the option

```
-I../urqmd
```

to indicate the compiler where to find the include files. The compiler procedures of the standard case (see begin of this subsection A.1) should be used to compile the `corsika_compilefile.f`.

The *urqmd1.3\_cors.tar.gz* file contains the UrQMD1.3\_cors source routines with slight modifications to adapt them for the use with CORSIKA. To compile these UrQMD source files one uses the (g)make command with the option *-f GNUmakefile\_corsika* (omitting the link step) in the *../urqmd/* subdirectory.

**NUPRIM option:** When compiling the CORSIKA-NUPRIM version of the *corsika\_compile-file.f*, the *HERWIG\_C.INC* include file should be available in the subdirectory from which you are calling the compiler to include it into the CORSIKA-HERWIG linking routines at the appropriate places. Alternatively you may use the *g77* resp. *gfortran* compiler with the option

```
-I../herwig
```

to indicate the compiler where to find the include file.

This include file is a simple copy of the *HERWIGnn.INC* file provided with the source code of the HERWIG model. In the *herwig/* subdirectory you give the command:

```
cp HERWIGnn.INC HERWIG_C.INC
```

Before compiling the HERWIG routines you should

- remove (erase or comment off) the function *HWRGEN* with entries *HWRSET* and *HWRGET* thus using the CORSIKA random generator (sequence 5) within the HERWIG routines;
- when using the NUPRIM option together with NEXUS or EPOS: Rename the *COMMON /JET/* within *herwig6510.f*.
- when using the NUPRIM option together with PYTHIA (CHARM option): The subroutines *UPINIT* and *UPEVNT* appear in both codes. Please comment off the calls to these routines and both dummy routines in the *herwig\_corsika.f* to enable the link step without problems.

The compiler procedures of the standard case (see begin of this subsection A.1) should be used for the CORSIKA-NUPRIM version as well as for HERWIG.

**C-file:** If you have specified the *TIMERC* option, the *PRESHOWER* option, or the *EFIELD* option you need the compilation of the C-files *timerc.c*, *preshw.c* with *utils.c* and *veto.c*, or *elfield.c*, but also the routines *tobuf.c* and *trapfpe.c* using a command like

```
cc -c timerc.c      (rsp. cc -c preshw.c ).
```

For the compilation of the C-routines of the *bernlshr* package see the instructions delivered with the *bernlshr* package [44].

## A.2 Linking

For linking of most CORSIKA versions typically a procedure is used (assuming QGSJET is employed) like:

```
g77 corsika_compilefile.o gheisha_2002d.o qgsjet-II-04.o trapfpe.o \  
-o corsika
```

For linking the large packages of FLUKA, UrQMD, DPMJET, EPOS, HERWIG, NEXUS, and VENUS you presumably have to give<sup>124</sup>

```
limit datasize unlimited  
limit stacksize unlimited
```

to overcome the small default values of many compilers which lead to an error stop during linking.

If the `pgf77` compiler has been used for compilation, you should also use `pgf77` for linking:

```
pgf77 corsika_compilefile.o gheisha_2002d.o qgsjet-II-04.o -o corsika
```

For linking **FLUKACERN** versions the FLUKA library has to be included in the link step:

```
g77 corsika_compilefile.o qgsjet-II-04.o trapfpe.o \  
-Iflukadirectory/include -Lflukadirectory/lib -lfluka -o corsika
```

In the run step of FLUKACERN versions the environment variable FLUKADATA has to be set for *bash/sh* shells

```
export FLUKADATA=$(fluka-config --data)
```

which is necessary to link the data files which will be read in by the FLUKACERN package.

For linking **FLUKA** versions the FLUKA library has to be included in the link step:

```
g77 corsika_compilefile.o qgsjet-II-04.o trapfpe.o \  
-Iflukadirectory/include -Lflukadirectory -lflukahp -o corsika
```

For linking FLUKA with `pgf77` you should use the option `-g77libs` in the link step to ensure that the runtime library<sup>125</sup> of `g77` is available as it is needed for running the FLUKA routines. In the run step of FLUKA versions the environment variable FLUPRO has to be set for *cs* shell

```
setenv FLUPRO flukadirectory
```

or for *bash/sh* shells

```
export FLUPRO=flukadirectory
```

which is necessary to link the data files which will be read in by the FLUKA package. Additionally the stacksize and datasize limits have to be overcome also in the run step.

---

<sup>124</sup>See footnote at DPMJET compilation page 160.

<sup>125</sup>Be aware that the FLUKA LINUX version and the runtime library are fitting together.

Linking of **URQMD** versions is performed easiest (e.g. for DEC-UNIX) by

```
f77 corsika_compilefile.o qgsjet-II-04.o ../urqmd/obj_urqcors/*.o \  
trapfpe.o -o corsika78010.QGS-II.urq13
```

to include the `.o` files of UrQMD which are stored into the subdirectory `../urqmd/obj_urqcors` by the GNU-make compilation procedure.

No other libraries are normally required.

If you have specified the `TIMERC`, the `PRESHOWER` option, or the `EFIELD` option you need the linking with the compiled `C`-files `timerc.o`, resp. `preshw.o` with `utils.o` and `veto.o`, `elfield.o`, but also the `tobuf.o` and `trapfpe.o` modules.

For linking the compiled `C`-routines of the bernlohr package see the instructions delivered with the bernlohr package.



## B Flow Diagram

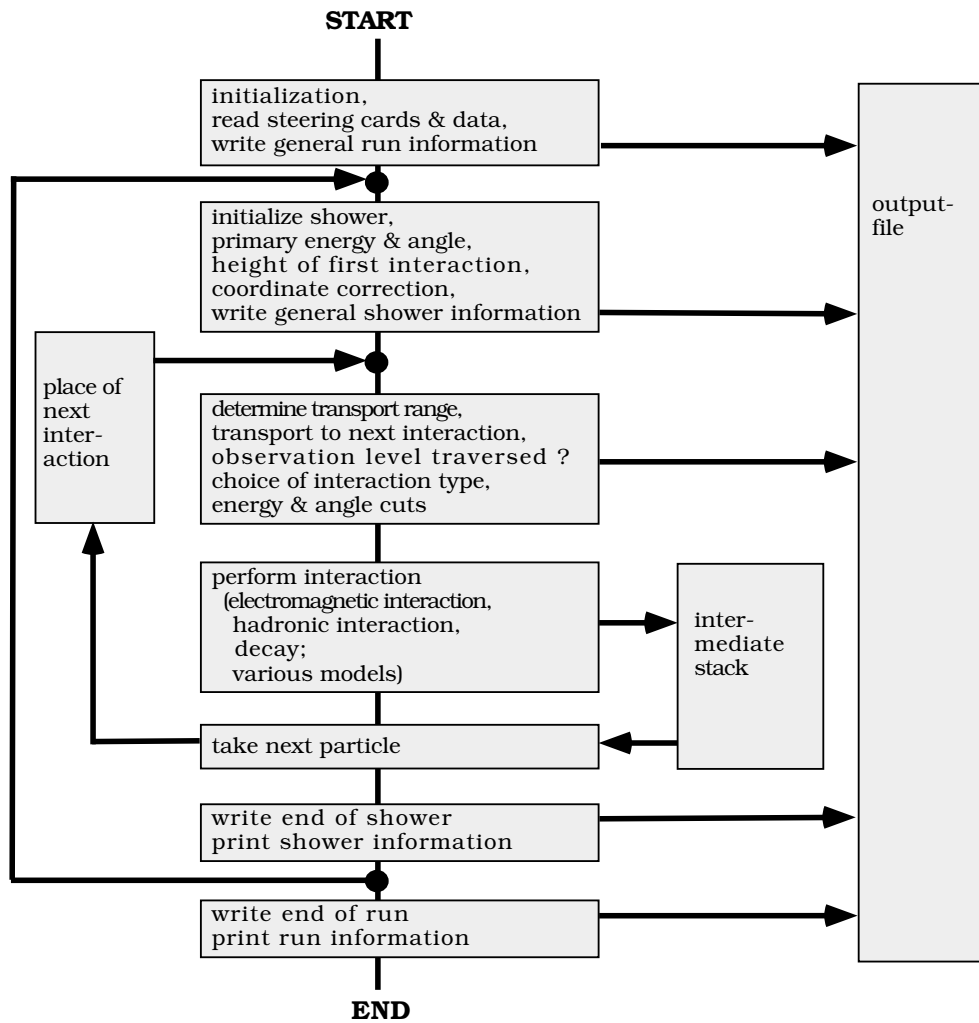


Figure 4: Simplified flow diagram of CORSIKA.

## C Sequence of Initializations

The sequence of initializations is shown for the QGSJET option (INTTEST option in brackets) as example, other interaction model options may differ slightly. The indentation gives the hierarchy of subroutines. Subroutine names are written in CAPITALS.

```
AAMAIN (CORSIKA main program)
  START
    write CORSIKA version and options
  PAMAF
    set particle masses and decay times
  DATAC
    read keywords from input
  set OBSLEV, atmospheres and layers
  initialize random generator
  read NUCNUCCS cross-section tables
  FILOPN
    open output files and external stack
  QGSINI
    initialize QGSJET-II
    QGSET (QGSJET-II)
    QGAINI (QGSJET-II)
      read qgsdat-II-04 and sectnu-II-04
    QGSSIGINI
  CGHINI
    initialize GHEISHA
  EGSIN1
    initialize EGS4
  EGSIN2
    read EGSDAT6.x.x
  calculate physical constants
  (set projectile and target for HSINI)
  (HSINI)
    (histogram initialization for INTTEST)
  INPRM
    check input parameters
    set various parameters: magn. field, Cherenkov...
    (print Cherenkov scatter table)
  MUPINI
    set  $\mu$ -parameters and  $\mu$ -cross-section tables
  write RUNHEADER
  write .dbase (rsp. .info) file
  ININKG
    initialize NKG parameters
  clear statistics arrays: multiplicity, elasticity, weight...
  shower loop
  .
  .
```

## D Particle Identification in Different Programs

Particle name	Particle mass <sup>126</sup> [ GeV ]	Particle life time <sup>126</sup> [ sec ]	Identification code					
			COR-SIKA	DPM-JET	PYTHIA HERWIG (PDG)	EPOS NEXUS (ISAJET)	QGS-JET	SIB-YLL
$\gamma$	0.	stable	1	7	22	10		1
$e^+$	.000511	stable	2	4	-11	-12		2
$e^-$	.000511	stable	3	3	11	12		3
$\mu^+$	.105658	$2.197 \cdot 10^{-6}$	5	10	-13	-14		4
$\mu^-$	.105658	$2.197 \cdot 10^{-6}$	6	11	13	14		5
$\pi^0$	.134977	$8.4 \cdot 10^{-17}$	7	23	111	110	0	6
$\pi^+$	.139570	$2.6033 \cdot 10^{-8}$	8	13	211	120	1	7
$\pi^-$	.139570	$2.6033 \cdot 10^{-8}$	9	14	-211	-120	-1	8
$K_L^0$	.49761	$5.18 \cdot 10^{-8}$	10	12	130	(230)	-5	11
$K^+$	.493677	$1.2384 \cdot 10^{-8}$	11	15	321	130	4	9
$K^-$	.493677	$1.2384 \cdot 10^{-8}$	12	16	-321	-130	-4	10
$n$	.939565	880.3	13	8	2112	1220	3	14
$p$	.938272	stable	14	1	2212	1120	2	13
$\bar{p}$	.938272	stable	15	2	-2212	-1120	-2	-13
$K_S^0$	.49761	$8.954 \cdot 10^{-11}$	16	19	310	(230)	5	12
$\eta$	.547862	$5.02 \cdot 10^{-19}$	17	31	221	220	10	23
$\Lambda$	1.11568	$2.632 \cdot 10^{-10}$	18	17	3122	2130	6	39
$\Sigma^+$	1.18937	$8.018 \cdot 10^{-11}$	19	21	3222	1130		34
$\Sigma^0$	1.192642	$7.4 \cdot 10^{-20}$	20	22	3212	1230		35
$\Sigma^-$	1.197449	$1.479 \cdot 10^{-10}$	21	20	3112	2230		36
$\Xi^0$	1.31486	$2.90 \cdot 10^{-10}$	22	97	3322	1330		37
$\Xi^-$	1.32171	$1.639 \cdot 10^{-10}$	23	98	3312	2330		38
$\Omega^-$	.1.67245	$8.21 \cdot 10^{-11}$	24	109	3334	3331		49
$\bar{n}$	.939565	880.3	25	9	-2112	-1220	-3	-14
$\bar{\Lambda}$	1.115683	$2.632 \cdot 10^{-10}$	26	18	-3122	-2130	-6	-39
$\bar{\Sigma}^-$	1.18937	$8.018 \cdot 10^{-11}$	27	99	-3222	-1130		-34
$\bar{\Sigma}^0$	1.192642	$7.4 \cdot 10^{-20}$	28	100	-3212	-1230		-35
$\bar{\Sigma}^+$	1.197449	$1.479 \cdot 10^{-10}$	29	101	-3112	-2230		-36
$\bar{\Xi}^0$	1.31486	$2.90 \cdot 10^{-10}$	30	102	-3322	-1330		-37
$\bar{\Xi}^+$	1.32171	$1.639 \cdot 10^{-10}$	31	103	-3312	-2330		-38
$\bar{\Omega}^+$	1.67245	$8.21 \cdot 10^{-11}$	32	115	-3334	-3331		-49

Table 20: Particles used in different programs (to be continued).

Particle name	Particle mass <sup>126</sup> [ GeV ]	Particle life time <sup>126</sup> [ sec ]	Identification code					
			COR-SIKA	DPM-JET	PYTHIA HERWIG (PDG)	EPOS NEXUS (ISAJET)	QGS-JET	SIB-YLL
$\eta'$	.95778	$3.32*10^{-21}$	48	95	331	330		24
$\phi$	1.0194	$1.54*10^{-22}$	49	96	333	331		33
$\omega$	.78265	$7.75*10^{-23}$	50	35	223	221		32
$\rho^{\circ}$	.769	$4.14*10^{-24}$	51	33	113	111	-10	27
$\rho^{+}$	.7665	$4.14*10^{-24}$	52	32	213	121		25
$\rho^{-}$	.7665	$4.14*10^{-24}$	53	34	-213	-121		26
$\Delta^{++}$	1.2305	$5.87*10^{-24}$	54	53	2224	1111		40
$\Delta^{+}$	1.2318	$5.02*10^{-24}$	55	54	2214	1121		41
$\Delta^{\circ}$	1.2331	$5.606*10^{-24}$	56	55	2114	1221		42
$\Delta^{-}$	1.2344	$5.*10^{-24}$	57	56	1114	2221		43
$\bar{\Delta}^{--}$	1.2309	$5.87*10^{-24}$	58	67	-2224	-1111		-40
$\bar{\Delta}^{-}$	1.2323	$5.02*10^{-24}$	59	68	-2214	-1121		-41
$\bar{\Delta}^{\circ}$	1.2336	$5.606*10^{-24}$	60	69	-2114	-1221		-42
$\bar{\Delta}^{+}$	1.2349	$5.*10^{-24}$	61	70	-1114	-2221		-43
$K^{*\circ}$	.89581	$1.398*10^{-23}$	62	37	313	231		30
$K^{*+}$	.89166	$1.296*10^{-23}$	63	36	323	131		28
$\bar{K}^{*-}$	.89166	$1.296*10^{-23}$	64	38	-323	-131		29
$\bar{K}^{*\circ}$	.89581	$1.398*10^{-23}$	65	39	-313	-231		31
$\nu_e$	0.	stable	66	5	12	11		15
$\bar{\nu}_e$	0.	stable	67	6	-12	-11		16
$\nu_{\mu}$	0.	stable	68	135	14	13		17
$\bar{\nu}_{\mu}$	0.	stable	69	136	-14	-13		18

Table 20: (continued) Particles used in different programs (to be continued).

<sup>126</sup>CORSIKA values; values in other programs may slightly differ.

Particle name	Particle mass <sup>126</sup> [ GeV ]	Particle life time <sup>126</sup> [ sec ]	Identification code					
			COR-SIKA	DPM-JET	PYTHIA HERWIG (PDG)	EPOS NEXUS	QGS-JET	SIB-YLL
$D^{\circ}$	1.8645	$4.101 \cdot 10^{-13}$	116	116	421	-140	8	71
$D^+$	1.8697	$1.040 \cdot 10^{-12}$	117	117	411	-240	7	59
$\overline{D}^-$	1.8697	$1.040 \cdot 10^{-12}$	118	118	-411	240	-7	60
$\overline{D}^{\circ}$	1.8645	$4.101 \cdot 10^{-13}$	119	119	-421	140	-8	72
$D_s^+$	1.9682	$5.00 \cdot 10^{-13}$	120	120	431	-340		74
$\overline{D}_s^-$	1.9682	$5.00 \cdot 10^{-13}$	121	121	-431	340		75
$\eta_c$	2.9804	$3.805 \cdot 10^{-23}$	122	122	441	440		73
$D^{*\circ}$	2.0067	$4. \cdot 10^{-22}$	123	123	423	-141		80
$D^{*+}$	2.0100	$6.86 \cdot 10^{-21}$	124	123	413	-241		78
$\overline{D}^{*-}$	2.0100	$6.86 \cdot 10^{-21}$	125	125	-413	241		79
$\overline{D}^{*\circ}$	2.0067	$4. \cdot 10^{-22}$	126	126	-423	141		81
$D_s^{*+}$	2.1121	$4. \cdot 10^{-22}$	127	127	433	-341		76
$\overline{D}_s^{*-}$	2.1121	$4. \cdot 10^{-22}$	128	128	-433	341		77
$J/\psi$	3.096916	$7.233 \cdot 10^{-21}$	130	130	443	441		83
$\tau^+$	1.77699	$2.906 \cdot 10^{-13}$	131	131	-15	-16		90
$\tau^-$	1.77699	$2.906 \cdot 10^{-13}$	132	132	15	16		91
$\nu_{\tau}$	0.	stable	133	133	16	15		92
$\overline{\nu}_{\tau}$	0.	stable	134	134	-16	-15		93
$\Lambda_c^+$	2.28646	$2.00 \cdot 10^{-13}$	137	137	4122	2140	9	89
$\Xi_c^+$	2.4679	$4.42 \cdot 10^{-13}$	138	138	4232	3140		87
$\Xi_c^{\circ}$	2.4710	$1.12 \cdot 10^{-13}$	139	139	4132	3240		88
$\Sigma_c^{++}$	2.45402	$2.95 \cdot 10^{-22}$	140	140	4222	1140		84
$\Sigma_c^+$	2.4529	$2. \cdot 10^{-22}$	141	141	4212	1240		85
$\Sigma_c^{\circ}$	2.45376	$3. \cdot 10^{-22}$	142	142	4112	2240		86
$\Xi_c'^+$	2.5757	$1. \cdot 10^{-23}$	143	143	4322	1340		
$\Xi_c'^{\circ}$	2.5780	$1. \cdot 10^{-23}$	144	144	4312	2340		
$\Omega_c^{\circ}$	2.6975	$6.9 \cdot 10^{-14}$	145	145	4332	3340		99
$\overline{\Lambda}_c^-$	2.28646	$2.00 \cdot 10^{-13}$	149	149	-4122	-2140	-9	-89
$\overline{\Xi}_c^-$	2.4679	$4.42 \cdot 10^{-13}$	150	150	-4232	-3140		-87
$\overline{\Xi}_c^{\circ}$	2.4710	$1.12 \cdot 10^{-13}$	151	151	-4132	-3240		-88
$\overline{\Sigma}_c^{--}$	2.45402	$2.95 \cdot 10^{-22}$	152	152	-4222	-1140		-84
$\overline{\Sigma}_c^-$	2.4529	$2. \cdot 10^{-22}$	153	153	-4212	-1240		-85
$\overline{\Sigma}_c^{\circ}$	2.45376	$3. \cdot 10^{-22}$	154	154	-4112	-2240		-86
$\overline{\Xi}_c'^-$	2.5757	$1. \cdot 10^{-23}$	155	155	-4322	-1340		
$\overline{\Xi}_c'^{\circ}$	2.5780	$1. \cdot 10^{-23}$	156	156	-4312	-2340		
$\overline{\Omega}_c^{\circ}$	2.6975	$6.9 \cdot 10^{-14}$	157	157	-4332	-3340		-99

Table 20: (continued) Particles used in different programs (to be continued).

Particle name	Particle mass <sup>126</sup> [ GeV ]	Particle life time <sup>126</sup> [ sec ]	Identification code					
			COR-SIKA	DPM-JET	PYTHIA HERWIG (PDG)	EPOS NEXUS	QGS-JET	SIB-YLL
$\Sigma_c^{*++}$	2.5184	$3.66*10^{-23}$	161	161	4224	1141		94
$\Sigma_c^{*+}$	2.5175	$4.*10^{-23}$	162	162	4214	1241		95
$\Sigma_c^{*\circ}$	2.5180	$5.06*10^{-23}$	163	163	4114	2241		96
$\overline{\Sigma}_c^{*-}$	2.5184	$3.66*10^{-23}$	171	171	-4224	1141		-94
$\overline{\Sigma}_c^{*0}$	2.5175	$4.*10^{-23}$	172	172	-4214	1241		-95
$\overline{\Sigma}_c^{*\circ}$	2.5180	$5.06*10^{-23}$	173	173	-4114	2241		-96
$B^\circ$	5.27961	$1.520*10^{-12}$	176		511			
$B^+$	5.27929	$1.638*10^{-12}$	177		521			
$\overline{B}^-$	5.27929	$1.638*10^{-12}$	178		-521			
$\overline{B}^\circ$	5.27961	$1.520*10^{-12}$	179		-511			
$B_s^\circ$	5.36679	$1.510*10^{-12}$	180		531			
$\overline{B}_s^\circ$	5.36679	$1.510*10^{-12}$	181		-531			
$B_c^+$	6.2751	$5.07*10^{-13}$	182		541			
$\overline{B}_c^-$	6.2751	$5.07*10^{-13}$	183		-541			
$\Lambda_b^\circ$	5.61951	$1.466*10^{-12}$	184		5122			
$\Sigma_b^-$	5.8155	$1.34*10^{-22}$	185		5112			
$\Sigma_b^+$	5.8113	$5.72*10^{-23}$	186		5222			
$\Xi_b^\circ$	5.7918	$1.464*10^{-12}$	187		5232			
$\Xi_b^-$	5.7944	$1.560*10^{-12}$	188		5132			
$\Omega_b^-$	6.0480	$1.57*10^{-12}$	189		5332			
$\overline{\Lambda}_b^\circ$	5.61951	$1.466*10^{-12}$	190		-5122			
$\overline{\Sigma}_b^+$	5.8155	$1.34*10^{-22}$	191		-5112			
$\overline{\Sigma}_b^-$	5.8113	$5.72*10^{-23}$	192		-5222			
$\overline{\Xi}_b^\circ$	5.7918	$1.464*10^{-12}$	193		-5232			
$\overline{\Xi}_b^+$	5.7944	$1.560*10^{-12}$	194		-5132			
$\overline{\Omega}_b^+$	6.0480	$1.57*10^{-12}$	195		-5332			

Table 20: (continued) Particles used in different programs.

## E Particle Arrays and Registers

current particle stack position	equivalence	mother particle stack position	grandmother particle stack position	meaning
CURPAR( 0)	(ITYPE)	SECPAR(17)	SECPAR(28)	particle identifier
CURPAR( 1)	GAMMA	SECPAR(18)	SECPAR(29)	$\gamma$ Lorentz factor (or energy)
CURPAR( 2)	COSTHE	SECPAR(19)	SECPAR(30)	$\cos(\theta)$ vertical direct. <sup>1</sup>
CURPAR( 3)	PHIX	SECPAR(20)	SECPAR(31)	$\sin(\theta) \cdot \cos(\phi)$ horizont. direct. <sup>1</sup>
CURPAR( 4)	PHIY	SECPAR(21)	SECPAR(32)	$\sin(\theta) \cdot \sin(\phi)$ horizont. direct. <sup>1</sup>
CURPAR( 5)	H	SECPAR(22)	SECPAR(33)	height [cm]
CURPAR( 6)	T	SECPAR(23)	SECPAR(34)	time [sec] (since first interaction)
CURPAR( 7)	X	SECPAR(24)		horizontal position [cm]
CURPAR( 8)	Y	SECPAR(25)		horizontal position [cm]
	(GEN)		SECPAR(35)	generation counter of mother particle
CURPAR( 9)	CHI		SECPAR(36)	$\chi$ depth to next interaction [g/cm <sup>2</sup> ]
CURPAR(10)	BETA			$\beta = v/c$ fraction of speed of light
CURPAR(11)	GCM			$\gamma$ Lorentz factor in cms-system
CURPAR(12)	ECM			energy in cms-system [GeV]
CURPAR(13)	WEIGHT	SECPAR(26)	SECPAR(37)	weight for thinning <sup>2</sup>
CURPAR(14)	HAPP			apparent height [cm] <sup>3</sup>
CURPAR(15)	COSTAP			$\cos(\theta^*)$ apparent zenith angle cosine <sup>3</sup>
CURPAR(16)	COSTEA			$\cos(\theta_E)$ angle at Earth center cosine <sup>3</sup>
CURPAR(17)				transverse momentum [GeV/c] <sup>4</sup>
CUTPAR(18)				random seed <sup>5</sup>
		SECPAR(27)	SECPAR(38)	(not used)
CURPAR(39)				ECT-flag <sup>5</sup>
CURPAR(40)				(not used)
CURPAR(41 - 46)				weights for different thin modes <sup>6</sup>

<sup>1</sup>  $\theta$  = zenith angle,  $\phi$  = azimuth angle  
<sup>2</sup> THIN option  
<sup>3</sup> CURVED option  
<sup>4</sup> INTTEST option  
<sup>5</sup> PARALLEL option  
<sup>6</sup> MULTITHIN option

Table 21: Meaning of particle array positions.

register / stack name	number of particles	content	located in
PRMPAR	1	primary particle	common PARPAR
CURPAR	1	current particle	common PARPAR
PROPAR	1	propagated particle	AAMAIN, BOX3, MUTRAC
SECPAR	1	secondary particle	common PARPAR
OUTPAR	1	output / propagated particle	common PARPAR
AMUPAR	1	additional muon info	common MUPART
CUTPAR	1	cutted particle <sup>1</sup>	common PARPAR
STACKINT	200 000	intermediate stack for thinning	common STACKINT
STACKI	680 624 <sup>2</sup> 512 <sup>3</sup> 50 000 <sup>1</sup>	intermediate stack for external file	common STACKF
STACKJ	50 000 <sup>1</sup>	intermediate stack	common STACKF
(MEXST)	$\infty$	particle stack	(scratch file)
DATAB	39	output buffer	common BUFFS
(STACKE)	60	em-particle stack	common STACKE
<sup>1</sup> PARALLEL option <sup>2</sup> THIN option <sup>3</sup> CURVED option			

Table 22: Particle registers / stacks with sizes.



## F Atmospheres

The atmosphere adopted consists of  $N_2$ ,  $O_2$ , and  $Ar$  with the volume fractions of 78.1%, 21.0%, and 0.9% [81]. The density variation of the atmosphere with altitude is modeled by 5 layers. In the lower four of them the density follows an exponential dependence on the altitude leading to a relation between the mass overburden  $T(h)$  of the atmosphere and the height  $h$  of the form

$$T(h) = a_i + b_i \cdot e^{-h/c_i} \quad i = 1, \dots, 4 \quad . \quad (1)$$

In the fifth layer the mass overburden decreases linearly with height

$$T(h) = a_5 - b_5 \cdot h/c_5 \quad .$$

The boundary of the atmosphere in this model is defined at the height where the mass overburden  $T(h)$  vanishes (which is at  $h = 112.8$  km for the U.S. standard atmosphere). Various atmospheres are foreseen: U.S. standard atmosphere parameterized according to J. Linsley [82], 7 typical atmospheres as measured above Stuttgart (about 60 km away from Karlsruhe) at various days of 1993 and transmitted by Deutscher Wetterdienst Offenbach (parameterized according to Ref. [83]), 4 South pole atmospheres (parameterized by D. Chirkin according to the MSIS-90-E model [84]), two South pole atmospheres by P. Lipari [85], and 12 monthly dependent atmospheres for the Pierre Auger Observatory experiment at Malargüe (Argentina) parameterized with GDAS-data by M. Will and B. Keilhauer [86, 87]. B. Keilhauer provided also a new parameterization of the U.S. standard atmosphere. For the South Pole the atmospheric parameters used for the IceTop experiment are available in 12 monthly dependent sets by S. de Ridder [88]. The parameters  $a_i$ ,  $b_i$ , and  $c_i$  are selected in a manner that the function  $T(h)$  is continuous at the layer boundaries and can be differentiated continuously. In Tables 23 - 49 the parameters for the various models are listed. Additional atmospheres [45] are available by the keyword ATMOSPHERE (page 83) in the ATMEXT option (page 48). User specific atmosphere parameters may be read in using the keywords ATMOD, ATMA, ATMB, ATMC, and ATMLAY.

### U.S. Standard Atmosphere

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-186.555305	1222.6562	994186.38
2	4 ... 10	-94.919	1144.9069	878153.55
3	10 ... 40	0.61289	1305.5948	636143.04
4	40 ... 100	0.0	540.1778	772170.16
5	> 100	0.01128292	1	10 <sup>9</sup>

Table 23: Parameters of the U.S. standard atmosphere (after Linsley).

## Middle Europe

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-118.1277	1173.9861	919546.
2	4 ... 10	-154.258	1205.7625	963267.92
3	10 ... 40	0.4191499	1386.7807	614315.
4	40 ... 100	$5.4094056 \cdot 10^{-4}$	555.8935	739059.6
5	> 100	0.01128292	1	$10^9$

Table 24: Parameters of the AT115 atmosphere (January 15, 1993).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-195.837264	1240.48	933697.
2	4 ... 10	-50.4128778	1117.85	765229.
3	10 ... 40	0.345594007	1210.9	636790.
4	40 ... 100	$5.46207 \cdot 10^{-4}$	608.2128	733793.8
5	> 100	0.01128292	1	$10^9$

Table 25: Parameters of the AT223 atmosphere (February 23, 1993).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-253.95047	1285.2782	1088310.
2	4 ... 10	-128.97714	1173.1616	935485.
3	10 ... 40	0.353207	1320.4561	635137.
4	40 ... 100	$5.526876 \cdot 10^{-4}$	680.6803	727312.6
5	> 100	0.01128292	1	$10^9$

Table 26: Parameters of the AT511 atmosphere (May 11, 1993).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-208.12899	1251.474	1032310.
2	4 ... 10	-120.26179	1173.321	925528.
3	10 ... 40	0.31167036	1307.826	645330.
4	40 ... 100	$5.591489 \cdot 10^{-4}$	763.1139	720851.4
5	> 100	0.01128292	1	$10^9$

Table 27: Parameters of the AT616 atmosphere (June 16, 1993).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-77.875723	1103.3362	932077.
2	4 ... 10	-214.96818	1226.5761	1109960.
3	10 ... 40	0.3721868	1382.6933	630217.
4	40 ... 100	$5.5309816 \cdot 10^{-4}$	685.6073	726901.3
5	> 100	0.01128292	1	$10^9$

Table 28: Parameters of the AT822 atmosphere (August 22, 1993).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-242.56651	1262.7013	1059360.
2	4 ... 10	-103.21398	1139.0249	888814.
3	10 ... 40	0.3349752	1270.2886	639902.
4	40 ... 100	$5.527485 \cdot 10^{-4}$	681.4061	727251.8
5	> 100	0.01128292	1	$10^9$

Table 29: Parameters of the AT1014 atmosphere (October 14, 1993).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-195.34842	1210.4	970276.
2	4 ... 10	-71.997323	1103.8629	820946.
3	10 ... 40	0.3378142	1215.3545	639074.
4	40 ... 100	$5.48224 \cdot 10^{-4}$	629.7611	731776.5
5	> 100	0.01128292	1	$10^9$

Table 30: Parameters of the AT1224 atmosphere (December 24, 1993).

## South pole

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0... 4	-137.656	1130.74	867358.
2	4... 10	-37.9610	1052.05	741208.
3	10... 40	0.222659	1137.21	633846.
4	40... 100	-0.000616201	442.512	759850.
5	> 100	0.00207722	1	$5.4303203 \cdot 10^9$

Table 31: Parameters of South pole atmosphere for March 31, 1997 (MSIS-90-E).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0... 4	-163.331	1183.70	875221.
2	4... 10	-65.3713	1108.06	753213.
3	10... 40	0.402903	1424.02	545846.
4	40... 100	-0.000479198	207.595	793043.
5	> 100	0.00188667	1	$5.9787908 \cdot 10^9$

Table 32: Parameters of South pole atmosphere for Jul. 01, 1997 (MSIS-90-E).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0... 4	-142.801	1177.19	861745.
2	4... 10	-70.1538	1125.11	765925.
3	10... 40	1.14855	1304.77	581351.
4	40... 100	-0.000910269	433.823	775155.
5	> 100	0.00152236	1	$7.4095699 \cdot 10^9$

Table 33: Parameters of South pole atmosphere for Oct. 01, 1997 (MSIS-90-E).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0 ... 4	-128.601	1139.99	861913.
2	4 ... 10	-39.5548	1073.82	744955.
3	10 ... 40	1.13088	1052.96	675928.
4	40 ... 100	-0.00264960	492.503	829627.
5	> 100	0.00192534	1	$5.8587010 \cdot 10^9$

Table 34: Parameters of South pole atmosphere for Dec. 31, 1997 (MSIS-90-E).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 2.67	-113.139	1133.10	861730.
2	2.67 ... 5.33	-79.0635	1101.20	826340.
3	5.33 ... 8.0	-54.3888	1085.00	790950.
4	8.0 ... 100.0	0.0000	1098.00	682800.
5	> 100.0	0.00421033	1	$2.6798156 \cdot 10^9$

Table 35: Parameters of South pole atmosphere for January (Lipari).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 6.67	-59.0293	1079.00	764170.
2	6.67 ... 13.33	-21.5794	1071.90	699910.
3	13.33 ... 20.0	-7.14839	1182.00	635650.
4	20.0 ... 100.0	0.0000	1647.10	551010.
5	> 100.0	0.000190175	1	$59.329575 \cdot 10^9$

Table 36: Parameters of South pole atmosphere for August (Lipari).

### Keilhauer's U.S. standard atmosphere

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 7.0	-149.801663	1183.6071	954248.34
2	7.0 ... 11.4	-57.932486	1143.0425	800005.34
3	11.4 ... 37.0	0.63631894	1322.9748	629568.93
4	37.0 ... 100.0	$4.35453690 \cdot 10^{-4}$	655.67307	737521.77
5	> 100.0	0.01128292	1.	$10^9$

Table 37: Parameters of the U.S. standard atmosphere (after Keilhauer).

### Malargüe (Argentina)

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.4	-136.72575606	1174.8298334	982815.95248
2	9.4 ... 15.3	-31.636643044	1204.8233453	754029.87759
3	15.3 ... 31.6	1.8890234035	1637.7703583	594416.83822
4	31.6 ... 100.0	$3.9201867984 \cdot 10^{-4}$	735.96095023	733974.36972
5	> 100.0	0.01128292	1.	$10^9$

Table 38: Parameters of January (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.2	-137.25655862	1176.0907565	981369.6125
2	9.2 ... 15.4	-31.793978896	1197.8951104	756657.65383
3	15.4 ... 31.0	2.0616227547	1646.4616955	592969.89671
4	31.0 ... 100.0	$4.1243062289 \cdot 10^{-4}$	755.18728657	731345.88332
5	> 100.0	0.01128292	1.	$10^9$

Table 39: Parameters of February (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.6	-132.36885162	1172.6227784	972654.0563
2	9.6 ... 15.2	-29.077046629	1215.3964677	742769.2171
3	15.2 ... 30.7	2.090501509	1617.0099282	595342.19851
4	30.7 ... 100.0	$4.3534337925 \cdot 10^{-4}$	769.51991638	728921.61954
5	> 100.0	0.01128292	1.	$10^9$

Table 40: Parameters of March (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 10.0	-129.9930412	1172.3291878	962396.5521
2	10.0 ... 14.9	-21.847248438	1250.2922774	711452.06673
3	14.9 ... 32.6	1.5211136484	1542.6248413	603480.61835
4	32.6 ... 100.0	$3.9559055121 \cdot 10^{-4}$	713.1008285	735460.83741
5	> 100.0	0.01128292	1.	$10^9$

Table 41: Parameters of April (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 10.2	-125.11468467	1169.9511302	947742.88769
2	10.2 ... 15.1	-14.591235621	1277.6768488	685089.57509
3	15.1 ... 35.9	0.93641128677	1493.5303781	609640.01932
4	35.9 ... 100.0	$3.2475590985 \cdot 10^{-4}$	617.9660747	747555.95526
5	> 100.0	0.01128292	1.	$10^9$

Table 42: Parameters of May (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 10.1	-126.17178851	1171.0916276	940102.98842
2	10.1 ... 16.0	-7.7289852811	1295.3516434	661697.57543
3	16.0 ... 36.7	0.81676828638	1455.3009344	612702.0632
4	36.7 ... 100.0	$3.1947676891 \cdot 10^{-4}$	595.11713507	749976.26832
5	> 100.0	0.01128292	1.	$10^9$

Table 43: Parameters of June (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.6	-126.17216789	1172.7340688	934649.58886
2	9.6 ... 16.5	-8.6182537514	1258.9180079	672975.82513
3	16.5 ... 37.4	0.74177836911	1450.0537141	614888.52458
4	37.4 ... 100.0	$2.9350702097 \cdot 10^{-4}$	583.07727715	752631.28536
5	> 100.0	0.01128292	1.	$10^9$

Table 44: Parameters of July (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.6	-123.27936204	1169.763036	931569.97625
2	9.6 ... 15.9	-10.051493041	1251.0219808	678861.75136
3	15.9 ... 36.3	0.84187346153	1436.6499372	617363.34491
4	36.3 ... 100.0	$3.2422546759 \cdot 10^{-4}$	627.42169844	746739.16141
5	> 100.0	0.01128292	1.	$10^9$

Table 45: Parameters of August (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.5	-126.94494665	1174.8676453	936953.91919
2	9.5 ... 16.2	-9.5556536981	1251.5588529	678906.60516
3	16.2 ... 37.2	0.74939405052	1440.8257549	618132.60561
4	37.2 ... 100.0	$2.9823116961 \cdot 10^{-4}$	606.31473165	750154.67709
5	> 100.0	0.01128292	1.	$10^9$

Table 46: Parameters of September (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.5	-133.13151125	1176.9833473	954151.404
2	9.5 ... 15.5	-13.973209265	1244.234531	692708.89816
3	15.5 ... 36.5	0.8378263431	1464.0120855	615439.43936
4	36.5 ... 100.0	$3.111742176 \cdot 10^{-4}$	622.11207419	747969.08133
5	> 100.0	0.01128292	1.	$10^9$

Table 47: Parameters of October (GDAS).



Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.6	-134.72208165	1175.7737972	964877.07766
2	9.6 ... 15.3	-18.172382908	1238.9538504	706199.57502
3	15.3 ... 34.6	1.1159806845	1505.1614366	610242.24564
4	34.6 ... 100.0	$3.5217025515 \cdot 10^{-4}$	670.64752105	741412.74548
5	> 100.0	0.01128292	1.	$10^9$

Table 48: Parameters of November (GDAS).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 9.6	-135.40825209	1174.644971	973884.44361
2	9.6 ... 15.6	-22.830409026	1227.2753683	723759.74682
3	15.6 ... 33.3	1.4223453493	1585.7130562	600308.13983
4	33.3 ... 100.0	$3.7512921774 \cdot 10^{-4}$	691.23389637	738390.20525
5	> 100.0	0.01128292	1.	$10^9$

Table 49: Parameters of December (GDAS).

### South Pole (for IceTop)

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 7.8	-91.6956	1125.71	821621.
2	7.8 ... 16.4	7.01491	1149.81	635444.
3	16.4 ... 40.4	0.505452	1032.68	682968.
4	40.4 ... 100.0	-.00181302	490.789	807327.
5	> 100.0	0.00207722	1.	5430320300.

Table 50: Parameters of South Pole January (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.0	-72.1988	1108.19	786271.
2	8.0 ... 10.6	22.7002	1159.77	599986.
3	10.6 ... 40.4	0.430171	1079.25	667432.
4	40.4 ... 100.0	-.0012030	523.956	780919.
5	> 100.0	0.00207722	1.	5430320300.

Table 51: Parameters of South Pole February (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 6.7	-63.7290	1102.66	764831.
2	6.7 ... 22.4	-1.02799	1093.56	660389.
3	22.4 ... 40.4	0.324414	1198.93	636118.
4	40.4 ... 100.0	-.000490772	589.827	734909.
5	> 100.0	0.00207722	1.	5430320300.

Table 52: Parameters of South Pole March (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 7.6	-69.7259	1111.70	766099.
2	7.6 ... 22.0	-2.79781	1128.64	641716.
3	22.0 ... 40.4	0.262692	1413.98	588082.
4	40.4 ... 100.0	-.0000841695	587.688	693300.
5	> 100.0	0.00207722	1.	5430320300.

Table 53: Parameters of South Pole April (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.4	-78.5551	1118.46	776648.
2	8.4 ... 20.0	-5.33239	1169.09	626683.
3	20.0 ... 39.7	0.312889	1577.71	553087.
4	39.7 ... 100.0	-.0000920472	452.177	696835.
5	> 100.0	0.00152236	1.	7409569900.

Table 54: Parameters of South Pole May (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.5	-92.6125	1129.88	791177.
2	8.5 ... 17.9	-8.56450	1191.98	618840.
3	17.9 ... 38.4	0.363986	1619.82	535235.
4	38.4 ... 100.0	.00207722	411.586	692253.
5	> 100.0	0.00152236	1.	5430320300.

Table 55: Parameters of South Pole June (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.5	-89.9639	1125.73	784553.
2	8.5 ... 15.9	-13.9697	1180.47	628042.
3	15.9 ... 37.5	0.441631	1581.43	531652.
4	37.5 ... 100.0	-0.0000146525	373.796	703417.
5	> 100.0	0.00207722	1.	5430320300.

Table 56: Parameters of South Pole July (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.5	-90.4253	1125.01	781628.
2	8.5 ... 14.4	-18.7154	1175.60	633793.
3	14.4 ... 37.5	0.513930	1518.03	533269.
4	37.5 ... 100.0	-0.00021565	299.006	737794.
5	> 100.0	0.00152236	1.	7409569900.

Table 57: Parameters of South Pole August (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.5	-91.6860	1125.53	786017.
2	8.5 ... 13.0	-23.3519	1169.77	645241.
3	13.0 ... 36.2	0.891302	1431.26	545022.
4	36.2 ... 100.0	-0.000765666	247.030	805419.
5	> 100.0	0.00207722	1.	5430320300.

Table 58: Parameters of South Pole September (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 3.1	451.616	849.239	225286.
2	3.1 ... 10.1	-85.5456	1113.16	789340.
3	10.1 ... 31.5	2.06082	1322.28	566132.
4	31.5 ... 100.0	-0.0010760	372.242	796434.
5	> 100.0	0.00207722	1.	5430320300.

Table 59: Parameters of South Pole October (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.5	-152.853	1174.09	891602.
2	8.5 ... 22.4	4.22741	1272.49	582119.
3	22.4 ... 32.4	1.38352	975.906	643130.
4	32.4 ... 100.0	-.00115014	481.615	783786.
5	> 100.0	0.00207722	1.	5430320300.

Table 60: Parameters of South Pole November (De Ridder).

Layer $i$	Altitude $h$ (km)	$a_i$ (g/cm <sup>2</sup> )	$b_i$ (g/cm <sup>2</sup> )	$c_i$ (cm)
1	0. ... 8.5	-100.386	1128.71	829352.
2	8.5 ... 22.0	5.43849	1198.10	612649.
3	22.0 ... 40.4	0.399465	858.522	706104.
4	40.4 ... 100.0	-.00175472	480.142	806875.
5	> 100.0	0.00207722	1.	5430320300.

Table 61: Parameters of South Pole December (De Ridder).

## G Atmospheric profiles from GDAS using gdastool

The `gdastool` script included with CORSIKA (in `src/utils/` subdirectory) downloads a section of the GDAS (Global Data Assimilation System) database and compiles a file containing Corsika settings ATMLAY, ATMA, ATMB, ATMC, and an altitude profile of the refractive index.

It takes as input:

- an observatory name (currently, valid names are “lofar” or “aera” or observatory latitude / longitude coordinates in degrees
- a UTC timestamp (seconds since Jan 1, 1970) for which to read out the atmospheric profile
- an output path and filename for the compiled profile
- an output path “gdaspath” for storing the downloaded section of GDAS.

Given coordinates and time stamp are rounded to the nearest 1 x 1 degree grid point, and to the nearest 3-hour time in GDAS.

To use the output file in CORSIKA, it should be included into the steering file with the `ATMFILE` keyword (see Sect. 4.23 page 83).

The full usage info for the script is :

```
usage: gdastool [-h] [-t UTCTIMESTAMP] [-o OUTPUT]
(--observatory \{lofar,aera\} | -c COORDINATES COORDINATES)
[-m MINHEIGHT] [-s INTERPOLATIONSTEPS] [-p GDASPATH] [-v] [-g]
```

Creates an atmosphere profile for CORSIKA/CoREAS from GDAS data. Downloads GDAS model data for the defined location and time and fits a 5-layer model of the atmosphere to the data. Based on the fit, a table for the refractive index is created for usage in CoREAS.

optional arguments:

```
-h, --help          show this help message and exit
-t UTCTIMESTAMP, --utctimestamp UTCTIMESTAMP
    UTC time stamp of the event
-o OUTPUT, --output OUTPUT
    Name of the outputfile.
--observatory \{lofar,aera\}
    Preset of observatory coordinates.
-c COORDINATES COORDINATES, --coordinates COORDINATES \
    COORDINATES
```

Coordinates of the observatory lat=-90..90 lon=0..360  
in deg, e.g. --coordinates 50.85 4.25 for Brussels.

-m MINHEIGHT, --minheight MINHEIGHT  
Minimum height for the interpolation. Default is -1.0 km.

-s INTERPOLATIONSTEPS, --interpolationSteps \ INTERPOLATIONSTEPS  
Step length for interpolation. Default is 1 m.

-p GDASPATH, --gdaspath GDASPATH  
path to local gdas file directory. If required file is not there, it will be downloaded.

-v, --verbose Set log level, -vv for debug.

-g, --createplot plot density profile.

Please contact

Pragati Mitra <pragati9163@gmail.com>,  
Arthur Corstanje <a.corstanje@astro.ru.nl> or  
Tobias Winchen <tobias.winchen@rwth-aachen.de>  
in case of questions or bugs.

Note that to run `gdastool`, Python version 2.7 or higher with the package `python-scipy` and `python-numpy` (and optionally `python-matplotlib`) should be installed on the system.

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